

Global solar coronal models driven with Alfvén and kink waves

Tom Van Doorselaere

Centre for mathematical Plasma Astrophysics, Department of Mathematics,
KU Leuven

4 July 2025

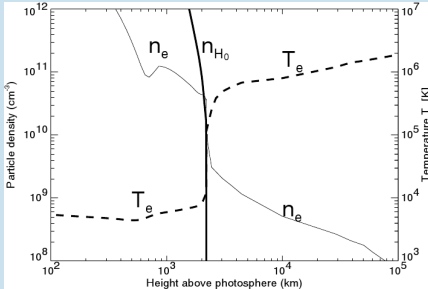
In collaboration with: Norbert Magyar, Valeria Sieyra, Marcel
Goossens, Max McMurdo, Luka Banović



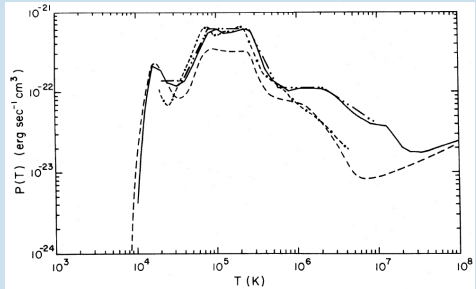
SOUL is a Methusalem project of the Flemish government and KU Leuven



Coronal heating



Aschwanden (2004)



Rosner et al. (1978)

Low density $\rightarrow$ corona is optically thin: energy losses as $n^2 P(t)$

$T = 10^6 \text{ K}, n = 10^{15} / \text{m}^3 \rightarrow$ energy loss $\approx 10^{-5} \text{ J/m}^3 \text{ s}$
internal energy $\approx 2 \cdot 10^{-2} \text{ J/m}^3$

Time scale for cooling $\approx 2000 \text{ s} = 33 \text{ min}$

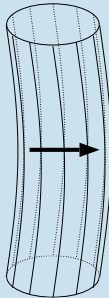
Need quick heat replenishing! $\rightarrow$ coronal heating problem

Alfvén vs. kink

VD et al. (2008): kink and Alfvén waves in coronal loops



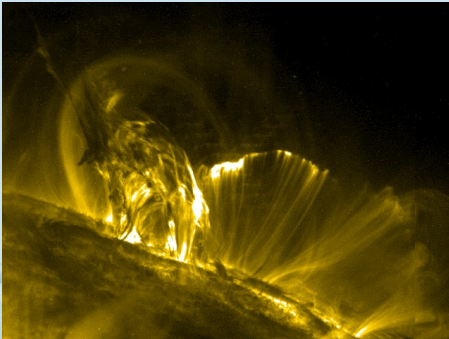
- incompressible
- transverse to $\vec{B}$
- torsional
- $V_A^2 = \frac{B^2}{\mu\rho}$
- homogeneous OK



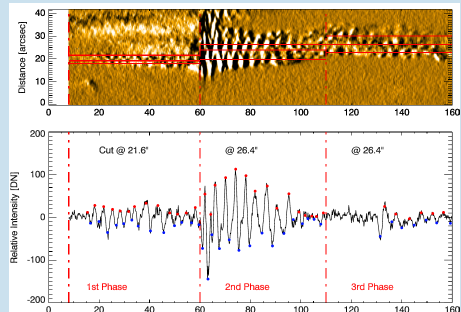
- incompressible
- transverse to $\vec{B}$
- kinking
- $V_k^2 = \frac{2B^2}{\mu(\rho_i + \rho_e)}$
- density (or $\vec{V}$ or $\vec{B}$) structure essential

Kink waves in loops

Decaying:
excited after flare
strong damping

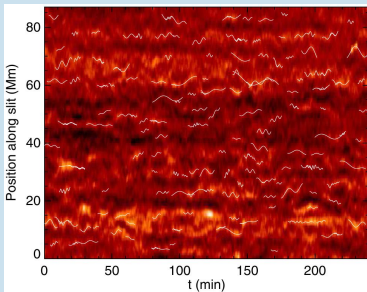
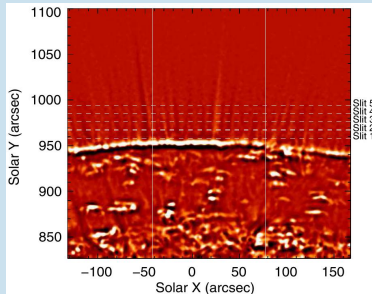


