The slow neutron capture process

Sergio Cristallo

INAF – Osservatorio Astronomico d'Abruzzo (Italy) INFN – Sezione di Perugia (Italy)

For the umpteenth time...

The Origin of the Solar System Elements



Graphic created by Jennifer Johnson http://www.astronomy.ohio-state.edu/~jaj/nucleo/ Astronomical Image Credits: ESA/NASA/AASNova





$\tau_{\beta} \gg \tau_{n} \quad \Leftrightarrow \quad N_{n} > 10^{20} \text{ n/cm}^{3}$

Unstable nucleus captures another neutron before decaying

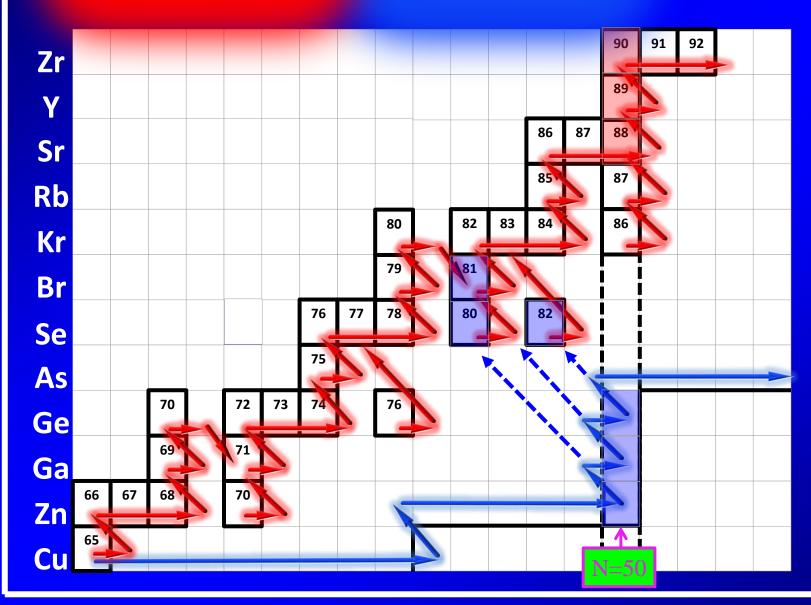
<u>The s-process</u>

 $\tau_{\beta} \leftrightarrow \tau_{n} \quad \Leftrightarrow \quad N_{n} \sim 10^{7} \text{ n/cm}^{3}$

Unstable nucleus <u>decays</u> before capturing another neutron

In principle one might expect to encounter astrophysical neutron fluxes in the large region between these two densities and have thereby intermediate processes between s and r. Such events are apparently not common, and it is one of the fortunate simplifications in the application theory of synthesis by neutron capture that the most common fluxes are either quite small or quite large... **if we ignore the i-process**.

Proton number (Z)

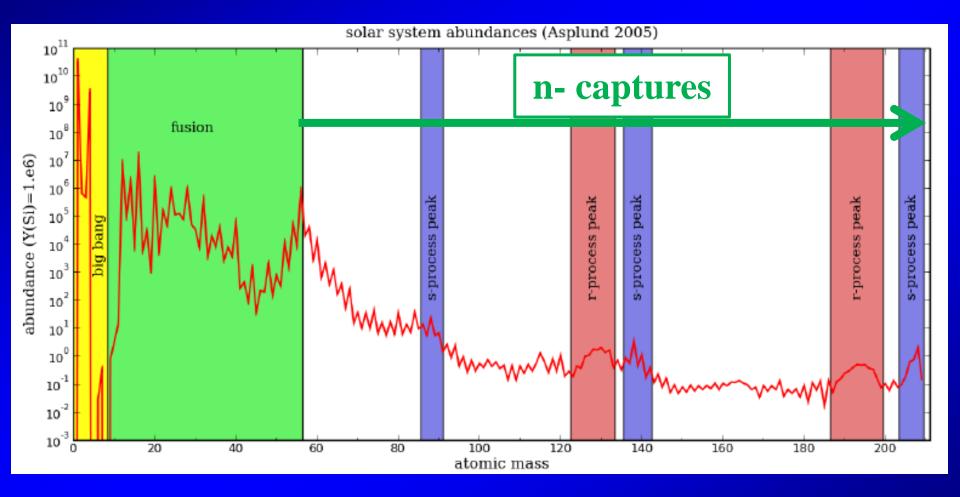


r process

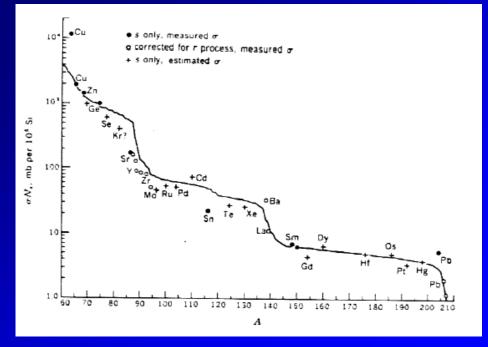
s process

Mass number (protons+neutrons) (A)⁴

Solar System Abundances



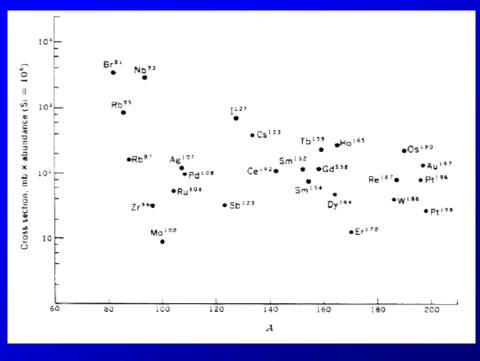
Cowan+2021



s-process

Easy to be reproduced with a series of neutron exposures (with an exponential distribution)

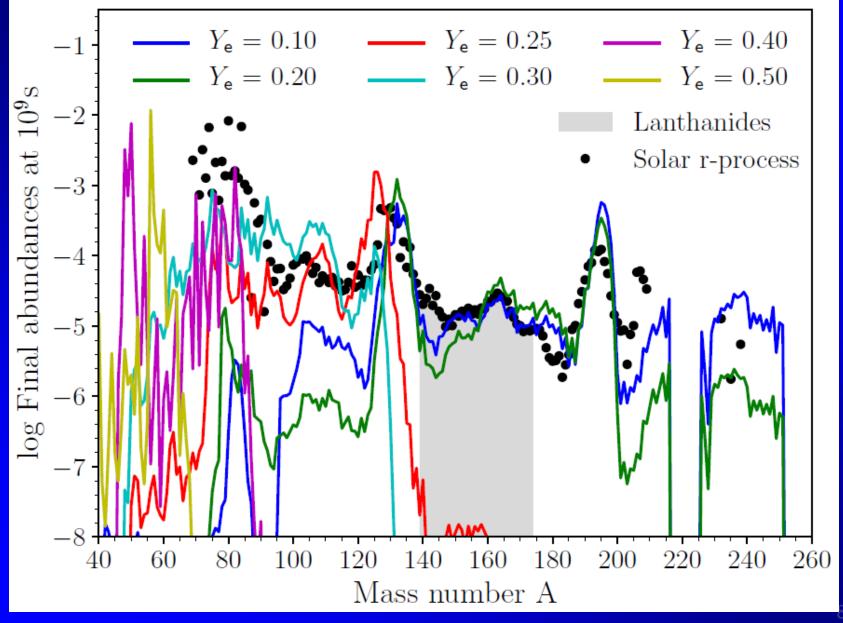
r-process Do you see any distribution?



How can we determine the r-process contribution to the solar distribution?

