

Solar neutrino flux fluctuations caused by g modes

Interdisciplinary Physics of the Sun (4th July, 2025)

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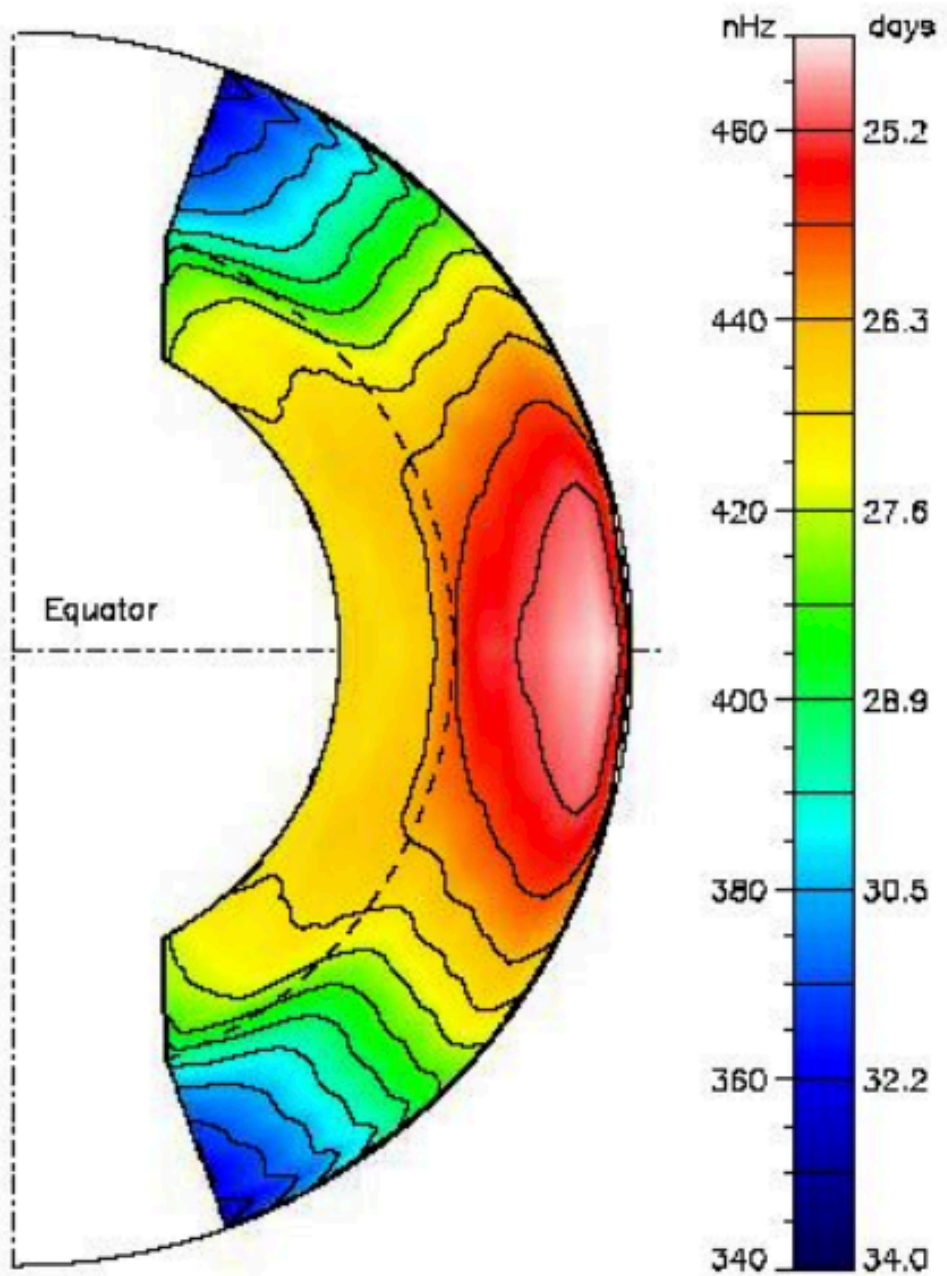
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Helioseismology: the study of solar oscillations to probe its interior

	restoring force	frequency (period)	sensitive to...	Can we detect it?
p mode	pressure	~ 3000 micro Hz (~ 5 min.)	outer envelope (conv. + upper rad.)	Yes (Leighton+1962)
f mode	buoyancy	~ 1500 micro Hz (~ 10 min.)	outer envelope	Yes (Leighton+1962)
inertial mode	Coriolis force	~ 0.1 micro Hz (~ months)	outer envelope (conv. + upper rad.)	Yes (Löptien+2018, Gizon+2021)
g mode	buoyancy	< 100 micro Hz	central region	No, so far...