Decayless:
no apparent excitation
no apparent damping



Nisticò et al. (2013)

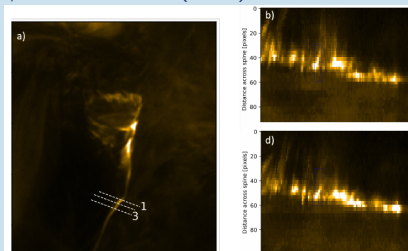
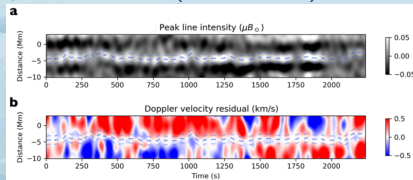
Kink and Alfvén waves in coronal plumes



Thurgood et al. (2014) ↑

↓ Petrova et al. (2024): SolO

↓ Morton et al. (2025, subm.): DKIST

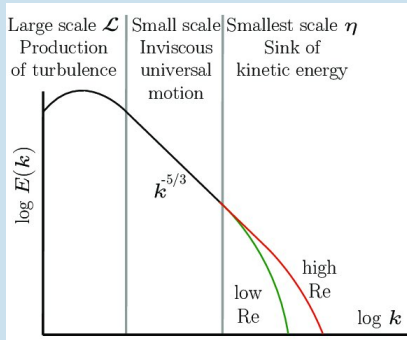


Alfvén turbulence

Propagation of Alfvén waves with Elsässer variables $\vec{z}^{\pm} = \vec{v} \pm \frac{\vec{B}'}{\sqrt{\mu\rho}}$.

Governing equations (incompressible MHD, no P'_T for Alfvén, $\vec{V}_0 = 0$):

$$\frac{\partial \vec{z}^{\pm}}{\partial t} \mp \vec{V}_A \cdot \nabla \vec{z}^{\pm} = -\vec{z}^{\mp} \cdot \nabla \vec{z}^{\pm}$$



Stein (2013)

AWSOM

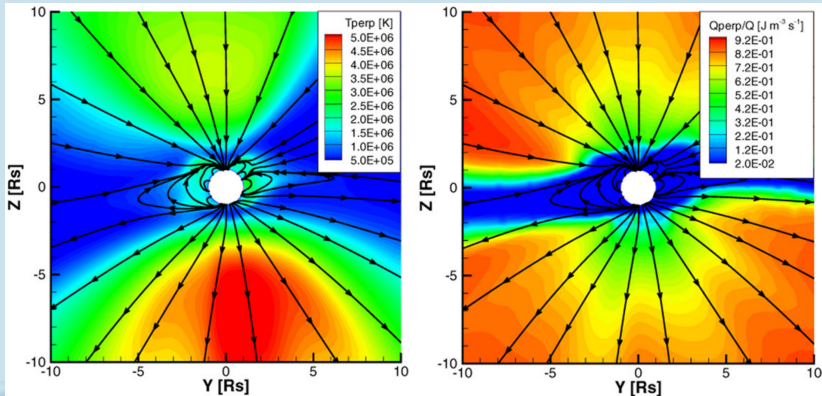
Van der Holst et al. (2014): Alfvén Wave Solar Model

- Starts from MHD/2-fluids
- two extra equations for $w_A^\pm$: advection, reflection, cascade
- includes extra force due to Alfvén wave pressure P_A
- includes extra heating terms due to cascade

$$E = \frac{p}{\gamma - 1} + \rho \frac{v^2}{2} + \frac{B^2}{2\mu} + W_A^+ + W_A^-$$

$$\begin{aligned} \frac{\partial W_A^\pm}{\partial t} + \nabla \cdot ((\vec{V} \mp \vec{V}_A) W_A^\pm) + \frac{W_A^\pm}{2} \nabla \cdot \vec{V} = & -\frac{1}{L_\perp} \frac{1}{\sqrt{\rho}} \sqrt{W_A^\mp} (W_A^\pm) \\ & \pm \mathcal{R} \sqrt{W_A^+ W_A^-} \end{aligned}$$

