## Advanced techniques and indirect methods

#### G.L. Guardo

on behalf of the AsFiN collaboration



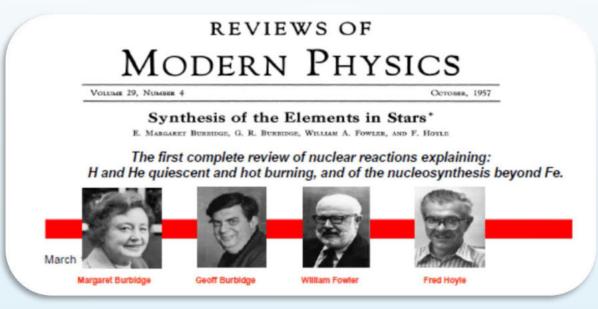




### **Experimental Nuclear Astrophysics**

... Everything starts from the B<sup>2</sup>FH review paper of 1957, the basis of the modern nuclear astrophysics

this work has been considered as the greatest gift of astrophysics to modern civilization



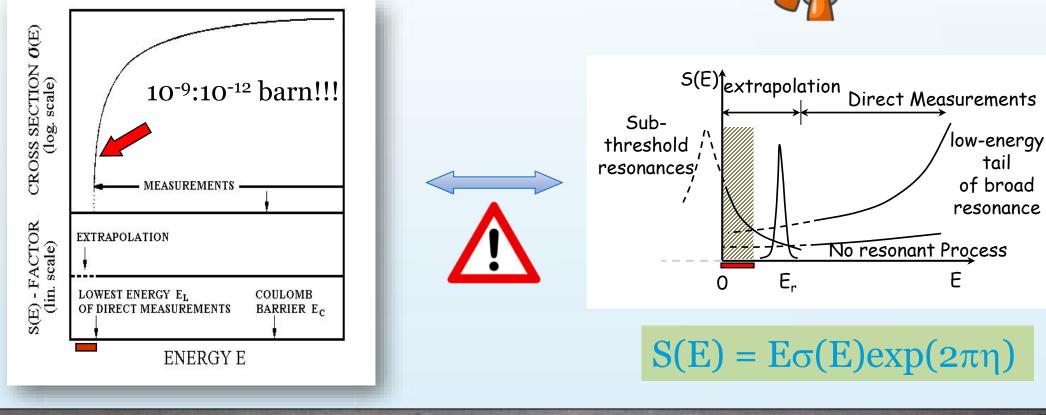
The elements composing everything from planets to life were forged inside earlier generations of stars!

Nuclear reactions responsible for both <u>ENERGY PRODUCTION</u> and <u>SYNYHESIS OF ELEMENTS</u>



### **Direct Measurements**

- Very small cross section values reflect in a faint statistic;
- Very low signal-to-noise ratio makes hard the investigation at astrophysical energies;
- > Instead of the cross section, the S(E)-factor is introduced



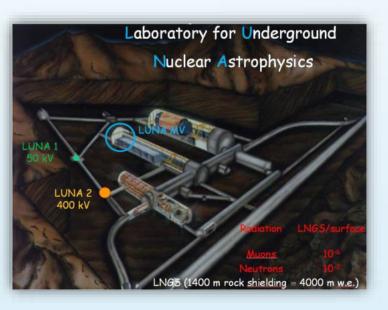


### **Direct Measurements**



- > Longer measurements
- > Higher beam currents
- > 4 $\pi$  detectors
- > Pure targets
- > Underground laboratories

Several efforts have been made in the last years in order to **improve the signal-tonoise ratio** for low-energy cross section measurement.





### Electron screening

ELECTRON

CLOUD

Ra

Due to the electron cloud surrounding the interacting ions the projectile feels a reduced barrier

> See A. Cvetinovic & M. Lipoglavsek talks

COULOMB

BARE

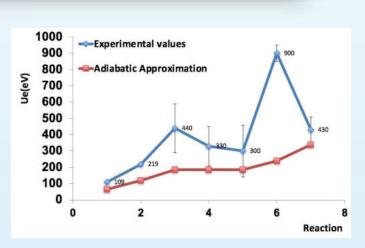
R<sub>n</sub>

NUCLEUS

SHIELDED NUCLEUS

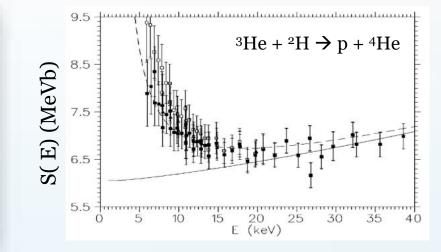
DISTANCE r

Reaction	Uad (eV)	$U_{exp}$ (eV)	Reference		
$^{6}\mathrm{Li}(\mathrm{p},\alpha)^{3}\mathrm{He}$	186	$440 \pm 150$	[Engstler et al.(1992)]		
$^{6}\text{Li}(d,\alpha)^{4}\text{He}$	186	$330 \pm 120$	[Engstler et al.(1992)]		
$H(^{7}Li,\alpha)^{4}He$	186	$300 \pm 160$	[Engstler et al.(1992)]		
<sup>2</sup> H( <sup>3</sup> He,p) <sup>4</sup> He	65	109±9	[Aliotta et al.(2004)]		
<sup>3</sup> He( <sup>2</sup> H,p) <sup>4</sup> He	120	219±7	[Aliotta et al.(2004)]		
$H(^{9}Be,\alpha)^{6}Li$	240	$900 \pm 50$	[Zahnow et al.(1997)]		
$H(^{11}B,\alpha)^8Be$	340	$430 \pm 80$	[Angulo et al. (1993)]		



PROJECTILE

Rc



Theory.vs.Experiment→ Far to be understood... <u>Stellar Plasma</u>



## Indirect Methods

#### \* Coulomb dissociation

G. Baur et al. Annu. Rev. Nucl. Part. Sci. 46,321,(1996) to determine the absolute S(E) factor of a radiative capture reaction  $A+x \rightarrow B+\gamma$ studying the reversing photodisintegration process  $B+\gamma \rightarrow A+x$ 

#### \* Asymptotic Normalization Coefficients (ANC)

A.M. Mukhamedzhanov et al.: PRC 56,1302,(1997) to determine the S(0) factor of the radiative capture reaction,  $A+x \rightarrow B+\gamma$  studying a peripheral transfer reaction into a bound state of the **B** nucleus

#### \* Trojan Horse Method (THM)

C. Spitaleri, *Problems of Fundamental Modern Physics, II,* (World Sci.,1991) C. Spitaleri et al., Phys. of Atomic Nuclei, 74 (2011) 1725

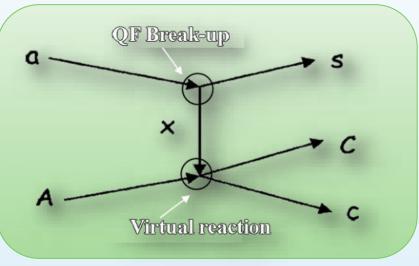
to determine the S(E) factor of a charged particle reaction  $A+x\rightarrow c+C$ 



### The Trojan Horse Method



The idea of the **THM** is to extract the cross section of an astrophysically relevant two-body reaction  $A+x \rightarrow c+C$  at low energies from a suitable three-body reaction  $a+A \rightarrow c+C+s$ 



**Quasi free kinematics is selected** 

- $\checkmark$  only x A interaction
- $\checkmark$  *s* = spectator (p<sub>s</sub>~0)

 $E_A > E_{Coul} \rightarrow$ 

- NO coulomb suppression
- NO electron screening
- NO centrifugal barrier

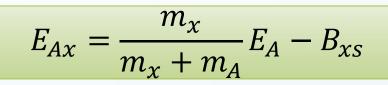
- THM Review paper  $\rightarrow$ 

Spitaleri C. et al., Prob. of Fund. Mod. Phys., 1991 Tumino A. et al., An. Rev. Nuc. and Part. Sci. 2021