1 = -S

Typical r-process distributions



Solar r-process residuals from Prantzos+2020

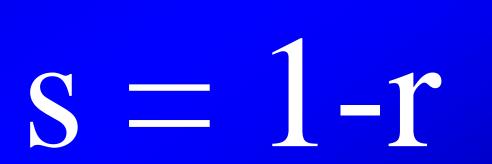
Various contributions to isotopic abundances

57 La 56 Ba 55 Cs 54 Xe 53 I 52 Te 51 Sb 50 Sn 49 In 48 Cd 47 Ag 46 Pd 45 Rh 45 Rh 42 Mo 41 Nb 40 Zr 39 Y	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	75 Re 74 W 73 Ta 72 Hf 71 Lu 70 Yb 69 Tm 68 Er 67 Ho 67 Ho	209 204 206 207 208 203 205 196 198 199 200 201 202 204 197 190 192 194 195 196 198 190 192 194 195 196 198 191 193 184 186 187 184 186 187 188 189 190 192 185 187 188 189 190 192 184 185 187 188 189 190 192 185 187 188 189 190 192 185 187 1 188 189 190 192 185 187 1 1 190 192 1 180 181 1
40 Zr 39 Y 38 Sr 37 Rb 36 Kr	90 91 92 94 96 89 84 86 87 88 85 87 88 85 87 78 80 82 83 84 86 S	66 Dy 65 Tb 64 Gd 63 Eu	156 158 160 161 162 163 164 159 152 154 155 156 157 158 160 151 153 155 156 157 158 160
35 Br 34 Se 33 As 32 Ge 31 Ga	79 81 5 74 76 77 78 80 82 r 75 70 72 73 74 76 p 69 71 Prantzos+2020 P	62 Sm 60 Nd 59 Pr 58 Ce	144 147 148 149 150 152 154 142 143 144 145 146 148 150 141 136 138 140 142



Arlandini+1999; Goriely+1999; Simmerer+2004; Bisterzo+2018; Prantzos+2020

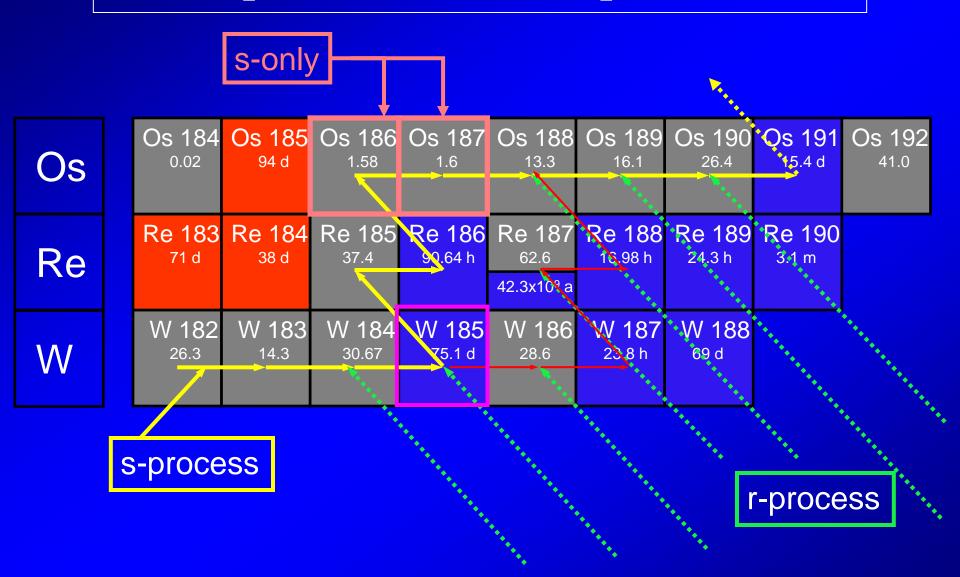
VS





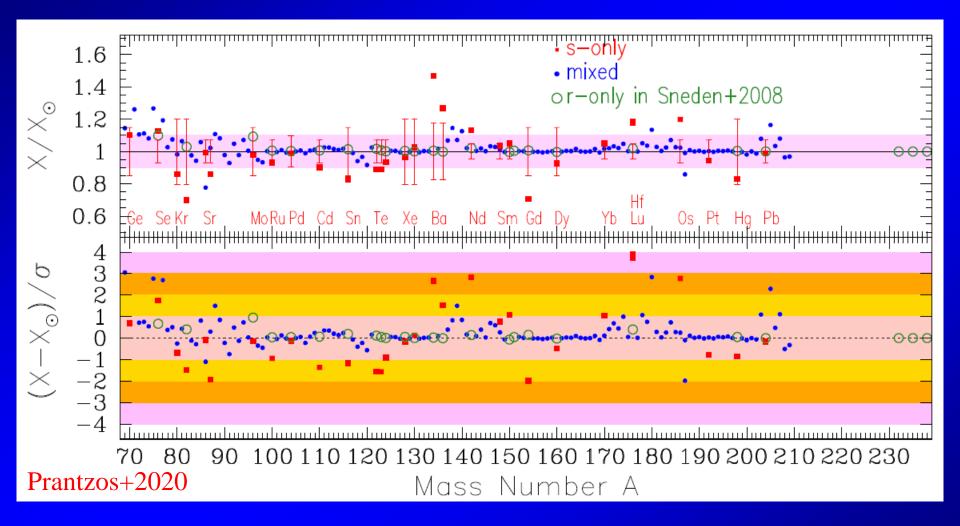
Busso,Kratz,Palmerini + 2022

How s-process neutron captures work?



Branching points: if $\tau_{\beta} \sim \tau_{\mu} \Rightarrow$ several paths are possible

The importance of s-only isotope distribution



They can be used to test stellar model & nucleosynthesis robustness.

Seeds for the s-process

Main seeds are ⁵⁶Fe nuclei... Why not the most abundant ¹H, ⁴He or ¹²C???

Seeds for the s-process

Main seeds are ⁵⁶Fe nuclei... Why not the most abundant ¹H, ⁴He or ¹²C???

10⁶ - 10⁹ v

Stable Unknown

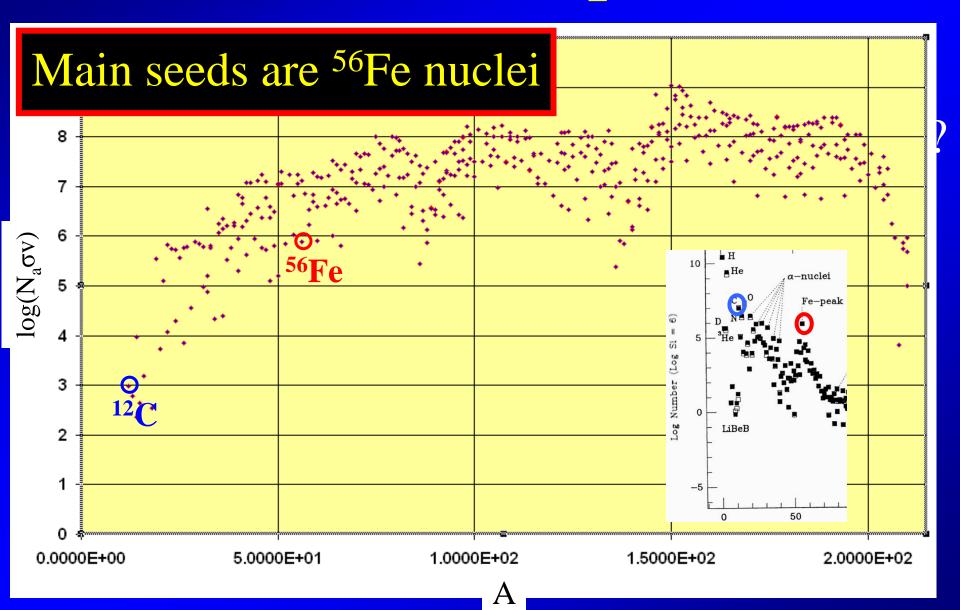
The reason lies in the nuclear structure of nuclei...and in the stars!!

RATE[H(n, γ)²H] \propto N(H) \downarrow 10⁻¹²



 $> 10^9 \text{ y}$

Seeds for the s-process



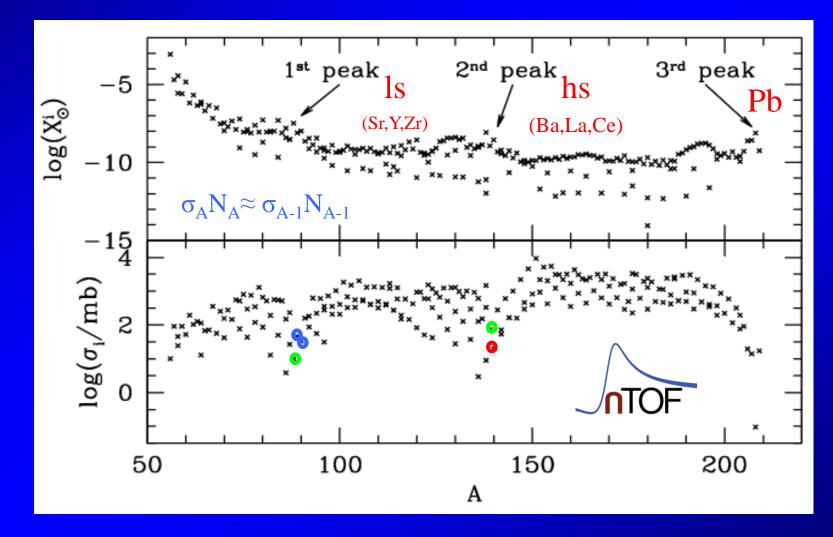
MAGIC NUCLEI



at neutron magic numbers

 \rightarrow

abundance curve for elements beyond iron

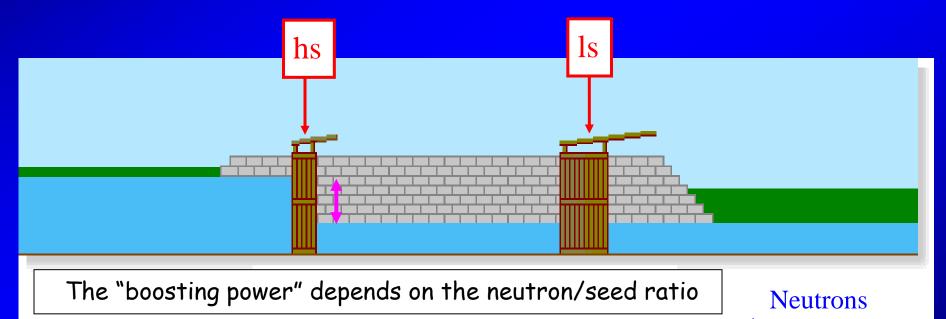


The three s-process peaks

 $1^{st} peak \rightarrow \underline{ls} elements (Sr,Y,Zr) [N=50]$

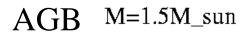
 $2^{nd} \text{ peak} \rightarrow \underline{\text{hs}} \text{ elements}$ (Ba,La,Ce,Nd,Sm) [N=82]