AWSOM

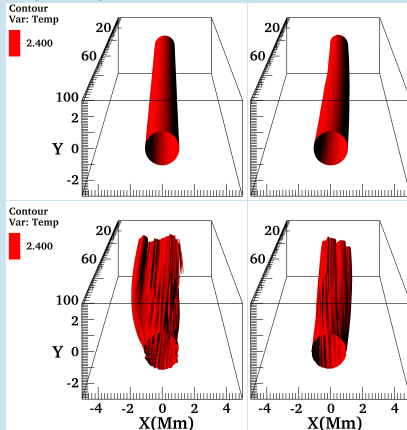


Van der Holst (private communication): problems in open-field regions, not enough driving of wind

Cooper Downs (private communication): extra heating term needed in low corona to match observations

Turbulence in transverse kink waves

Karampelas et al. (2019)



Non-linear damping of kink waves: power law and heating

How to add kink heating to AWSOM?

Q-variables

Van Doorselaere et al. (2024): Introduce new Q-variables

$$\vec{Q}^{\pm} = \vec{V} \pm \alpha \vec{B},$$

($\vec{V}$ is velocity, $\vec{B}$ is magnetic field, α is parameter)

→ extension to Elsässer variables

$$(\vec{Z} = \vec{V} \pm \vec{B}/\sqrt{\mu\rho}, \text{ i.e. } \alpha = 1/\sqrt{\mu\rho})$$

→ $\alpha\vec{B}$ is phase speed of wave

$$\frac{D^{\pm}}{Dt} = \frac{\partial}{\partial t} + \vec{Q}_0^{\pm} \cdot \nabla = \frac{\partial}{\partial t} + (\vec{V}_0 \pm \alpha \vec{B}_0) \cdot \nabla$$

Allows to track other waves than Alfvén waves

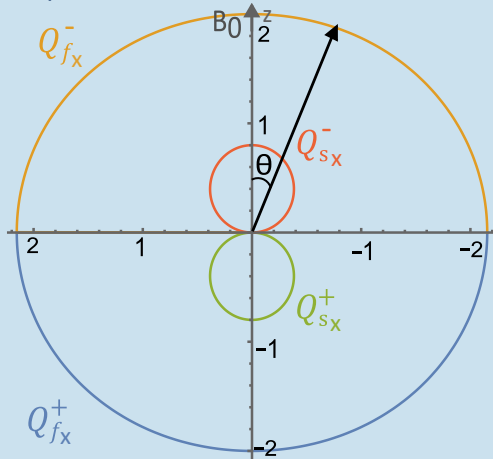
e.g. kink waves in inhomogeneous plumes

Split between propagating up or down

→ Generalisation of: characteristics, Riemann invariants,
Frieman-Rosenbluth, Elsässer variables

Q-variables for slow and fast waves

Numerically compute Q for slow and fast waves



Q split between propagation directions!

Induction equation

Linearised induction equation:

$$\frac{\partial \vec{B}'}{\partial t} = \nabla \times \left((\vec{V} + \vec{v}) \times \vec{B} \right)$$

Take $\vec{B} = B(x, y)\vec{e}_z$, $\vec{V} = V(x, y)\vec{e}_z$, Fourier analyse in z and t :

$$\frac{\partial \vec{B}'_{\perp}}{\partial t} = B(x, y) \frac{\partial \vec{v}_{\perp}}{\partial z} - V(x, y) \frac{\partial \vec{B}'_{\perp}}{\partial z}, \quad -i\omega \vec{B}'_{\perp} = ik_z B(x, y) \vec{v}_{\perp} - ik_z V \vec{B}'_{\perp}$$

Rearrange using $\vec{v} = \frac{1}{2}(\delta \vec{Q}^+ + \delta \vec{Q}^-)$ and $\vec{B}' = \frac{1}{2\alpha}(\delta \vec{Q}^+ - \delta \vec{Q}^-)$:

$$(\omega - k_z V + \alpha k_z B) \delta \vec{Q}_{\perp}^+ = (\omega - k_z V - \alpha k_z B) \delta \vec{Q}_{\perp}^-$$

Choice of α splits waves between directions. For kink waves:

$$V_k^2 = \frac{2B^2}{\mu(\rho_i + \rho_e)}, \text{ so } \alpha = \sqrt{\frac{2}{\mu(\rho_i + \rho_e)}}$$

UAWSOM

Van Doorselaere et al. (2025):