### **Theoretical Approach**

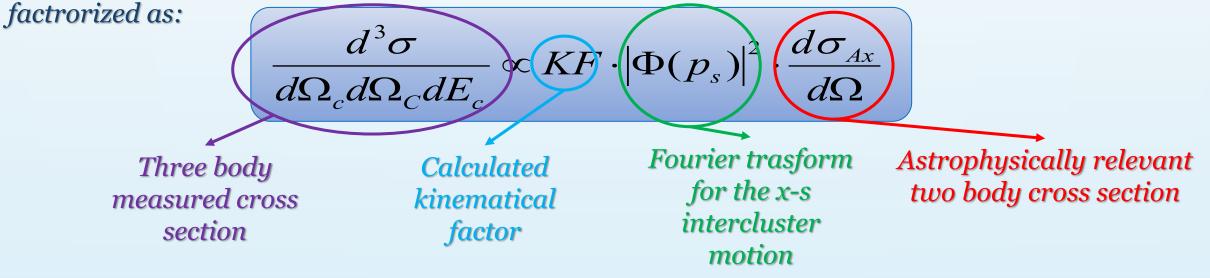
The TH-nucleus is chosen because of:

- its large amplitude in the  $a=x \oplus s$  cluster configuration;
- its relatively low-binding energy;
- Its known *x-s* momentum distribution  $|\Phi(p_s)|^2$  in *a*.



**B**<sub>x-s</sub> plays a key role in compensating for the beam energy thanks to the *x-s intercluster motion* inside a, it is possible to span an energy range of several hundreds of keV with <u>only one beam energy</u>

In the <u>Plane Wave Impulse Approximation</u> (PWIA) the cross section of the three body reaction can be



### **TH cross section**

#### Virtual nature of x particle $\rightarrow$ A+x interaction is off-energy shell

Cross section of the bare nucleus but NO <u>absolute value</u>  $\rightarrow$  <u>normalization to direct data available at higher energies</u>

Standard R-Matrix approach cannot be applied to extract the resonance parameters  $\rightarrow$  Modified R-Matrix is introduced instead

$$\frac{\mathrm{d}^2 \sigma}{\mathrm{d} E_{xA} \mathrm{d} \Omega_s} = \mathrm{NF} \sum_i (2\mathrm{J}_i + 1)$$
$$\times \left| \sqrt{\frac{\mathrm{k}_{\mathrm{f}}(E_{xA})}{\mu_{cC}}} \frac{\sqrt{2P_{l_i}(k_{cC}R_{cC})}M_i(p_{xA}R_{xA})\gamma^i_{cC}\gamma^i_{xA}}}{D_i(E_{xA})} \right|^2$$

where:

- $M_i(p_{xa}R_{xa})$  describes the transfer amplitude for the QF-process;
- $\gamma_{xa}$  and  $\gamma_{Cc}$  represents the reduced partial widths for the resonant excited states that are the <u>same</u> of the direct measurements

## <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study Astrophysical Scenario

#### Weak component s-process

<sup>17</sup>O(n,a)<sup>14</sup>C and <sup>17</sup>O(a,n)<sup>20</sup>Ne since they act as a neutron poison and a recycle channel during s-process nucleosinthesys in massive stars (M>8M<sub>SUN</sub>)



## <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study Astrophysical Scenario

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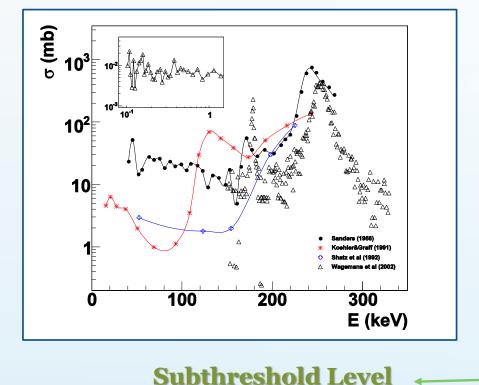


Temperature → 0.8<T<sub>8</sub><11 K Energy range→ ~0-100 keV



## <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study

• Status of the art



Suppressed due to the centrifugal barrier

**Available in literature** 

• R. M. Sanders, Phys. Rev., 104, 1434 (1956) INVERSE REACTION <sup>14</sup>C(α,n)<sup>17</sup>O

\* P.E.Koehler & S.M.Graff, Phys. Rev., C44(6), 2788 (1991)

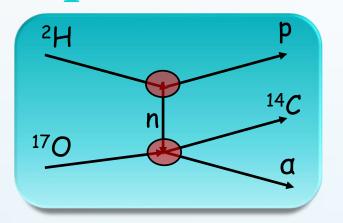
- H. Schatz et al., Astroph. J., 413, 750 (1993)
- ∆ J. Wagemans et al., Phys. Rev., C65(3), 34614 (2002)

E <sub>c.m.</sub> (keV)	<sup>18</sup> O* (MeV)	Jπ	
 -7	8.039	1-	
75	8.125	5⁻	
166	8.213	2+	
 236	8.282	3-	

F. Ajzenberg-Selove, Nucl. Phys., A475, 1 (1987)

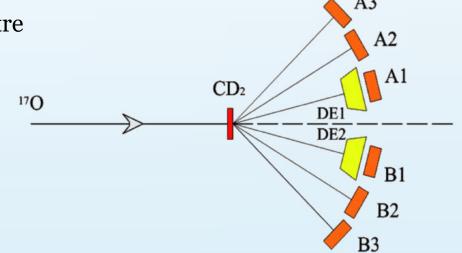


# <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study Experimental setup



The reaction <sup>17</sup>O(n,α)<sup>14</sup>C was studied via the <sup>2</sup>H(<sup>17</sup>O,α<sup>14</sup>C)p, V<sub>coul</sub>=2.3 MeV;
 The deuteron is the TH nucleus. Strong cluster n+p; B=2.2 MeV, |p<sub>s</sub>|=0 MeV/c.

- ✓ Experiments performed at ISNAP at the University of Notre Dame (USA) and LNS of Catania;
- ✓  $E_{\text{beam}}(^{17}\text{O})=43.5 \text{ MeV};$
- ✓ Target thickness  $CD_2 \sim 150 \ \mu g/cm^2$ ;
- $\checkmark\,$  IC filled with ~50 mbar isobutane gas;
- $\checkmark\,$  Angular position to cover the QF angular region
- ✓ Symmetric set-up in order to increase the statistic.

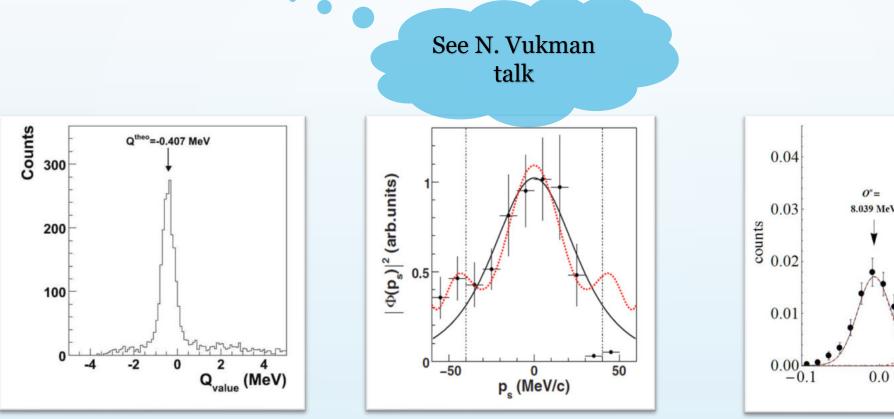




### <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study

Data reduction & analisys

» Channel selection



» *QF* mechanism selection

» Three-body yield determination

0.1

Ecm (MeV)

