

Constraining the Astrophysical γ Process: Cross Section Measurement of the ⁸²Kr(p,γ)⁸³Rb Reaction in Inverse Kinematics

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Introduction





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U.S. Department of Energy Office of Science National Science Foundation Michigan State University W. Rapp, J. Görres, M. Wiescher, H. Schatz, and F. Käppeler. Astrophys J, 653:474, 2006.

A. Tsantiri, Russbach School on Nuclear Astrophysics, 14/3/2023, Slide 4

γ-summing technique with SuN



- <u>Summing Nal(TI) SuN</u>: Large size, high efficiency γ-ray detector
- 8 optically isolated segments
- 24 PMTs
- Sum of Segments (SoS) → Information about individual γ-rays
- Total Absorption Spectrum (TAS) \rightarrow Information about total excitation energy E_x





Experiment at ReA NSCL



Experimental Setup (without SuNSCREEN)

E. Klopfer et al, Nucl. Instrum. Meth. Phys. Res. A 788, 5 (2015)



Background Subtraction

Background contributions:

- Cosmic rays → SuNSCREEN veto
- Room Background \rightarrow Pulsed Beam
- Interaction of the ⁸²Kr beam with the beam line and the gas cell \rightarrow gas cell full and empty runs 1000 Full gas cell





Analysis Overview





Efficiency and Yield Determination with RAINIER and GEANT4

- RAINIER: Simulates the de-excitation of a compound nucleus through $\boldsymbol{\gamma}$ cascades
- To describe our compound nucleus we use combinations of nuclear level densities $\rho(E_x E_y)$ and γ ray strength functions $\gamma SF(E_y)$ that can replicate our SoS





Efficiency and Yield Determination with RAINIER and GEANT4

- GEANT4: account for detector's geometry
- Fit to experimental TAS, SoS and multiplicity using χ^2 minimization code





Results

- Standard statistical model calculations tend to overproduce the cross section
- Based on experimental data in neighboring nuclei, theory appears to consistently overestimates reaction rates in this mass region





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But we managed to constrain the product of the NLD and γSF, and therefore we should be able to accurately reproduce our extracted cross section with TALYS!



Results

- Standard statistical model calculations (default TALYS) tend to overproduce the cross section
- A better description of the experimental data can be obtained with the suggested combinations of NLD and γSF (Optimized TALYS)
- The cross section at lower energies is dependent on the choice of OMP and width fluctuation correction





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Summary and Conclusion

- Systematic study of (p,γ) reactions allows for constrains on theoretical models used in astrophysical applications
- The ⁸²Kr(p,γ)⁸³Rb reaction cross section was measured for the first time in inverse kinematics
- A better description of the experimental data can be obtained with the suggested combinations of NLD and γSF
- The cross section at lower energies is dependent on the choice of OMP and width fluctuation correction
- A. Tsantiri, A. Palmisano-Kyle, A. Spyrou, P. Mohr, et al. [accepted, in press] Cross Section Measurement of the ⁸²Kr(p,γ)⁸³Rb Reaction in Inverse Kinematics, Phys. Rev. C.
- Next step: ⁷³As(p,γ)⁷⁴Se FRIB experiment e22505 scheduled for the summer [PI: A. Palmisano-Kyle]

Thank you for your attention!



Backup: Theoretical Investigation





Backup: What about OMP?





Backup: A little bit of Hauser-Feshbach

$$\langle \sigma_{12}(E) \rangle = \frac{\lambda^2}{2\pi} \sum_{J,\pi} W_{12} \omega_J \frac{T_{1,J,\pi} T_{2,J,\pi}}{T_{J,\pi}}$$

- λ : de Broglie wavelength
- ω : function of angular momenta J_1 , J_2 and J_r
- T : transmission coefficients (function of energy width Γ and level density ρ)
- W₁₂: width fluctuation correction (set to 1 through independence hypothesis)

$$T_{\gamma}(E, J, \pi) = \sum_{0}^{\nu_{r}} T_{\gamma}^{\nu}(E, J, \pi, E_{r}^{\nu}, J_{r}^{\nu}, \pi_{r}^{\nu}) + \int_{E_{r}^{\nu_{r}}}^{E} \sum_{J_{r}, \pi_{r}} T_{\gamma}^{\nu}(E, J, \pi, E_{r}^{\nu}, J_{r}^{\nu}, \pi_{r}^{\nu}) \cdot \rho(E_{r}, J_{r}, \pi_{r}) dE_{r}$$

Experimentally known discrete levels
$$T_{\gamma}(E_{\gamma}, J, \Pi) = 2\pi \underbrace{f(E_{\gamma})}_{\gamma \text{SF}} E_{\gamma}^{2\lambda+1}$$



Backup: NLD models

Constant Temperature Model (Gilbert, Cameron and Ericson) Parameters: T, E₀

$$\rho_{\rm CT}(E_x) = \frac{1}{T} \exp\left(\frac{E_x - E_0}{T}\right)$$

Bach Shifted Fermi Gas Model (Bethe) Parameters: α , E_1

$$\rho_{\rm BSFG}(E_x) = \frac{1}{12\sqrt{2}\alpha^{1/4}} \frac{\exp\left[2\sqrt{\alpha}E_x - E_1\right]}{(E_x - E_1)^{5/4}}$$



Backup: NLD and **ySF** parameters chosen

| LD Model | LD Model Details | | Upbend in γSF | |
|--------------|------------------|----------------------------|---------------|----------------|
| CT default | T = 0.824 | E ₀ = -1.16 [1] | No | |
| BSFG default | α = 10.17 | E ₁ = -0.54 [1] | No | |
| СТ | T = 0.824 | E ₀ = -2.2 | No | |
| СТ | T = 0.861 | E ₀ = -3.34 [2] | No | |
| BSFG | α = 10.17 | E ₁ = -1.6 | No | |
| BSFG | α = 10.17 | $E_1 = -0.54$ | a = 1.5 | c = 8.7E-8 [3] |
| BSFG | α = 10.17 | $E_1 = -0.54$ | a = 1.0 | c = 1.0E-7 |

<u>**vSF chosen**</u>: Generalized Lorentzian of the form of Kopecky-Uhl [4] <u>**Upbend**</u> added of the form: $f_{upbend} = c \cdot \exp(-a \cdot E_{\gamma})$

[1] T. von Egidy and D. Bucurescu, Phys. Rev. C 80, 054310 (2009)
[2] R. Hoffman, F. Dietrich, R. Bauer, K. Kelley, and M. Mustafa 10.2172/15014588 (2004)
[3] M. Guttormsen, R. Chankova, U. Agvaanluvsan, and et. al., Phys. Rev. C 71, 044307 (2005).
[4] J. Kopecky and M. Uhl, Phys. Rev. C 41, 1941 (1990)