UAWSOM: Uniturbulence and AWSOM

→ make equation for kink wave heating

- Linearise Q-MHD equations
(VD et al. 2024)
- incompressible $\delta\rho = 0$
- only perp. components $\delta\vec{Q}_{\perp}^{\pm}$
- Scalar multiplication with $\rho_0\delta\vec{Q}_{\perp}^{\pm}/2$
→ energy density kink $W_k^{\pm} = \frac{\rho_0}{4}\langle(Q^{\pm})^2\rangle$
- Average over cross-section
→ **Remove small scales**

$$\frac{\partial W_k^{\pm}}{\partial t} + \nabla \cdot (\vec{Q}_k^{\mp} W_k^{\pm}) + \frac{W_k^{\pm}}{2} \nabla \cdot \vec{V} = -\frac{1}{L_{\perp,VD}} \frac{1}{\sqrt{\rho_e}} (W_k^{\pm})^{3/2}$$

Perpendicular length $L_{\perp,VD}$ scales with R , ζ and filling factor f

UAWSOM

Van Doorselaere et al. (2025):

UAWSOM: Uniturbulence and AWSOM

→ make equation for kink wave heating

- Linearise Q-MHD equations
(VD et al. 2024)
- incompressible $\delta\rho = 0$
- only perp. components $\delta\vec{Q}_{\perp}^{\pm}$
- Scalar multiplication with $\rho_0\delta\vec{Q}_{\perp}^{\pm}/2$
→ energy density kink $W_k^{\pm} = \frac{\rho_0}{4}\langle(Q^{\pm})^2\rangle$
- Average over cross-section

→ Remove small scales

Compare with
Alfvén wave:

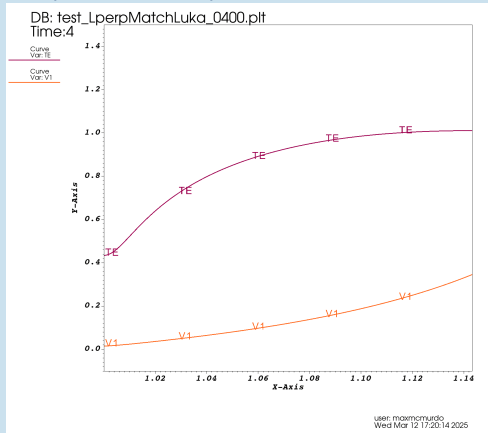
$$\begin{aligned} & \text{RHS}_{\text{Alfvén}} \\ &= -\frac{1}{L_{\perp,A}} \sqrt{\frac{W_A^{\mp}}{\rho_0}} W_A^{\pm} \end{aligned}$$

$$\frac{\partial W_k^{\pm}}{\partial t} + \nabla \cdot (\vec{Q}_k^{\mp} W_k^{\pm}) + \frac{W_k^{\pm}}{2} \nabla \cdot \vec{V} = -\frac{1}{L_{\perp,VD}} \frac{1}{\sqrt{\rho_e}} (W_k^{\pm})^{3/2}$$

Perpendicular length $L_{\perp,VD}$ scales with R , ζ and filling factor f

UAWSOM in MPI-AMRVAC

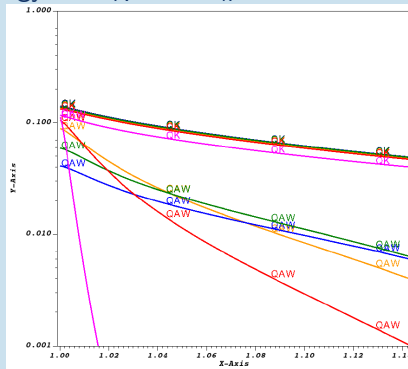
McMurdo et al. (2025, poster):



Kink wave heated equilibrium. No background heating needed.

UAWSOM in 1D

McMurdo et al. (2025, poster):
compare Alfvén and kink heating (in $\text{erg cm}^{-3} \text{s}^{-1}$)
Equal energy in W_A and W_k



Kink wave heating larger than Alfvén wave heating

$$W_A^- \leq W_A^+ < W_k^+$$



Conclusions

- Solar coronal heating problem
- Potential wave heating contributions?
- Coronal structures → kink waves
- New wave tracking formalism: Q -variables
- UAWSOM: kink wave heating model
- kink waves heat more than Alfvén waves