0"=

8.125 MeV

0"=

8.213 MeV

 $O^* =$ 

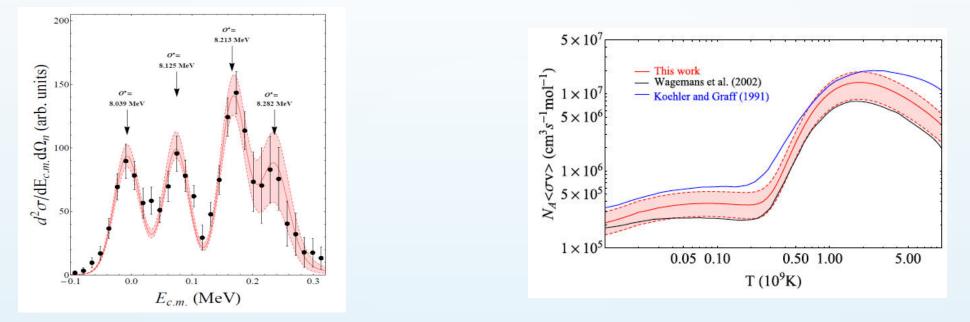
8.282 MeV

0.2

As Fi N

0.3

## <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study Data results



$\Gamma_{n}$ (eV)	$\Gamma_{\alpha}$ (eV)	$\Gamma_{\rm TOT}$ (eV)	Γ <sub>wag.</sub> (eV)
$0,01\pm0,001$	2362±307	2362±307	2400
0,05±0,006	36±5	36±5	-
86±11	2171±282	2257±293	2258±135
1714±446	13021±3386	14735±3832	14739±590
	0,01±0,001 0,05±0,006 86±11	0,01±0,001     2362±307       0,05±0,006     36±5       86±11     2171±282	0,01±0,001     2362±307     2362±307       0,05±0,006     36±5     36±5       86±11     2171±282     2257±293

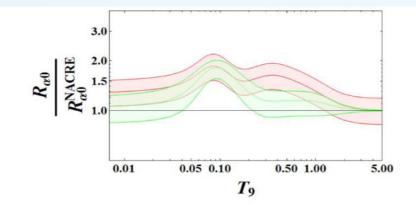
Guardo et al., Phys. Rev. C, 95, 025807, 2017



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#### Three open channels:

- ${}^{19}F(p,\alpha_0){}^{16}O$
- ${}^{19}F(p, \alpha_{\pi}){}^{16}O$
- ${}^{19}F(p, \alpha_{\gamma}){}^{16}O$  —



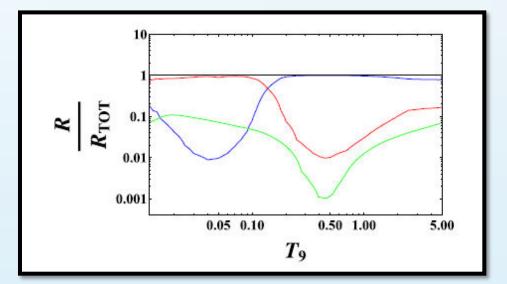
New Improved Indirect Measurement of the <sup>19</sup>F(p,  $\alpha$ )<sup>16</sup>O Reaction at Energies of

**Astrophysical Relevance** 

I. Indelicato<sup>1</sup>, M. La Cognata<sup>1</sup>, C. Spitaleri<sup>1,2</sup>, V. Burjan<sup>3</sup>, S. Cherubini<sup>1,2</sup>, M. Gulino<sup>1,4</sup>, S. Hayakawa<sup>5</sup>, Z. Hons<sup>3</sup>, V. Kroha<sup>3</sup>, L. Lamia<sup>1</sup>, M. Mazzocco<sup>6,7</sup>, J. Mrazek<sup>3</sup>, R. G. Pizzone<sup>1</sup>, S. Romano<sup>1,2</sup>, E. Strano<sup>6,7</sup>, D. Torresi<sup>6,7</sup>, and A. Tumino<sup>1,4</sup> <sup>1</sup>INFN, Laboratori Nazionali del Sud, Catania, Italy: indelicato@Ins.infn.it <sup>2</sup>Dipartimento di Fisica e Astronomia dell'Università degli studi di Catania, Catania, Italy

<sup>3</sup> Nuclear Physics Institute of ASCR, Rez near Prague, Czech Republic
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 <sup>6</sup> INFN, Sezione di Padova, Padova, Italy
 <sup>7</sup> Dipartimento di Fisica dell'Università di Padova, Padova, Italy
 *Received 2017 May 30; revised 2017 July 3; accepted 2017 July 3; published 2017 August 8*

**Figure 15.** Ratio of the reaction rate calculation obtained from the THM astrophysical factor (red band) to the rate recommended in NACRE (Angulo et al. 1999).  $T_9$  is the temperature in GK ( $T_9 = T/10^9$  K). The black line corresponds to  $R/R_{\alpha0}^{nacre} = 1$ . For comparison, the  $R/R_{\alpha0}^{nacre}$  ratio given in La Cognata et al. (2015) is shown as a green band.





THE ASTROPHYSICAL JOURNAL, 860:61 (11pp), 2018 June 10  $\odot$  2018. The American Astronomical Society. All rights reserved.

https://doi.org/10.3847/1538-4357/aac207

#### The <sup>19</sup>F( $\alpha$ , p)<sup>22</sup>Ne Reaction at Energies of Astrophysical Relevance by Means of the Trojan Horse Method and Its Implications in AGB Stars

G. D'Agata<sup>1,2</sup>, R. G. Pizzone<sup>1</sup>, M. La Cognata<sup>1</sup>, I. Indelicato<sup>1</sup>, C. Spitaleri<sup>1,2</sup>, S. Palmerini<sup>3,4</sup>, O. Trippella<sup>3,4</sup>, D. Vescovi<sup>3,4</sup>, S. Blagus<sup>5</sup>, S. Cherubini<sup>1,2</sup>, P. Figuera<sup>1</sup>, L. Grassi<sup>5</sup>, G. L. Guardo<sup>1</sup>, M. Gulino<sup>1,6</sup>, S. Hayakawa<sup>1,7</sup>, R. Kshetri<sup>1,8</sup>, L. Lamia<sup>1,2</sup>, M. Lattuada<sup>1,2</sup>, T. Mijatovic<sup>5</sup>, M. Milin<sup>9</sup>, D. Miljanic<sup>5,10</sup>, L. Prepolec<sup>5</sup>, G. G. Rapisarda<sup>1</sup>, S. Romano<sup>1,2</sup>, M. L. Sergi<sup>1,2</sup>, N. Skukan<sup>5</sup>, N. Soic<sup>5</sup>, V. Tokic<sup>5</sup>, A. Tumino<sup>1,6</sup>, and M. Uroic<sup>5</sup>

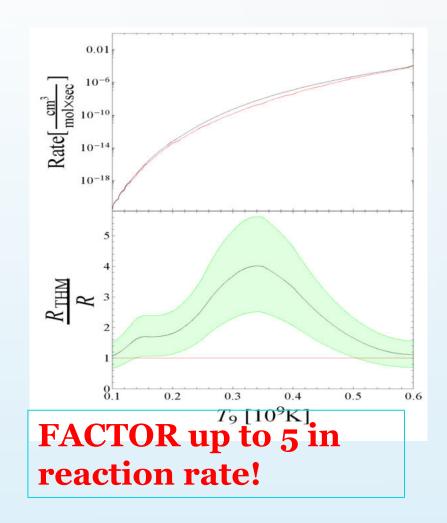
THE ASTROPHYSICAL JOURNAL, 836:57 (6pp), 2017 February 10 © 2017. The American Astronomical Society. All rights reserved.

