Studying plasma induced variation of bound state β-decays in PANDORA

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Stars in a Bottle: PANDORA



What we need & What we have



Steps involved



Step 1: Selecting isotope and transitions



Step 2: Calculating lepton phase volume

The lepton phase volume quantifies the number of ways a decay can occur. The phase volume changes with variations in atomic configuration, depending on type of decay

Information needed: level probability distribution (LPD), orbital occupancy, orbital electron wavefunction, decay energy and shape factor



Step 2 (a): Decay energetics

 $Q = Q_0 + (E_{X,K}^* - E_{Y,K'}^*) + (\epsilon^{i,j} - \epsilon^{i',j'}) + (\Delta_X - \Delta_Y)$

The decay energy depends on not just the difference in nuclear masses, but on the overall system energy which includes atomic/ionic energy



⁷Be - ⁷Li level coupling schematic for K- and L-shell capture (neutral ion)

All configurations should have at
least one K-shell electronNo selection on daughter configuration
K-shell vacant states autoionising

Information needed: energy of different atomic configurations of parent system and coupling with daughter system



All configurations should have at least one L-shell electron

No selection on daughter configuration L-shell vacant states autoionising

K-SHELL CAPTURE

L-SHELL CAPTURE

Step 2 (a): Decay energetics

$$Q = Q_0 + (E_{X,K}^* - E_{Y,K'}^*) + (\epsilon^{i,j} - \epsilon^{i',j'}) + (\Delta_X - \Delta_Y)$$



⁷Be decay Q-value as a function of charge state and level for neutral and 1+ (left and 2+ and 3+ (right)

Step 2 (b): Shape factor

Conservation of total angular momentum implies that only certain electron orbitals can interact with the nucleus, depending on the spin and parity of the decay

Information needed: spin-parity of electron orbitals and decay transition

$$S_{(m)x} = \begin{cases} 1 & \text{for } m = a, nu \text{ and } x = ns_{1/2}, np_{1/2} \\ q^2 & \text{for } m = u \text{ and } x = ns_{1/2}, np_{1/2} \\ 9/R^2 & \text{for } m = u \text{ and } x = np_{3/2}, nd_{3/2} \\ 0 & \text{otherwise.} \end{cases}$$

Step 2 (c):
$$[f_x \text{ or } g_x]^2$$

The probability of electron capture from bound states depends on the square of the radial component of the orbital wavefunction evaluated on the nuclear surface

Information needed: formalism for radial wavefunctions of different orbitals

Radial component of Dirac equation – Coupled differential equations

$$\frac{\mathrm{d}P(r)}{\mathrm{d}r} = -\frac{\kappa}{r}P(r) - \left(2c + \frac{V-\epsilon}{c}\right)Q(r) \qquad \frac{\mathrm{d}Q(r)}{\mathrm{d}r} = \frac{\kappa}{r}Q(r) + \left(\frac{V-\epsilon}{c}\right)P(r)$$

$$\mathbf{For \, V = \mathbf{Z/r \ (in \ atomic \ units)}}$$

$$P(r) = \left(1 - \frac{\epsilon}{c^2}\right)^{1/2} \xi\left(\frac{\rho}{N}\right)^{\gamma} \mathrm{e}^{-\rho/2N} \left[-n_r F_1 + (N-\kappa)F_2\right]$$

$$Q(r) = \left(\frac{\epsilon}{c^2}\right)^{1/2} \xi\left(\frac{\rho}{N}\right)^{\gamma} \mathrm{e}^{-\rho/2N} \left[n_r F_1 + (N-\kappa)F_2\right]$$

- N = apparent principal quantum number
- *n_r* = number of nodes in orbital spatial distribution
- $\xi = normalisation$
- *ρ* = 2Zr = radial function
- *ε* = quantised electron energy
- $\kappa = -(j+1/2)a, a=\pm 1$
- F_{1} , F_{2} = confluent hypergeometric functions