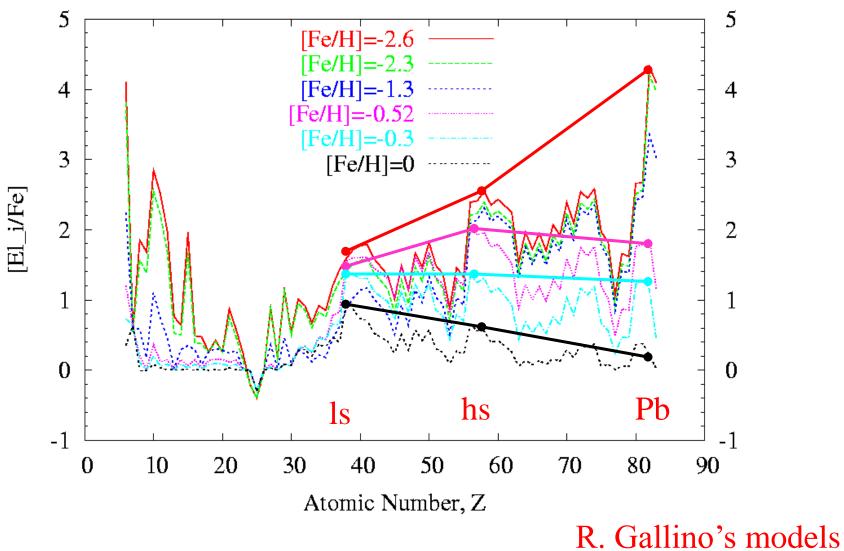
 $3^{rd} \text{ peak} \rightarrow \underline{\text{lead}} (^{208}\text{Pb}) [N=126 \& P=82]$



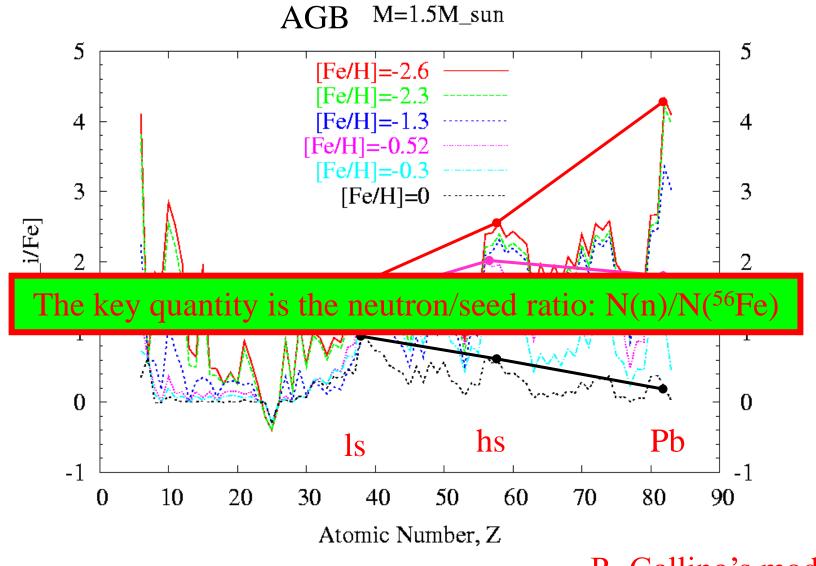
A sluice system with opening bulkheads

SURFACE DISTRIBUTION



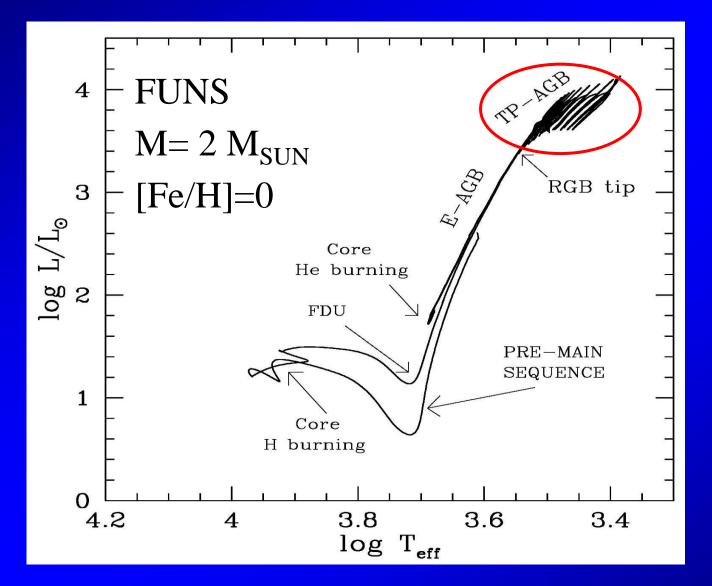


SURFACE DISTRIBUTION



R. Gallino's models

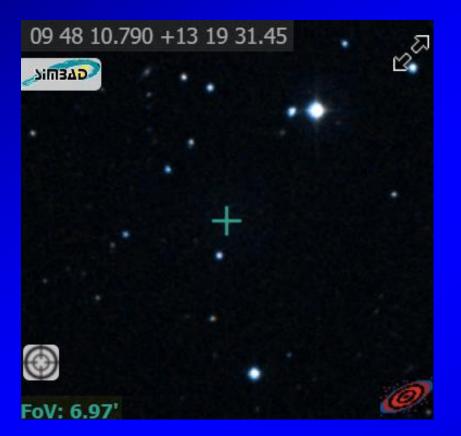
Asymptotic Giant Branch (AGB) stars



 $\tau_{\rm MS} \approx 1 \ {\rm Gyr}$ $\tau_{\rm AGB} \approx 1 \ {\rm Myr}$

Can we see an AGB star in the sky?

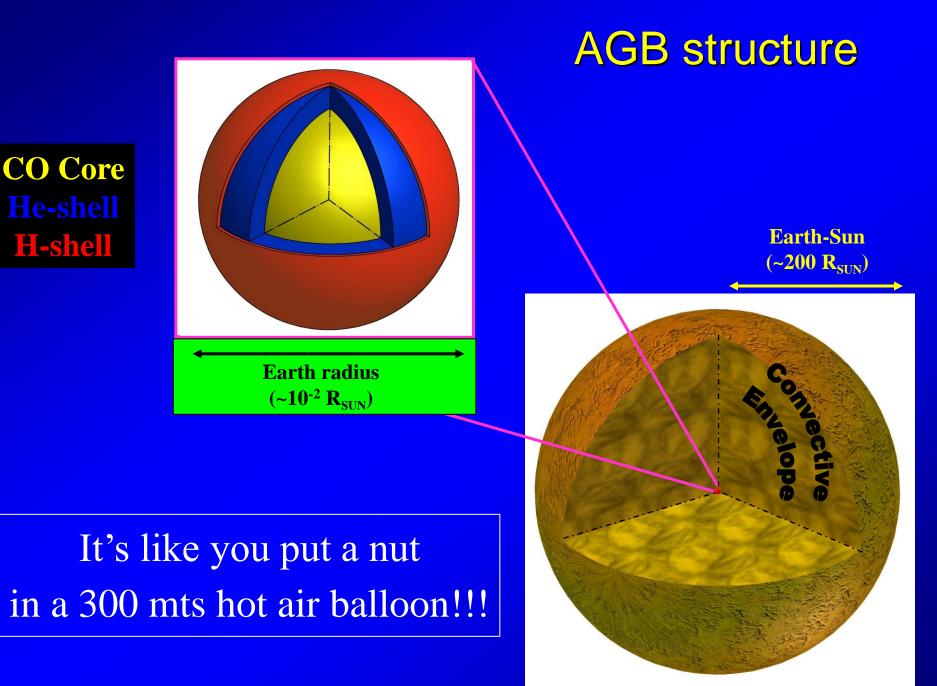
CW Leonis





OPTICAL BAND

INFRARED BAND



Where do s-process neutrons come from?