https://doi.org/10.3847/1538-4357/836/1/57

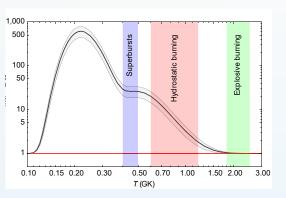


#### First Measurement of the ${}^{19}F(\alpha, p){}^{22}Ne$ Reaction at Energies of Astrophysical Relevance

R. G. Pizzone<sup>1</sup>, G. D'Agata<sup>1,2</sup>, M. La Cognata<sup>1</sup>, I. Indelicato<sup>1</sup>, C. Spitaleri<sup>1,2</sup>, S. Blagus<sup>3</sup>, S. Cherubini<sup>1,2</sup>, P. Figuera<sup>1</sup>, L. Grassi<sup>3</sup>,
G. L. Guardo<sup>1</sup>, M. Gulino<sup>1,4</sup>, S. Hayakawa<sup>1,5</sup>, R. Kshetri<sup>1,6</sup>, L. Lamia<sup>1,2</sup>, M. Lattuada<sup>1,2</sup>, T. Mijatović<sup>3</sup>, M. Milin<sup>7</sup>, D. Miljanić D.<sup>3,8</sup>,
L. Prepolec<sup>3</sup>, G. G. Rapisarda<sup>1</sup>, S. Romano<sup>1,2</sup>, M. L. Sergi<sup>1</sup>, N. Skukan<sup>3</sup>, N. Soić<sup>3</sup>, V. Tokić<sup>3</sup>, A. Tumino<sup>1,4</sup>, and M. Uroić<sup>3</sup>





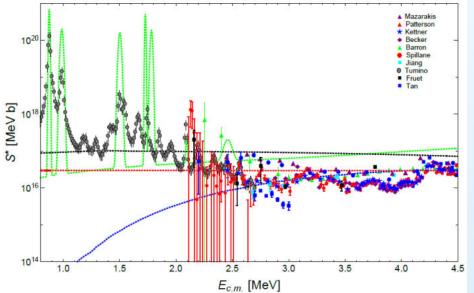


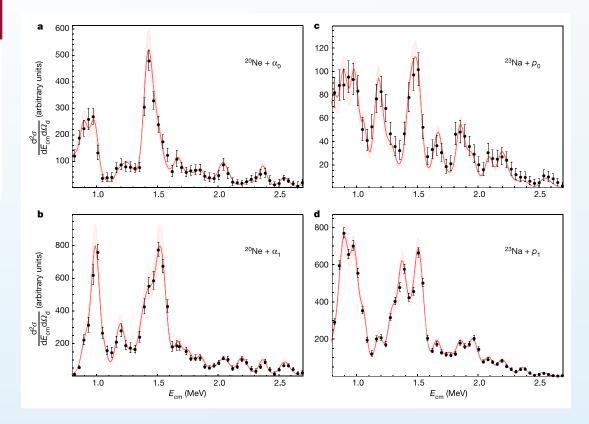
#### nature

Letter | Published: 23 May 2018 An increase in the  ${}^{12}C + {}^{12}C$  fusion rate from resonances at astrophysical energies

A. Tumino 🎒, C. Spitaleri, M. La Cognata, S. Cherubini, G. L. Guardo, M. Gulino, S. Hayakawa, I. Indelicato, L. Lamia, H. Petrascu, R. G. Pizzone, S. M. R. Puglia, G. G. Rapisarda, S. Romano, M. L. Sergi, R. Spartá & L. Trache

Nature 557, 687–690 (2018) Download Citation ±





Blue => hindrance effect Red => Caughland & Fowler 1988 Black => Godbey 2019

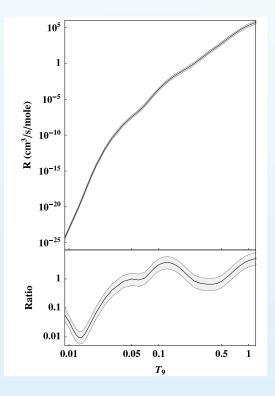
See A. Spiridon talk

### **Recent results of THM** *THM* successfully applied to RIBs

#### Experiment in CNS RIKEN and Texas A&M

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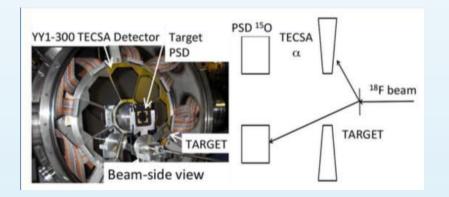


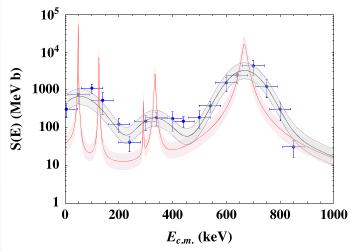


#### A Trojan Horse Approach to the Production of <sup>18</sup>F in Novae

M. La Cognata<sup>1</sup>, R. G. Pizzone<sup>1</sup>, J. José<sup>2,3</sup>, M. Hernanz<sup>3,4</sup>, S. Cherubini<sup>1,5</sup>, M. Gulino<sup>1,6</sup>, G. G. Rapisarda<sup>1,5</sup>, and C. Spitaleri<sup>1,5</sup>

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 <sup>6</sup> Facoltà di Ingegneria ed Architettura, Kore University, Viale delle Olimpiadi, 1, I-94100 Enna, Italy
 *Received 2017 May 10; revised 2017 July 31; accepted 2017 August 3; published 2017 August 31*





**Figure 2.** R-matrix analysis of the THM astrophysical factor (blue points) as in Figure 1. The evaluated uncertainty in the R-matrix fit is reported as a shadowed gray area and as a red band for the corresponding deconvoluted S(E)-factor.

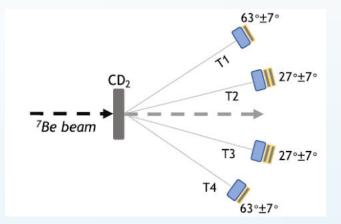
#### **Application of THM with RIBs and neutron induced reactions**

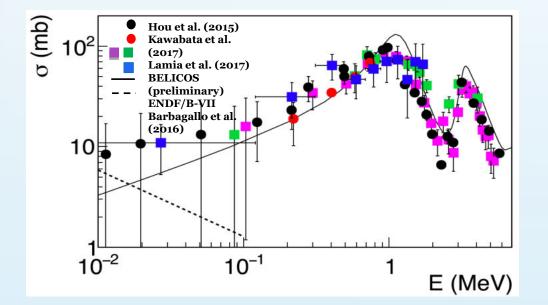
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CrossMark

On the Determination of the <sup>7</sup>Be(n,  $\alpha$ )<sup>4</sup>He Reaction Cross Section at BBN Energies

L. Lamia<sup>1,2</sup>, C. Spitaleri<sup>1,2</sup>, C. A. Bertulani<sup>3</sup>, S. Q. Hou<sup>3,4</sup>, M. La Cognata<sup>2</sup>, R. G. Pizzone<sup>2</sup>, S. Romano<sup>1,2</sup>, M. L. Sergi<sup>2</sup>, and A. Tumino<sup>2,5</sup> <sup>1</sup> Dipartimento di Fisica e Astronomia, Università degli Studi di Catania, Catania, Italy <sup>2</sup> NFN—Laboratori Nazionali del Sud, Catania, Italy <sup>3</sup> Department of Physics and Astronomy, Texas A&M University-Commerce, Commerce, TX 75428, USA <sup>4</sup> Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, China <sup>5</sup> Facoltà di Ingegneria e Architettura, Università degli Studi di Enna "Kore", Enna, Italy *Received 2017 September 12; revised 2017 October 20; accepted 2017 October 24; published 2017 November 30* 