Step 2 (c): $[f_x \text{ or } g_x]^2$ (benchmarking)



[5] V. M. Burke and I. P. Grant, Proceedings of the Physical Society 90 (1967) 297[6] S. Liu, C. Gao and C. Xu, Physical Review C 104 (2021), 024304

Step 2 (c): $[f_x \text{ or } g_x]^2$



Larger of f_x^2 or g_x^2 as calculated for ⁷Be taking R = $R_0A^{1/3}$

Step 2 (d): Occupancy and Level Dependent Lepton Phase Volumes



Step 2 (e): Level probability distributions

The ion CSD and LPD strongly depends on electron density and temperature

Information needed: CSD and LPD of ⁷Be for various n_e and T_e (calculated using FLYCHK)





LPD of ⁷Be⁰⁺ and ⁷Be³⁺ for different temperatures as calculated by FLYCHK (no effect of density) CSD and LPD calculated using grid of density and temperature values in FLYCHK under LTE approximation



CSD of ⁷Be for different temperatures as calculated by FLYCHK (no effect of density)

Step 3: Calculating decay rates

⁷Be is expected to show a rise in $T_{1/2}$ with increase in T_{e} high temperatures lead to greater ionisation and fewer electrons left to capture



The methodology also benchmarked with decay rates in fully stripped ¹⁶⁶Dy in Liu et al

[6] S. Liu, C. Gao and C. Xu, Physical Review C 104 (2021), 024304

Results



Conclusions and Next Steps

- Radio-isotopes like ⁷Be can show modification of lifetimes when ionised or excited
- Results show increase in lifetime due to absence of bound electrons for capture at eV temperatures, much lower than stellar temperatures
- 7Be in stellar interior proceeds by continuum capture [8]
- The recipe for calculating in-plasma bound state decay and bound state capture is ready and tested

$$f_{IF(m)}^{*} = \underbrace{\sum_{ij} p_{ij}}_{x(ij)} \sum_{x(ij)} \sigma_{x} \frac{\pi}{2} [g_{x} \, or \, f_{x}]^{2} (Q(ij)/m_{e}c^{2})^{2} S_{(m)x(ij)}$$

- CSD and LPD calculated under LTE approximation here, but can be calculated under NLTE approximation in ECR plasma
- 3D Particle-in-Cell Monte Carlo code ready for the purpose
- Number of levels *ij* can be reduced by level grouping based on similarity in lepton phase volume contribution

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Alessio Galatà



Alberto Mengoni

...and the PANDORA collaboration

THANK YOU FOR YOUR ATTENTION!



Additional Content



Additional Content







Additional Content

$$\begin{split} \frac{\mathrm{d}n_{i}^{R}}{\mathrm{d}t} &= \left[n_{i-1}^{R}n_{e}\gamma_{i-1,i} - n_{i}^{R}n_{e}\gamma_{i,i+1}\right] + \left[n_{i+1}^{R}\sum_{j=0}^{N-1}n_{j}E_{i+1,i} + n_{i-1}^{R}\sum_{j=1}^{N}n_{j}E_{i-1,i}\right] - \\ &\left[n_{i}^{R}\sum_{j=0}^{N-1}n_{j}E_{i,i-1} + n_{i}^{R}\sum_{j=1}^{N}n_{j}E_{i,j+1}\right] - \frac{n_{i}^{R}}{\tau_{i}^{R}} \\ \frac{\mathrm{d}n_{ii'}^{R}}{\mathrm{d}t} &= \left[\sum_{i,j' < ii'}n_{ij'}^{R}n_{e}C_{ij'ii'} - n_{ii'}^{R}\sum_{i,j' > ii'}n_{e}C_{ij'ij'}\right] + \left[\sum_{i,j' > ii'}n_{ij'}^{R}A_{ij'ii'} - n_{ii'}^{R}\sum_{i,j' < ii'}A_{ii'ij'}\right] \end{split}$$