Free neutrons are NOT abundant in the major phases of nuclear burnings.

Neutrons are liberated to some extent by secondary reactions during helium burning in <u>Asymptotic Giant Branch (AGB) stars</u>, as well as during <u>core-He</u> and shell-C burnings of massive stars.

Major neutron sources of the s-process $^{13}C(\alpha,n)^{16}O$ $^{22}Ne(\alpha,n)^{25}Mg$

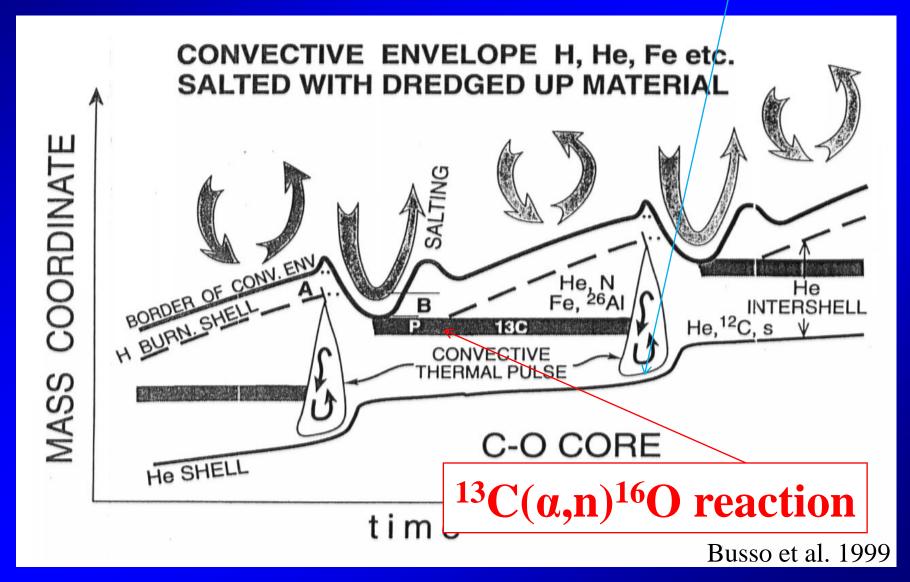
The nuclear paths

¹³C: main source for the Main component ${}^{12}C(p,\gamma){}^{13}N(\beta^+){}^{13}C$

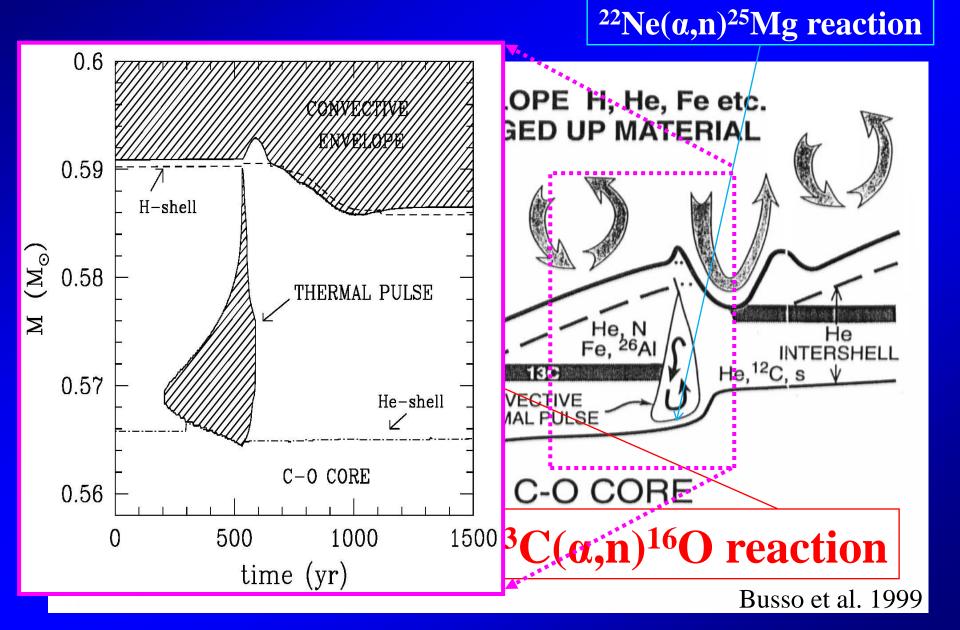
²²Ne: main source for the Weak component ¹⁴N(α,γ)¹⁸F(β^+)¹⁸O(α,γ)²²Ne

The s-process in AGB stars

$^{22}Ne(\alpha,n)^{25}Mg$ reaction



The s-process in AGB stars

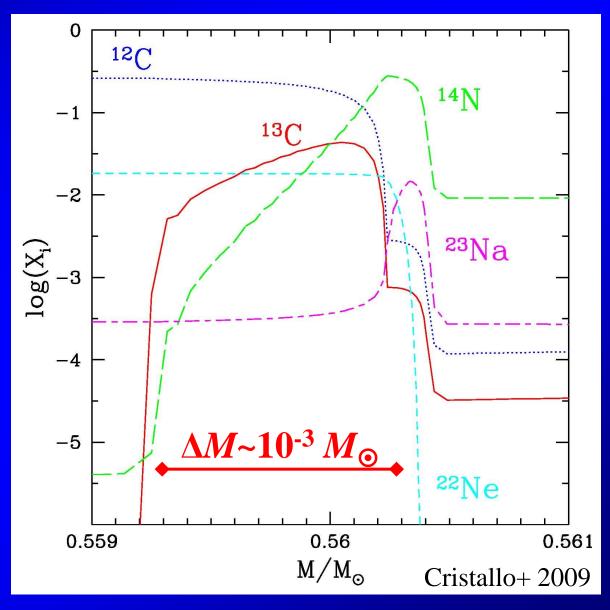


The formation of the ¹³C pocket



¹⁴N strong neutron poison via
¹⁴N(n,p)¹⁴C reaction





The ¹³C pocket in stellar evolutionary models

- ✓ Opacity induced overshoot (SC+...)
- ✓ Convective Boundary Mixing + Gravity Waves (Battino+ 2017)

The ¹³C pockets in post-process calculations:

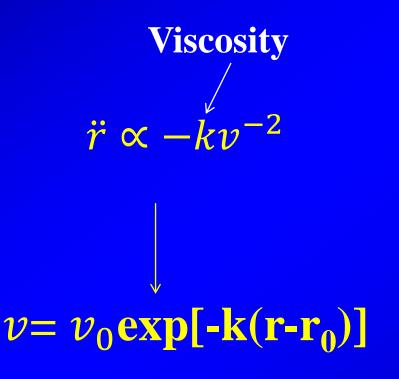
✓ n-zones profile (Gallino+...)
 ✓ Exponential hydrogen profile (Lugaro+...)
 ✓ Magnetic-induced mixing (Trippella+ 2014)

How does the ¹³C pocket change?

 Rotation-induced mixing (Herwig+ 2003; Siess+ 2004; Piersanti+ 2013)

Opacity induced overshoot: a ballistic approach

Let's assume that the deceleration is proportional to the square of the velocity, as it happens to a body moving in a sufficiently dense fluid:



 $v = v_{bce} \cdot exp (-d/\beta H_p)$

- V_{bce} is the convective velocity at the inner border of the convective envelope (*CE*)
- d is the distance from the CE
- *H_p* is the scale pressure height
- **β** = 0.1

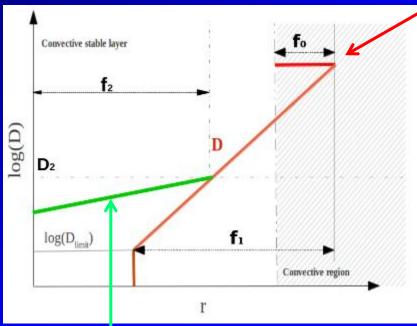
F.R.U.I.T.Y.