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#### Cross-section Measurement of the Cosmologically Relevant ${}^{7}Be(n, \alpha)^{4}He$ Reaction over a Broad Energy Range in a Single Experiment

L. Lamia<sup>1,2</sup>, M. Mazzocco<sup>3,4</sup>, R. G. Pizzone<sup>2</sup>, S. Hayakawa<sup>5</sup>, M. La Cognata<sup>2</sup>, C. Spitaleri<sup>1,2</sup>, C. A. Bertulani<sup>6</sup>, A. Boiano<sup>7</sup>, C. Boiano<sup>8</sup>, C. Broggini<sup>4</sup>, A. Caciolli<sup>3,4</sup>, S. Cherubini<sup>1,2</sup>, G. D'Agata<sup>1,2,1,3</sup>, H. da Silva<sup>9</sup>, R. Depalo<sup>3,4</sup>, F. Galtarossa<sup>10</sup>, G. L. Guardo<sup>1,2</sup>, M. Gulino<sup>2,11</sup>, I. Indelicato<sup>1,2</sup>, M. La Commara<sup>7,12</sup>, G. La Rana<sup>7,12</sup>, R. Menegazzo<sup>4</sup>, J. Mrazek<sup>13</sup>, A. Pakou<sup>14</sup>, C. Parascandolo<sup>7</sup>, D. Piatti<sup>3,4</sup>, D. Pierroutsakou<sup>7</sup>, S. M. R. Puglia<sup>2</sup>, S. Romano<sup>1,2</sup>, G. G. Rapisarda<sup>2</sup>, A. M. Sánchez-Benítez<sup>15</sup>, M. L. Sergi<sup>2</sup>, O. Sgouros<sup>2,14</sup>, F. Soramel<sup>3,4</sup>, V. Soukeras<sup>2,14</sup>, R. Spartá<sup>1,2</sup>, E. Strano<sup>3,4</sup>, D. Torresi<sup>2</sup>, A. Tumino<sup>2,11</sup>, H. Yamaguchi<sup>5</sup>, and G. L. Zhang<sup>16</sup>



### Advantages of THM

A - It is possible measure the bare nucleus cross section  $s_b$  (or the bare nucleus Astrophysical Factor  $S_b(E)$ ) at Gamow energy for reactions involving charged particles and neutron.

B - <u>No extrapolation</u>

C - It is possible to measure excitation function in a "relatively" short time because typical order of magnitude for a three-body cross-section is mb;

D - One of the few ways to measure the electron screening effect: comparison with direct data;

E - Application to the radioactive beam measurements.

### Main limitation of THM

- A Preliminary study of quasi-free mechanism and tests of validity are necessary.
- B Presence of different 3-body reaction mechanisms

(Sequential Decay – Quasi-Free)

C - Absolute cross section isnot easily measurable

D - The excitation functions at energies around Coulomb barrier must be known from direct measurements;

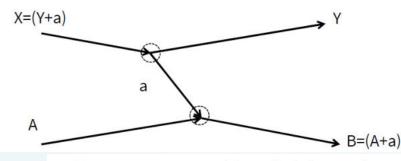
- D Measurements with high angular and energy resolutions are needed;
- E Theoretical analysis is needed: PWIA, MPWBA, DWBA...

TH Method is complementary to direct measurements as well as other indirect methods.



### **Asymptotic Normalitation Coefficient**

Studies performed by means of «simple» transfer reactions



In Distorted Wave Born Approximation, the transition amplitude between the states before and after the reactions can be written as:

$$M(E_i, \vartheta_{c.m.}) = \sum_{M_a} \left\langle \chi_f^{(-)} I_{Aa}^B \left| \Delta V \right| I_{Ya}^X \chi_i^{(+)} \right\rangle$$

Using DWBA we were able to find the ANC's coefficients from the spettroscopic factors. This gives us some advantages:

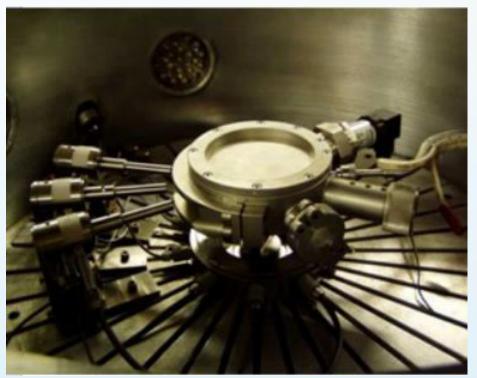
- For perihperal reactions, ANCs have small dependance from the potential
- $R_{l_B, j_B, l_x, j_x}$  is nearly indipendent from  $b^2$
- ANCs are defined in the nuclear «exterior», so are «observable»

$$\frac{d\sigma}{d\Omega} = \sum_{j_B, j_x} (C^B_{Aa, l_B, j_B})^2 (C^X_{Ya, l_x, j_x})^2 \frac{\sigma^{DWBA}_{l_B, j_B, l_x, j_x}}{b^2_{Aa, l_B, j_b} b^2_{Ya, j_x, j_x}} = \sum_{j_B, j_x} (C^B_{Aa, l_B, j_B})^2 (C^X_{Ya, l_x, j_x})^2 R_{l_B, j_B, l_x, j_x}$$



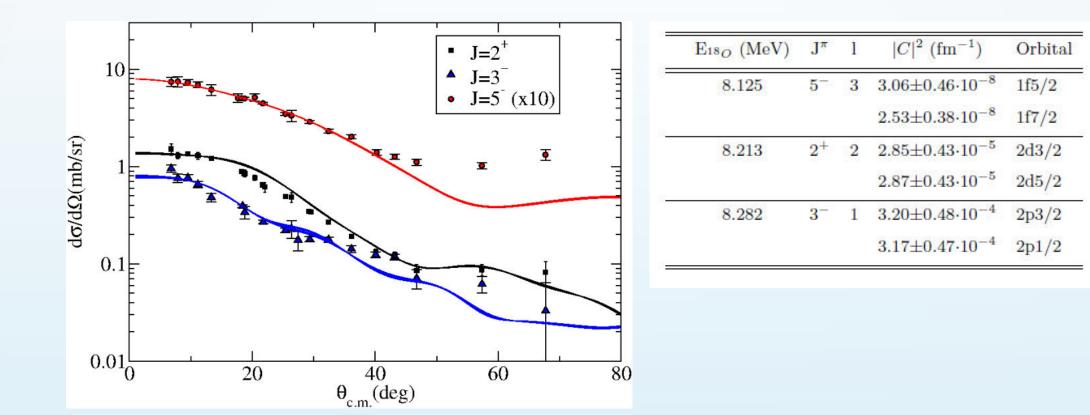
# <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study • Experimental setup

- ✓ Experiments performed at NPI-CAS at Rez, Prague (Czech Republic);
- ✓  $E_{\text{beam}}(^{2}\text{H})= 16.3 \text{ MeV};$
- ✓ Gas target <sup>17</sup>O, 90% pure;
- ✓ 8 point-like silicon telescope;
- ✓ Angular coverage: 6-67 degree in lab system