Ex. 2-d rotational profile
of the Sun



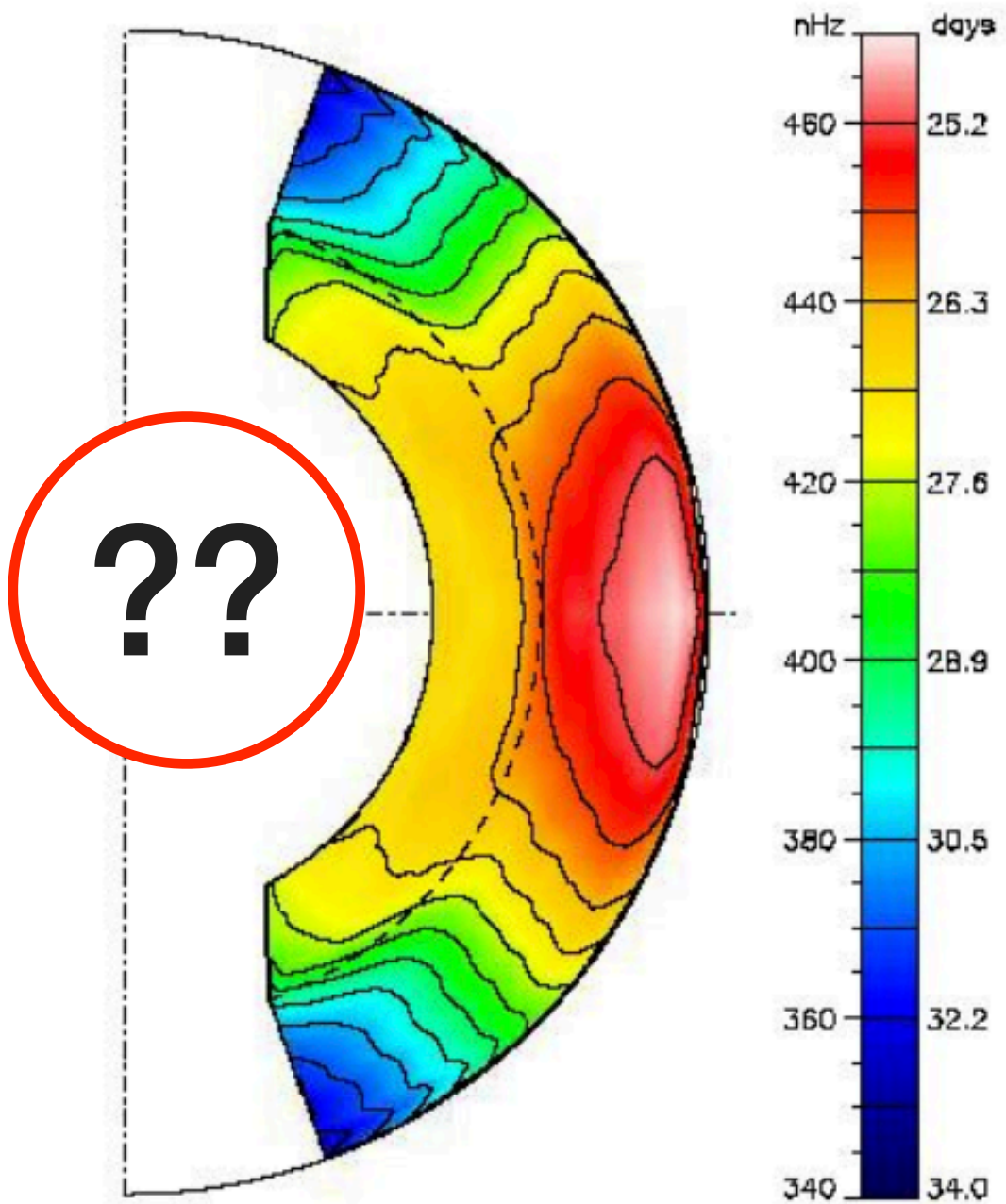
(Thompson+1996)