		P. all sugar	
MODEL SELECTION	OUTPUT SELECTION	OUTPUT FORMAT	
Mass (M₀) ❤	Nuclides Properties	Multiple Table format ⁽¹⁰⁾	Single Table format ⁽¹¹⁾
 Metallicity (<i>Z</i>) ⁽¹⁾	O Elements (3,4) Z: All Isotopes (5) A: All Z: All s-process (6): [hs/ls], [Pb/hs], Instance Instance	 All Dredge Up Episodes⁽¹²⁾ Final Composition 	Final Composition
 0 ∨ ¹³ C Pocket ⁽⁹⁾ Standard ∨	● Net ⁽⁸⁾ Yields ⁽⁷⁾ A: All Z: All ● Total	• Final	Final

On line at www.oa-abruzzo.inaf.it/fruity

Convective Boundary Mixing + Gravity Waves

Battino+ 2016

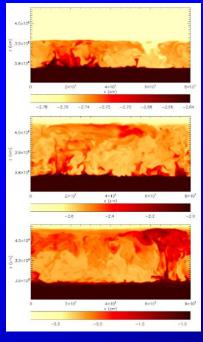


Kelvin-Helmholtz (shear) instability

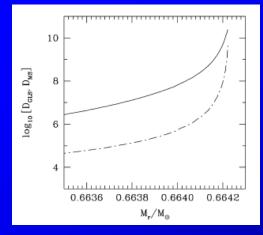
Casanova+ 2016

Depending on the velocity difference across the interface, K-H instability may induce mixing if:

N²/(dv/dr)²<0.25

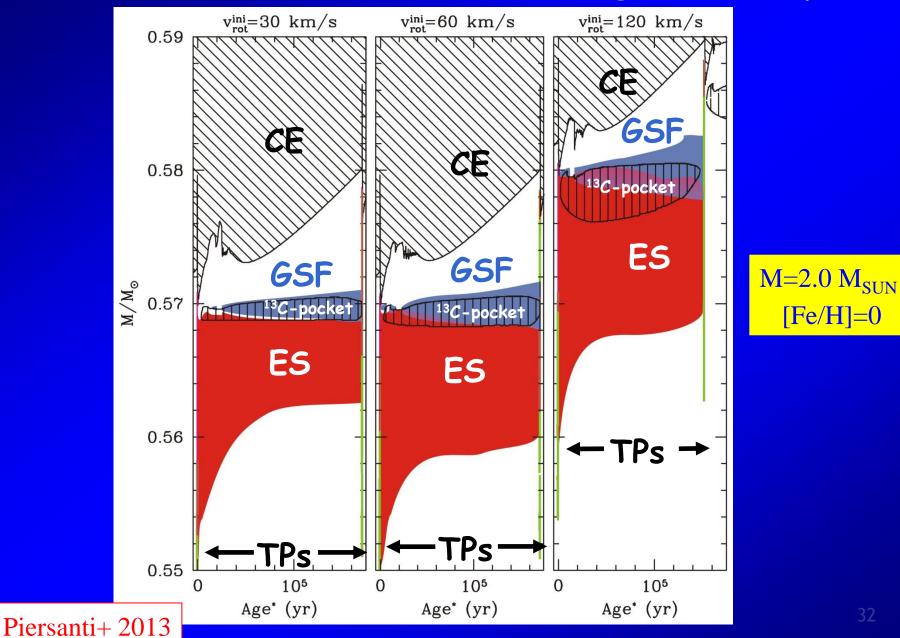


Gravity waves



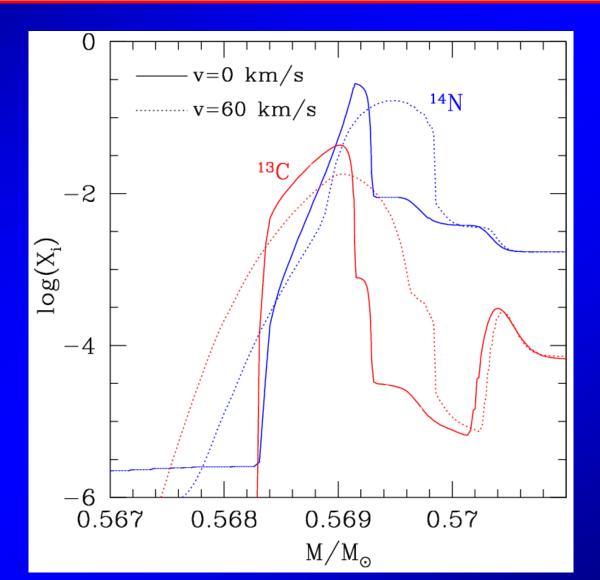
Denissenkov & Tout 2003

Rotation induced instabilities during the AGB phase



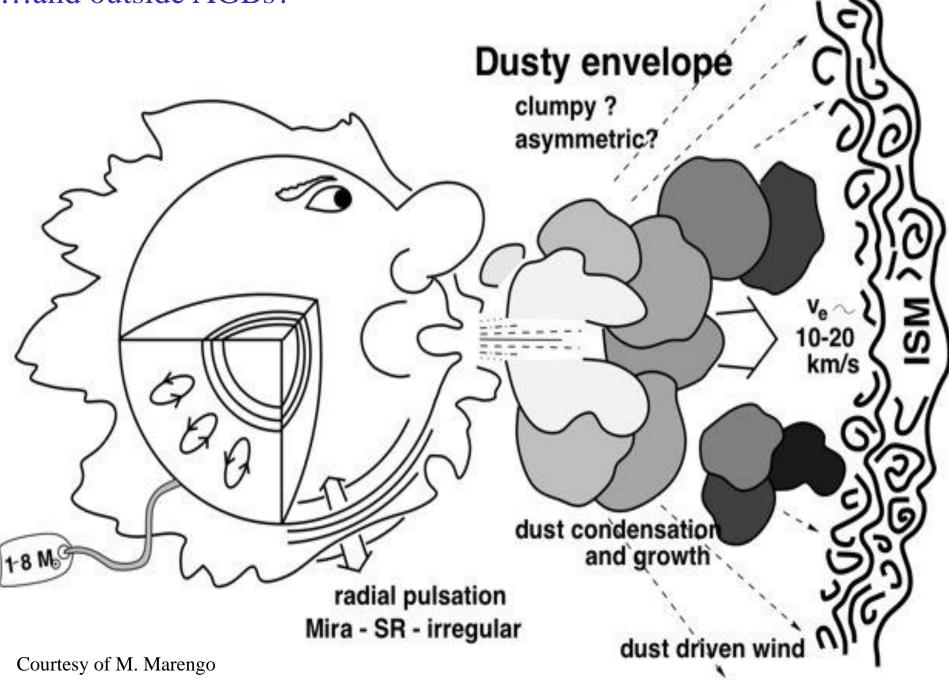
NET EFFECT

It mixes ¹⁴N in ¹³C-rich layers (and viceversa), thus implying a decrease of the local neutron density and an increase of the iron seeds. As a consequence, the surface s-process distributions change.



33

...and outside AGBs?





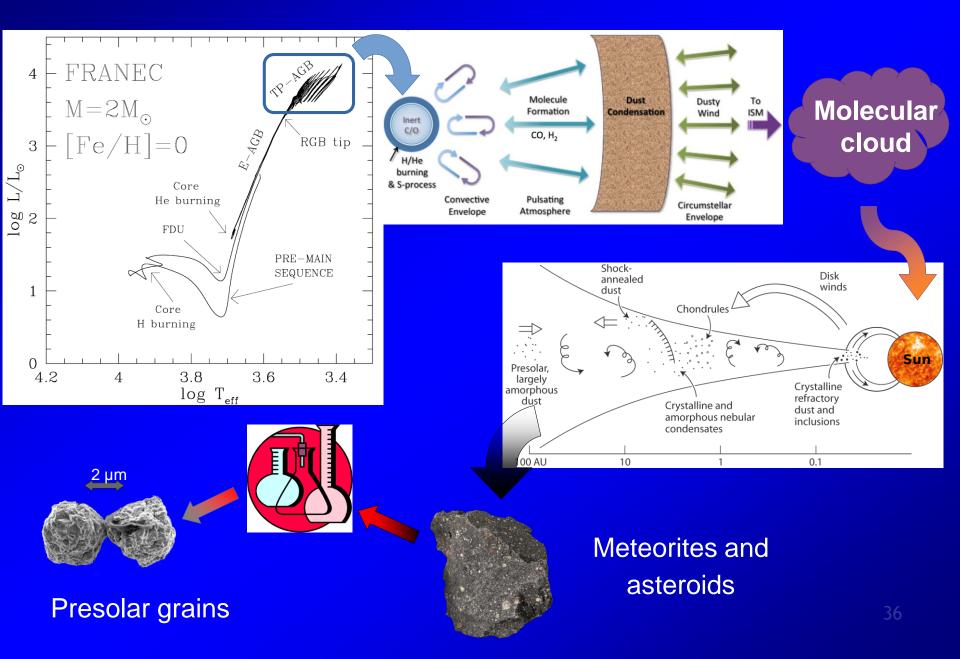
Allende (Mexico, 1969)