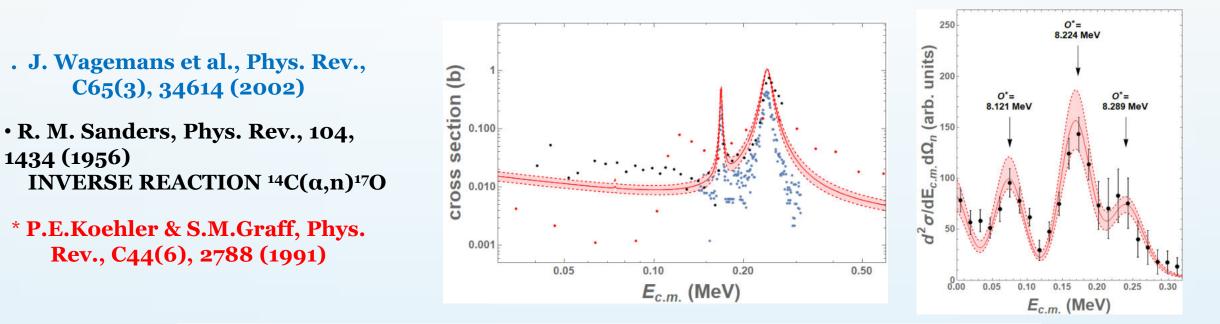


## <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study ANC extraction from <sup>17</sup>O(d,p)<sup>18</sup>O transfer reaction





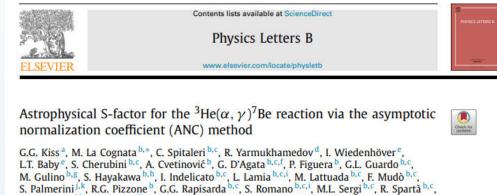
# <sup>17</sup>O(n,α)<sup>14</sup>C reaction: a case study Direct and THM data comparison



$E_{cm}$ (keV)	Koehler & Graff (1991)	Wagemans et al. (2002)	Guardo et al. (2017)	This Work
75	-	-	$36\pm5~{\rm eV}$	$33\pm5~{\rm eV}$
178	$1280{\pm}1000~{\rm eV}$	$2258{\pm}235~{\rm eV}$	$2260{\pm}300~{\rm eV}$	$2150{\pm}323~{\rm eV}$
244	$8000 \pm 1000 \text{ eV}$	$14739 \pm 590 \text{ eV}$	$14700 \pm 3800 \text{ eV}$	$16670 \pm 2500 \text{ eV}$



### **Recent results of ANC**

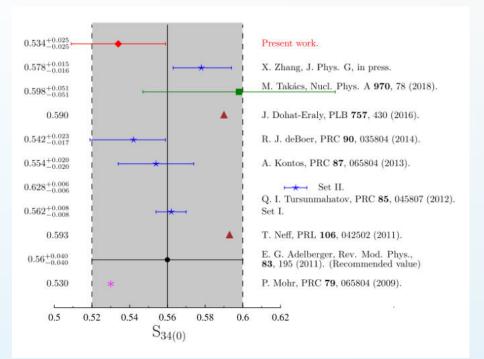


S. Palmerini<sup>J,k</sup>, R.G. Pizzone<sup>b</sup>, G.G. Rapisarda<sup>b,C</sup>, S. Romano<sup>b,C,i</sup>, M.L. Sergi<sup>b,c</sup>, R. Spartà<sup>b,c</sup> O. Trippella<sup>J,k</sup>, A. Tumino<sup>b,g</sup>, M. Anastasiou<sup>e</sup>, S.A. Kuvin<sup>e</sup>, N. Rijal<sup>e</sup>, B. Schmidt<sup>e</sup>, S.B. Igamov<sup>d</sup>, S.B. Sakuta<sup>1</sup>, K.I. Tursunmakhatov<sup>d,m</sup>, Zs. Fülöp<sup>a</sup>, Gy. Gyürky<sup>a</sup>, T. Szücs<sup>a</sup>, Z. Halász<sup>a</sup>, E. Somorjai<sup>a</sup>, Z. Hons<sup>f</sup>, J. Mrázek<sup>f</sup>, R.E. Tribble<sup>n</sup>, A.M. Mukhamedzhanov<sup>n</sup>

#### PHYSICAL REVIEW C 104, 015807 (2021)

#### Indirect determination of the astrophysical *S* factor for the ${}^{6}\text{Li}(p, \gamma) {}^{7}\text{Be}$ reaction using the asymptotic normalization coefficient method

G. G. Kiss,<sup>1,\*</sup> M. La Cognata<sup>(0)</sup>,<sup>2,†</sup> R. Yarmukhamedov,<sup>3,‡</sup> K. I. Tursunmakhatov,<sup>3,4</sup> I. Wiedenhöver,<sup>5</sup> L. T. Baby,<sup>5</sup>
S. Cherubini,<sup>2,6</sup> A. Cvetinović,<sup>2</sup> G. D'Agata,<sup>7</sup> P. Figuera,<sup>2</sup> G. L. Guardo,<sup>2,6</sup> M. Gulino,<sup>2,8</sup> S. Hayakawa,<sup>9</sup> I. Indelicato,<sup>2,6</sup> L. Lamia,<sup>2,6,10</sup> M. Lattuada,<sup>2,6</sup> F. Mudò,<sup>2,6</sup> S. Palmerini,<sup>11,12</sup> R. G. Pizzone,<sup>2</sup> G. G. Rapisarda,<sup>2,6</sup> S. Romano,<sup>2,6,10</sup> M. L. Sergi,<sup>2,6</sup> R. Spartà,<sup>2,6</sup> C. Spitaleri,<sup>2,6</sup> O. Trippella,<sup>11,12</sup> A. Tumino,<sup>2,8</sup> M. Anastasiou,<sup>5</sup> S. A. Kuvin,<sup>5</sup> N. Rijal,<sup>5</sup> B. Schmidt,<sup>5</sup> S. B. Igamov,<sup>3</sup> S. B. Sakuta,<sup>13</sup> Zs. Fülöp,<sup>1</sup> Gy. Gyürky,<sup>1</sup> T. Szücs,<sup>1</sup> Z. Halász,<sup>1</sup> E. Somorjai,<sup>1</sup> Z. Hons,<sup>7</sup> J. Mrázek,<sup>7</sup> R. E. Tribble,<sup>14</sup> and A. M. Mukhamedzhanov<sup>14</sup>



Astrophysical factor at Gamow energies for the Sun was extracted via the ANC method. For the  ${}^{3}\text{He}({}^{4}\text{He},\gamma){}^{7}\text{Be}$  case, in good agreement with previous experiments



### **Recent results of ANC**

PHYSICAL REVIEW C 103, 015806 (2021)

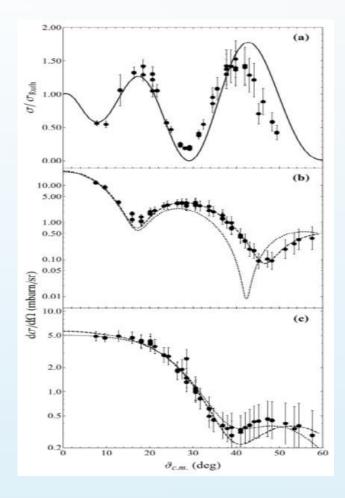
#### $^{26}$ Si $(p, \gamma)$ $^{27}$ P direct proton capture by means of the asymptotic normalization coefficients method for mirror nuclei

G. D'Agata<sup>(1,\*)</sup> A. I. Kilic<sup>(1,\*)</sup> V. Burjan<sup>(1,\*)</sup> J. Mrazek<sup>(3,\*)</sup> V. Glagolev,<sup>1</sup> V. Kroha,<sup>1</sup> G. L. Guardo,<sup>2</sup> M. La Cognata,<sup>2</sup> L. Lamia,<sup>2,3,4</sup> S. Palmerini,<sup>5,6</sup> R. G. Pizzone,<sup>2</sup> G. G. Rapisarda,<sup>2</sup> S. Romano,<sup>2,3,4</sup> M. L. Sergi,<sup>2,3</sup> R. Spartà<sup>(3)</sup>,<sup>2,3</sup> C. Spitaleri,<sup>2</sup> I. Siváček,<sup>1,7</sup> and A. Tumino<sup>2,8</sup>

TABLE IV. The <sup>26</sup>Si( $p, \gamma$ ) <sup>27</sup>P reaction rate (in  $\frac{cm^4}{mot ssic}$ ) for the ground-state direct capture and the first excited state resonant contribution, extracted using Eqs. (12)–(14). The present data are compared with the ones coming from [26], derived in the same way but using the different  $\Gamma_{\gamma}$ ,  $\Gamma_{p}$ , and S(0).