Meteorites

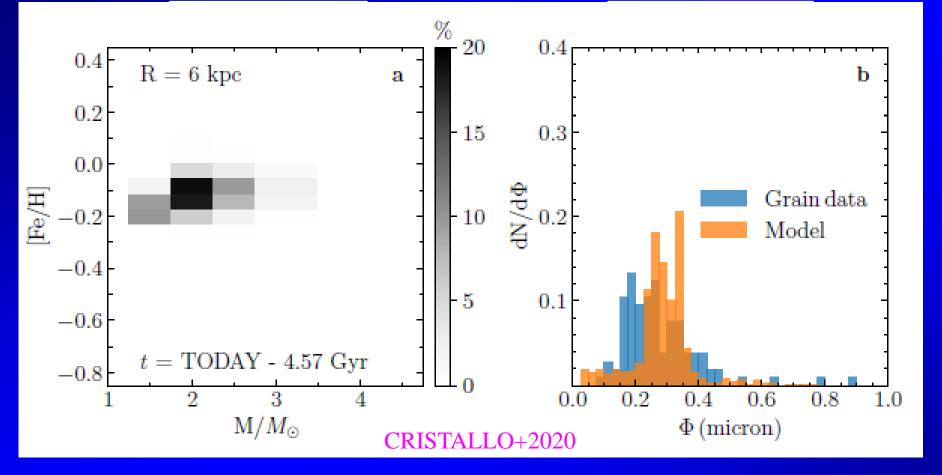
Murchison (Australia, 1969)



AGB stars and presolar SiC grains



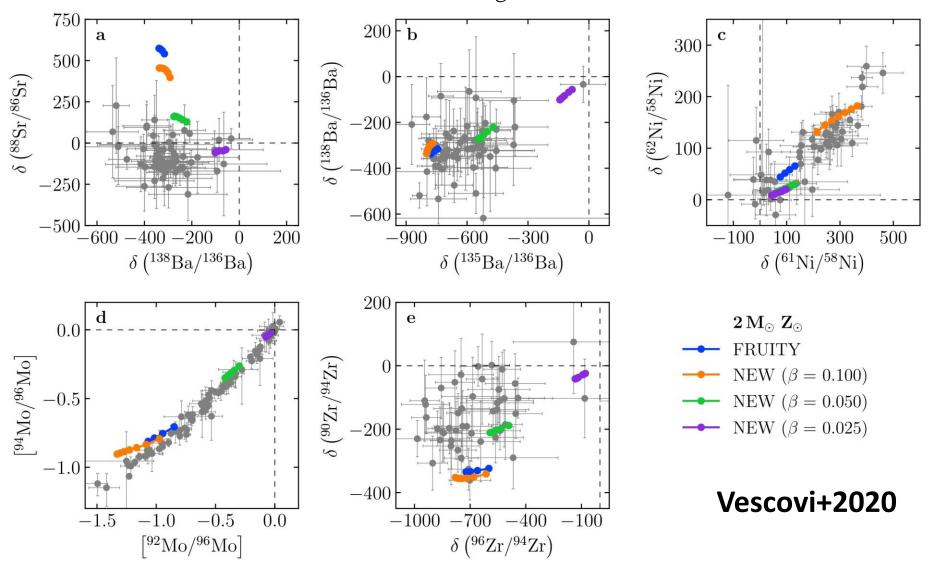
Important question is: which stars are progenitor of pre-solar SiC grains?



Our models predict that the SiC production at the epoch of the Solar System formation is dominated by contributions from AGB stars with $M\approx 2M_{SUN}$ and $Z\approx Z_{SUN}$, which are thus likely the parent stars of presolar SiC grains identified in extraterrestrial materials.

SiC Grains I

- Isotopic data including Ni, Sr, Zr, Mo, and Ba isotope ratios in presolar SiC grains
- Stellar models with same initial mass (2 M_{\odot}) and solar metallicity



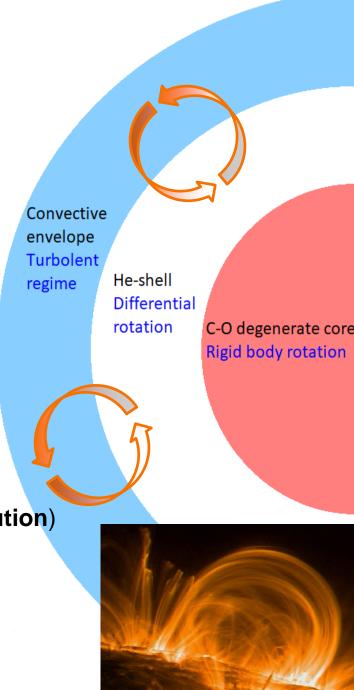
The ¹³C-pocket: formation

• Protons can penetrate into the He-rich region at each TDU (Third Dredge-Up) phenomenon

Which is the physical mechanism? TOP-DOWN MECHANISMS

- Opacity-induced overshoot (Cristallo+2009,2011,2015)
- Convective Boundary Mixing (Battino+2016)
 BOTTOM-UP MECHANISMS
- Magnetic fields (Trippella+2016; Palmerini+2018)
- MagnetoHydroDynamics (MHD) solutions
- (Nucci & Busso 2014):
 - No numerical approximations (exact analytic solution)
 - → Simple geometry: toroidal magnetic field

$$\rho(r) = \frac{\rho_p}{r_p^k} r^k$$



Magnetic-buoyancy-induced mixing

→ **Magnetic** contribution (Vescovi+2020) to the dowflow velocity $v_{d'}$, acting when the density distribution is $\rho \propto r^{k}$:

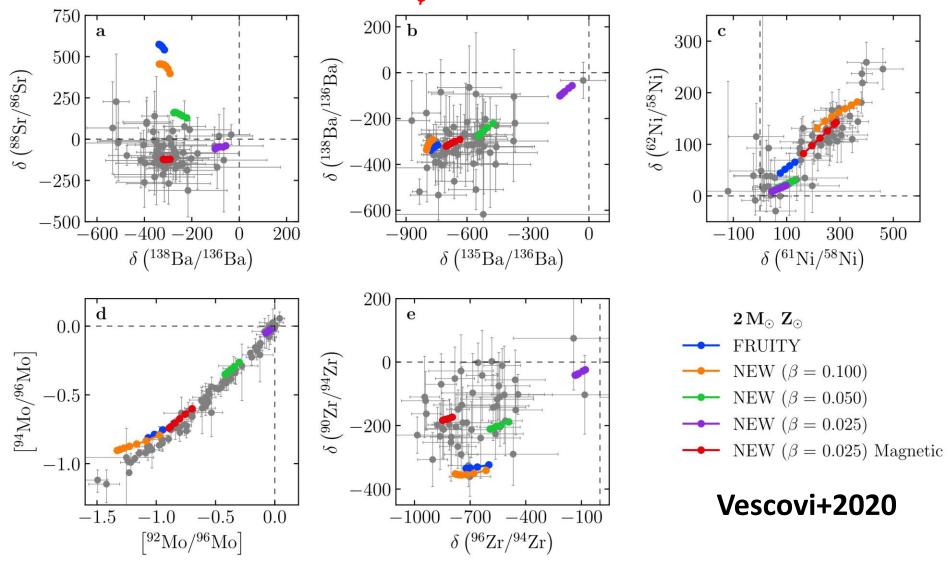
$$\rightarrow v_d(r) = u_p \left(\frac{r_p}{r}\right)^{k+2}$$

- Parameters:
 - Layer "p" at the deepest coordinate from which buoyancy starts
 - (can be identified from the corresponding critical toroidal B_o value)
 - Starting velocity u_p of the buoyant material

$$\implies B_{\varphi} \gtrsim \left(4\pi\rho r N^2 H_{\rm p} \frac{\eta}{K}\right)^{1/2}$$

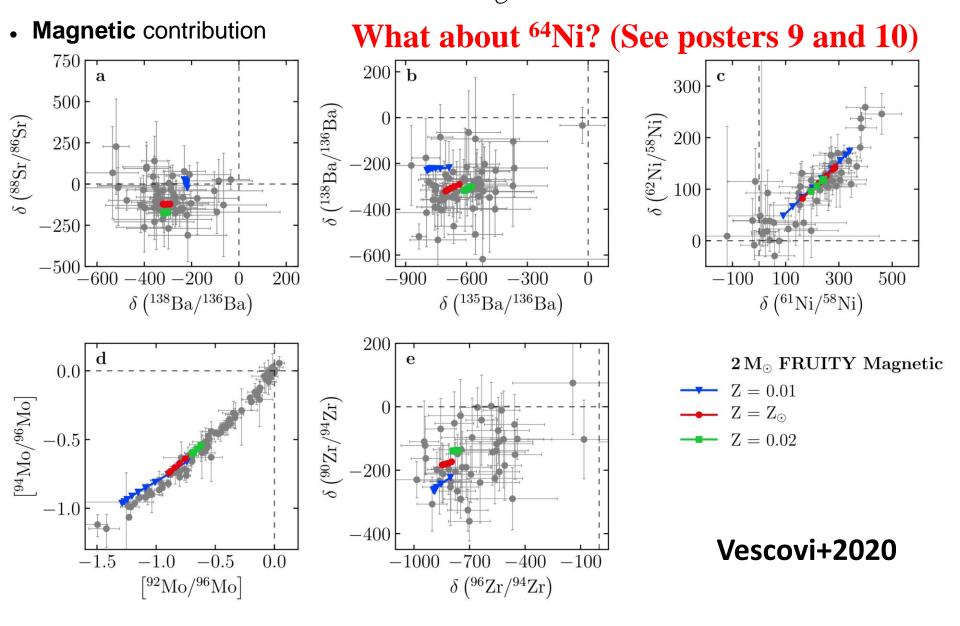
SiC Grains II

- **Magnetic** contribution account for SiC data!!
- Best fit for $u_p = 5 \times 10^{-5}$ cm/s and $B_{\omega} = 5 \times 10^4$ G

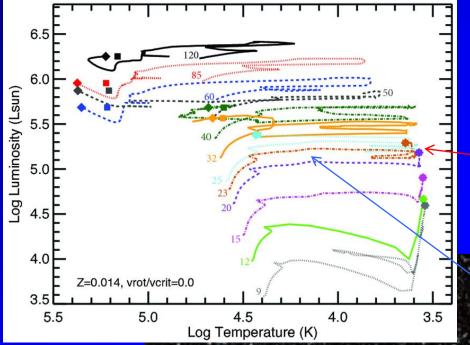


SiC Grains III

• Stellar models with same initial mass (2 M_o) and close-to-solar metallicity



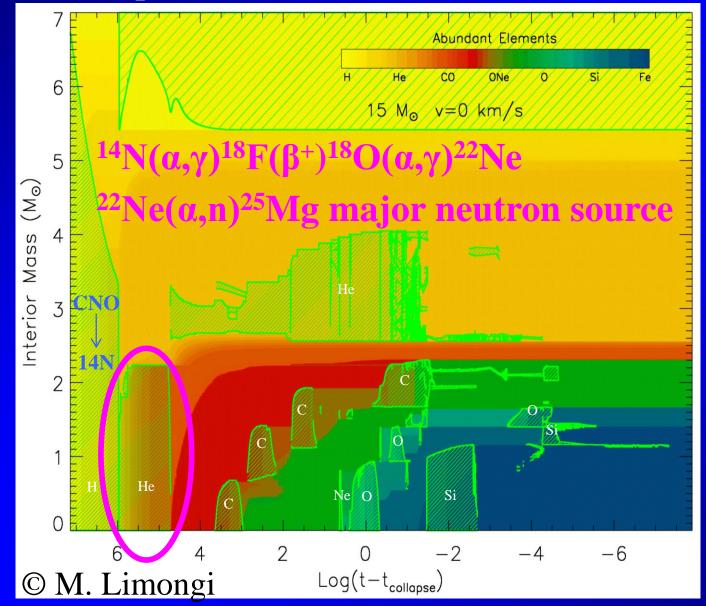
The weak s-process in massive stars



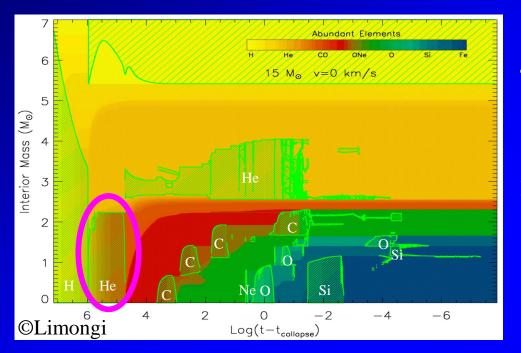




The weak s-process and the evolution of massive stars



Core He-burning phase

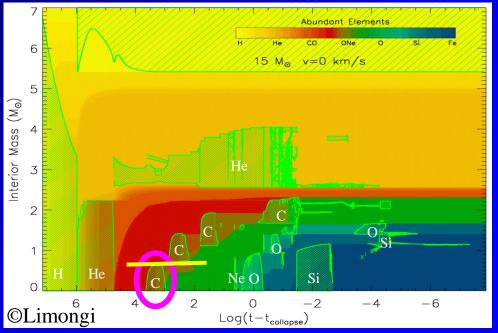


 $3a \rightarrow {}^{12}C$ $12C(a,\gamma)^{16}O$ $14N(a,\gamma)^{18}F(\beta^{+})^{18}O(a,\gamma)^{22}Ne$ $\tau \approx 1 Myr$

When $T \sim 3x 10^8$ K the ²²Ne(α ,n)²⁵Mg is efficiently activated

The resulting neutron density is low (~10⁶ n/cm³) Similar to the s-process

Core C-burning phase



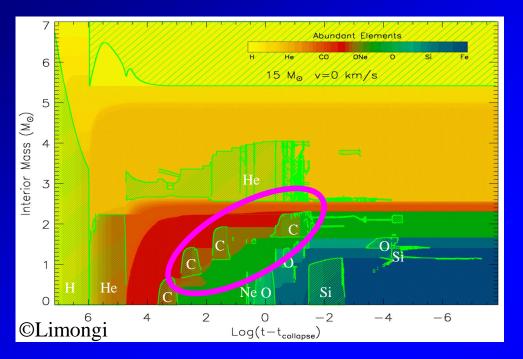
 $\frac{12C(12C,\alpha)^{20}Ne}{12C(12C,p)^{23}Na}$ $\frac{12C(12C,p)^{23}Na}{12C(12C,n)^{23}Mg^{*}}$ $\tau \approx 1 \text{ Kyr}$

Some ²²Ne is left after He burning

All (α ,n) channels are activated: ¹³C(α ,n)¹⁶O - ¹⁷O(α ,n)²⁰Ne ¹⁸O(α ,n)²¹Ne - ²¹Ne(α ,n)²⁴Mg ²²Ne(α ,n)²⁵Mg - ²⁵Mg(α ,n)²⁸Si ²⁶Mg(α ,n)²⁹Si

The resulting neutron density is very high, BUT...

Shell C-burning phase



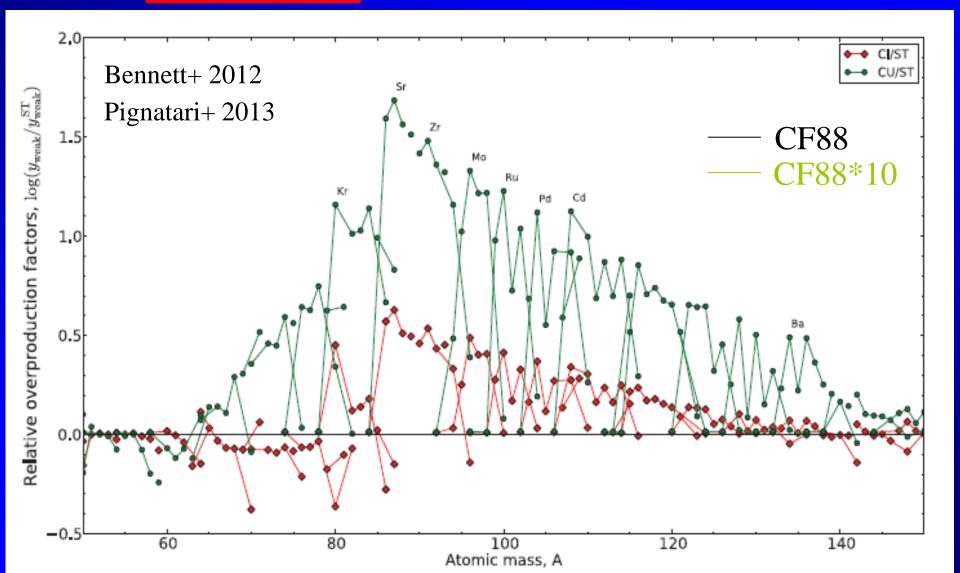
All (α ,n) channels are activated: ¹³C(α ,n)¹⁶O - ¹⁷O(α ,n)²⁰Ne ¹⁸O(α ,n)²¹Ne - ²¹Ne(α ,n)²⁴Mg ²²Ne(α ,n)²⁵Mg - ²⁵Mg(α ,n)²⁸Si ²⁶Mg(α ,n)²⁹Si ¹²C(¹²C,α)²⁰Ne
¹²C(¹²C,p)²³Na
¹²C(¹²C,n)²³Mg*

Why not the ${}^{13}C(\alpha,n){}^{16}O?$ Because at T~1x10⁹ K the ${}^{13}N(\gamma,p){}^{12}C*$ works!!