	Reaction rate ground state (this work)				Reaction rate first excited state (this work)			
T <sub>9</sub>	Lower Limit	Value	Upper Limit	Ref. [26]	Lower Limit	Value	Upper Limit	Ref. [26]
0.1	$8.48 \times 10^{-14}$	$9.57 \times 10^{-14}$	$1.25 \times 10^{-13}$	$6.68 \times 10^{-14}$	$2.09 \times 10^{-12}$	$2.91 \times 10^{-12}$	$3.84 \times 10^{-12}$	$1.41 \times 10^{-12}$
0.2	$2.71 \times 10^{-09}$	$3.06 \times 10^{-09}$	$4.01 \times 10^{-09}$	$2.13 \times 10^{-09}$	$7.63 \times 10^{-05}$	$1.06 \times 10^{-04}$	$1.40 \times 10^{-04}$	5.15×10-05
0.3	$4.02 \times 10^{-07}$	$4.53 \times 10^{-07}$	$5.94 \times 10^{-07}$	$3.17 \times 10^{-07}$	$1.94 \times 10^{-02}$	$2.71 \times 10^{-02}$	$3.58 \times 10^{-02}$	1.32×10-02
0.4	$9.28 \times 10^{-06}$	$1.05 \times 10^{-05}$	$1.37 \times 10^{-05}$	$7.31 \times 10^{-06}$	$2.74 \times 10^{-01}$	$3.81 \times 10^{-01}$	$5.03 \times 10^{-01}$	$1.85 \times 10^{-01}$
0.5	$8.57 \times 10^{-05}$	$9.66 \times 10^{-05}$	$1.27 \times 10^{-04}$	$6.75 \times 10^{-05}$	1.24	1.72	2.28	8.38×10-0
0.6	$4.64 \times 10^{-04}$	$5.24 \times 10^{-04}$	$6.86 \times 10^{-04}$	$3.66 \times 10^{-04}$	3.23	4.49	5.93	2.18
0.7	$1.78 \times 10^{-03}$	$2.01 \times 10^{-03}$	$2.63 \times 10^{-03}$	$1.40 \times 10^{-04}$	6.17	8.58	$1.13 \times 10^{+01}$	4.17
0.8	$5.39 \times 10^{-03}$	$6.08 \times 10^{-03}$	$7.96 \times 10^{-03}$	$4.25 \times 10^{-03}$	9.76	$1.36 \times 10^{+01}$	$1.79 \times 10^{+01}$	6.60
0.9	$1.37 \times 10^{-02}$	$1.54 \times 10^{-02}$	$2.02 \times 10^{-02}$	$1.08 \times 10^{-02}$	$1.37 \times 10^{+01}$	$1.90 \times 10^{+01}$	$2.51 \times 10^{+01}$	9.23
1.0	$3.05 \times 10^{-02}$	$3.44 \times 10^{-02}$	$4.51 \times 10^{-02}$	$2.40 \times 10^{-02}$	$1.76 \times 10^{+01}$	$2.44 \times 10^{+01}$	$3.23 \times 10^{+01}$	1.19×10+0
1.2	$1.14 \times 10^{-01}$	$1.28 \times 10^{-01}$	$1.68 \times 10^{-02}$	$8.97 \times 10^{-02}$	$2.47 \times 10^{+01}$	$3.44 \times 10^{+01}$	$4.54 \times 10^{+01}$	$1.67 \times 10^{+01}$
1.4	$3.24 \times 10^{-01}$	$3.66 \times 10^{-01}$	$4.79 \times 10^{-01}$	$2.56 \times 10^{-01}$	$3.04 \times 10^{+01}$	$4.23 \times 10^{+01}$	$5.59 \times 10^{+01}$	$2.06 \times 10^{+01}$
1.6	$7.67 \times 10^{-01}$	$8.65 \times 10^{-01}$	1.13	$6.04 \times 10^{-01}$	$3.47 \times 10^{+01}$	$4.82 \times 10^{+01}$	$6.36 \times 10^{+01}$	$2.34 \times 10^{+0}$
1.8	1.58	1.78	2.34	1.25	$3.75 \times 10^{+01}$	$5.22 \times 10^{+01}$	$6.89 \times 10^{+01}$	2.53×10+0
2.0	2.95	3.32	4.36	2.32	$3.93 \times 10^{+01}$	$5.47 \times 10^{+01}$	$7.22 \times 10^{+01}$	2.66×10+0

#### REACTION RATE CALCULATED AT GAMOW ENERGIES!!!





### Laser-induced Experiments

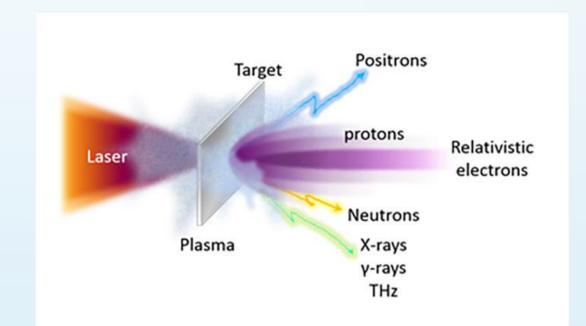
An effective approach to improve our knowledge on fusion reactions in stars can be to mimic the stellar conditions by generating a controlled laboratory plasma with thermodynamical status not too different from the stellar conditions.



### Laser-induced Experiments

An effective approach to improve our knowledge on fusion reactions in stars can be to mimic the stellar conditions by generating a controlled laboratory plasma with thermodynamical status not too different from the stellar conditions.

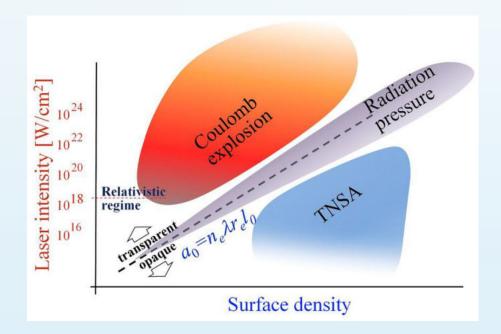
The advent of novel techniques in laser amplification on one hand opened the way to the the implementation of high-power highintensity lasers in many facilities around the world and on the other it forced a steep development on the diagnostics side, de facto unlocking the feasibility of **new paradigms of research**.





### **Nuclear Astrophysics with Laser Beams**

The most common scenario on laser-ion acceleration relies on focusing a high intensity laser pulse into a target made of the species of interest. Depending on the intensity and on the surface density of the target, different acceleration mechanisms are involved. The most studied scenario today relies on the Target Normal Sheath Acceleration mechanism originated by high-intensity laser pulses focused on thin targets.



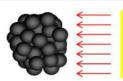
These methods will open the way for a new approach to study Nuclear Astrophysics Reaction such as:

- deuterium- deuterium
- deuterium-<sup>3</sup>He
- proton-lithium
- proton-boron
- $12C^{-12}C$
- 16O-16O
- and much more....

## **Coulomb Explosion**

#### THE COULOMB EXPLOSION PARADIGMA

The interaction of ultra-short laser pulses with an expanding gas mixture at controlled temperature and pressure inside a vacuum chamber causes the formation of plasmas with multi-keV temperature. These energies overlap with the typical temperatures of stellar environments where thermonuclear reactions occur, thus making this paradigm a *perfect scenario for nuclear astrophysics research*.