The resulting neutron density is very high: 10^{11} - 10^{12} n/cm³ 47

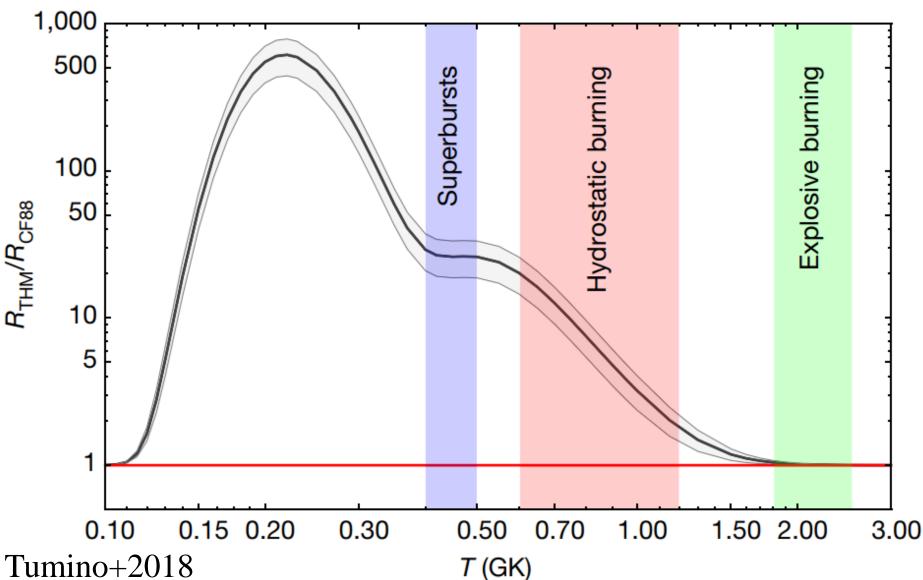
Uncertainties of the weak s-process: cross sections

$^{12}C(^{12}C,x)x - ^{22}Ne(\alpha,x)x - ^{12}C(\alpha,\gamma)^{16}O$



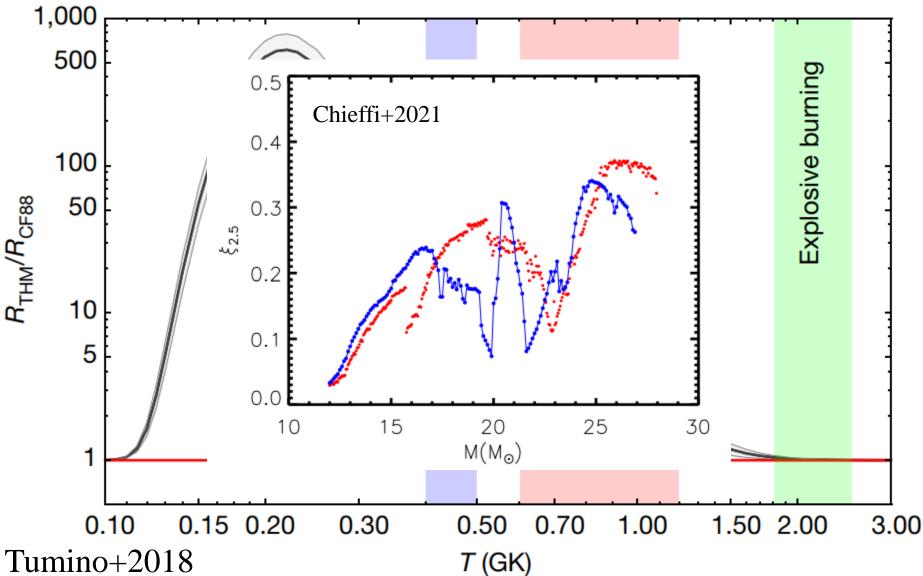
Uncertainties of the weak s-process: cross sections

 $12C(12C.x)x - 22Ne(\alpha.x)x - 12C(\alpha.y)^{16}O$



Uncertainties of the weak s-process: cross sections





Uncertainties of the weak s-process: stellar modelling

Convection - Rotation

Strong production of primary ¹⁴N at low metallicities

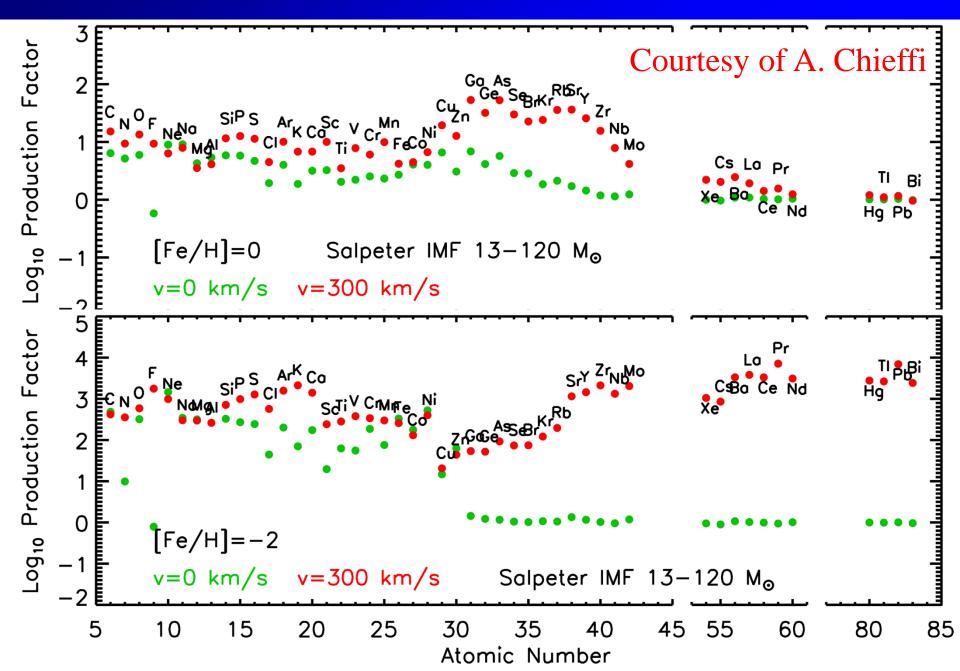
$^{13}{\rm C}/^{14}{\rm N}\simeq~5.7\cdot10^{-3}$

In any case the dominant source is the $^{22}Ne(\alpha,n)^{25}Mg$

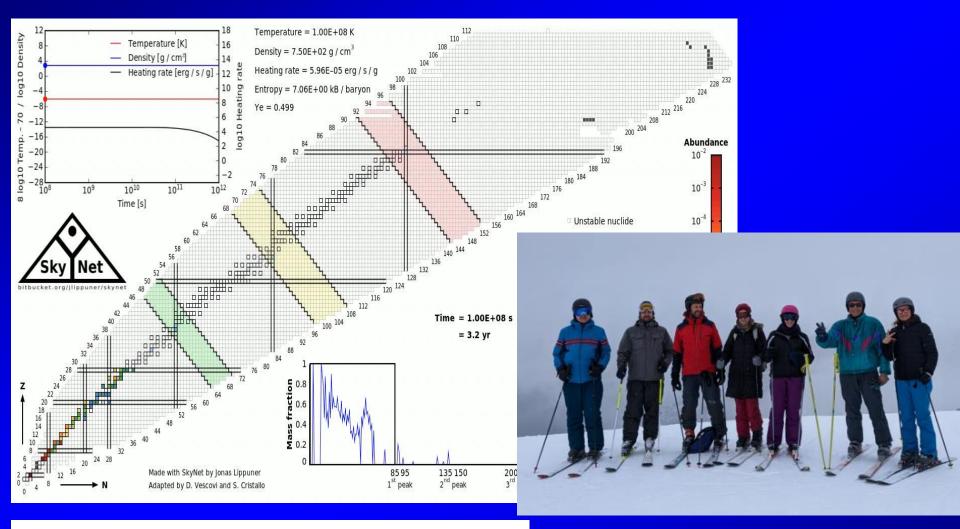
He burning

Courtesy of A. Chieffi

The effect of rotation: differences in the stellar ejecta



THAT'S ALL FALKS!!!



SAPIENZA UNIVERSITÀ DI ROMA

Phd in ASTRONOMY ASTROPHYSICS AND SPACE SCIENCE



ISTITUTO NAZIONALE DI ASTROFISICA OSSERVATORIO ASTRONOMICO D'ABRUZ