#### Step 1

Clusters are irradiated by high intensity laser pulse (~10<sup>16</sup>~10<sup>18</sup> W/cm<sup>2</sup>).

#### Step 2

Laser pulse energy is first absorbed by electrons via heating mechanisms such as rapid collisional heating.

Step 3 Electrons escape from the cluster and leave positive charge build-up on the cluster.

Step 4

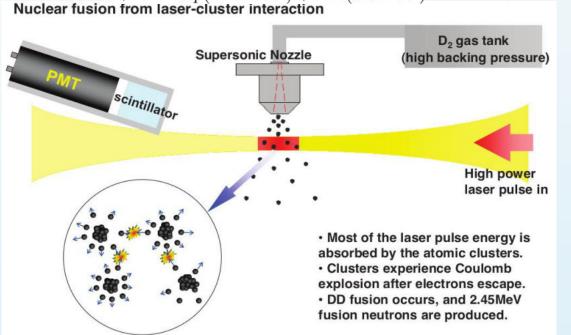
The cluster "explodes" and deuterons acquire multi-keV kinetic energy.

#### Example: deuterium-deuterium fusion

 $d + d \rightarrow {}^{3}\text{He}(0.82MeV) + n(2.45MeV)$ 

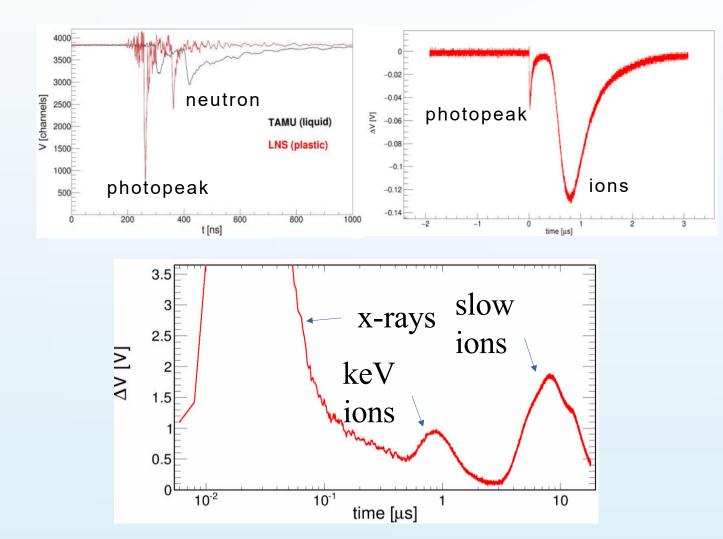
 $d+d \rightarrow p(3.02 MeV) + t(1.01 MeV)$ 

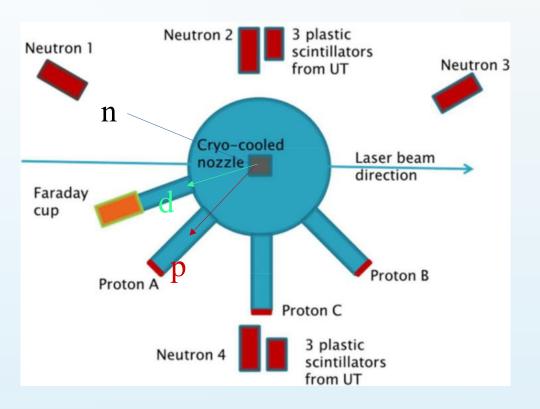
 $d + {}^{3}\text{H}e \rightarrow p(14.7MeV) + {}^{4}\text{H}e(3.6MeV)$ 



Kinetic Energy < 10<sup>2</sup> keV Density ~ 10<sup>18</sup> atoms/cm<sup>3</sup> 10<sup>5</sup>-10<sup>7</sup> neutrons per shot

### **Experimental approach**





### **CE Results**

 PHYSICAL REVIEW C

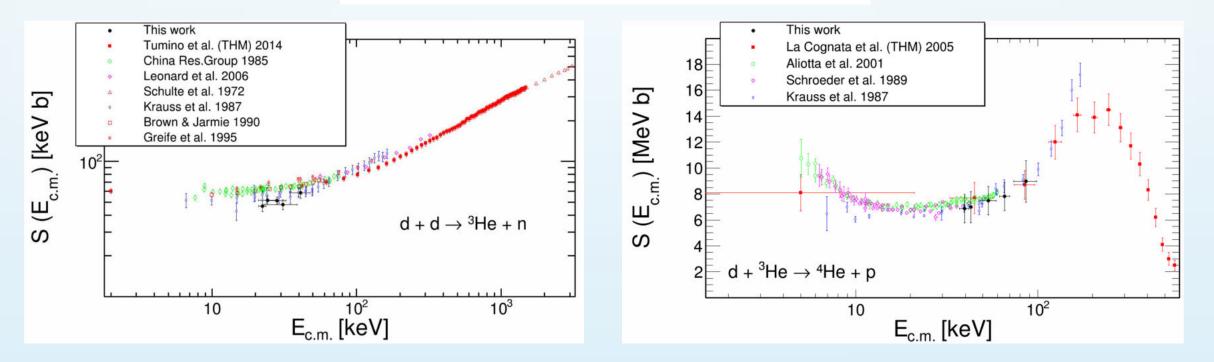
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 Model-independent determination of the astrophysical S
 factor in laser-induced fusion plasmas

D. Lattuada, M. Barbarino, A. Bonasera, W. Bang, H. J. Quevedo, M. Warren, F. Consoli, R. De Angelis, P. Andreoli, S. Kimura, G. Dyer, A. C. Bernstein, K. Hagel, M. Barbui, K. Schmidt, E. Gaul, M. E. Donovan, J. B. Natowitz, and T. Ditmire Phys. Rev. C **93**, 045808 – Published 19 April 2016



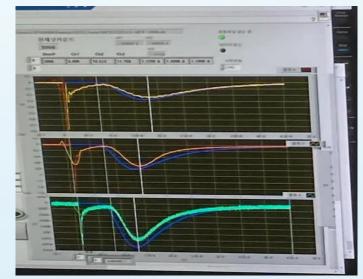


### **GIST Experiment**

Over 3000 shots, with T control between kT~1 keV and kT~40 keV





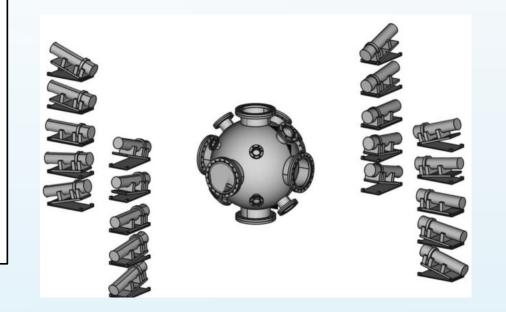




### Versatile Array for Laser-induced Astrophysics Research

Science-driven, portable, cost-efficient

.cryo-cooled supersonic nozzle
.compact interaction chamber
.neutron ToF detectors (plastic/liquid scintillators)
.charged particle ToF detectors (SiC/CVD
diamond detectors + FCs)
.2 TPS
.(CR39 for checks/normalization)





### The AsFiN Collaboration







#### THE 12<sup>TH</sup> EUROPEAN SUMMER SCHOOL ON EXPERIMENTAL NUCLEAR ASTROPHYSICS

Primordial Nucleosynthesis and Early Stars, Stellar Evolution, Hydrostatic and Explosive Nucleosynthesis, Plasmas in Stars and Laboratories, Detectors and Facilities for Nuclear Astrophysics, Experiments with rare and radioactive isotopes, Indirect Methods

https://agenda.infn.it/e/essena2024

#### 16-22 JUNE 2024 ACI TREZZA (CT)