(Talks by Rachel, Aldo, and Gaël)

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Difficulty in solar g-mode detection

- There have been many attempts of solar g-mode detection, but ... (see a review by Appourchaux+2010)
- The main difficulty is that amplitudes of solar g modes are very small at the surface
 - e.g., a few mm/s in line-of-sight velocity (vlos)
- * Actually, it is theoretically difficult to estimate solar g-mode amplitudes (Belkacem+2022, Pincon+2021)
- We need a few more decades observations (in vlos) to achieve signal-to-noise ratio required for the firm detection of solar g modes (Appourchaux+2010) (→ how about solar neutrinos...??)

(Discussions on solar g-mode detection by LISA)

Why solar neutrino for solar g-mode detection?

- The idea is:
 - solar g modes perturb T and ρ in the central region
 - δT and $\delta \rho$ cause the perturbation in the nuclear reaction rate $\delta \varepsilon$
 - $\delta \varepsilon$ is related to the neutrino flux density fluctuation $\delta \phi$(e.g. Gough 1991, Kumar and Bahcall 1992 in the context of solar neutrino problem)
- Lopes and Turck-Chièze (2014) evaluated flux fluctuations caused by g modes (in the case of ^8B neutrinos), putting some constraints on solar g-mode amplitudes
 - How about other types of neutrinos, such as ^7Be , ^{13}N , ^{15}O , and ^{17}F neutrinos?
(suggested by a Kamiokande experimentalist, Yuuki Nakano)

(Talks by Michael and Francesco)

Goals

- To evaluate neutrino flux fluctuations caused by solar g modes (for various neutrinos such as ^7Be , ^{13}N , ^{15}O , and ^{17}F neutrinos)
- Can we detect solar g modes via solar neutrino flux measurements by neutrino detectors?
(Actually, the answer is not positive one...)

Assumptions in our formulation

- Linear adiabatic oscillations:

$$\delta \ln \rho = (\Gamma_3 - 1)^{-1} \delta \ln T: \quad \delta \rho \rightleftharpoons \delta T$$

- Approximation for nuclear reaction rates (following Lopes & Turck-Chièze 2014):

$$\varepsilon \propto \rho^\beta T^\eta: \quad \delta \varepsilon \rightleftharpoons \delta \rho, \delta T$$

- A linear relation between $\delta \ln \phi$ and $\delta \ln \varepsilon$:

→ $\delta \phi$ is determined by, e.g., δT alone

- Negligence of time-delay effect:

- All the neutrinos produced at a certain time in the solar interior are assumed to arrive at an observational point at the same time

The first-order perturbation

- The neutrino flux (of the equilibrium state) is:

$$\Phi_0 = \int \phi_0 \rho_0 dV \quad (\phi_0 = \phi_0(r) \text{ is the neutrino flux density in units of } \text{cm}^{-2} \text{ s}^{-1} \text{ g}^{-1})$$

- The first-order neutrino flux fluctuation is:

$$\Delta \Phi_{\text{osc}} = \int \phi' \rho_0 dV + \int \phi_0 \rho' dV$$

- The eigenfunctions may be given as:

$$\phi'(r, \theta, \psi, t) = \sum_{n\ell m} \text{Re}[\phi'_{n\ell}(r) Y_{\ell}^m(\theta, \psi) e^{-i(\omega_{n\ell m} t + \delta_{n\ell m})}]$$

$$\rho'(r, \theta, \psi, t) = \sum_{n\ell m} \text{Re}[\rho'_{r,n\ell}(r) Y_{\ell}^m(\theta, \psi) e^{-i(\omega_{n\ell m} t + \delta_{n\ell m})}]$$

The first-order perturbation... is **zero**!

- The neutrino flux fluctuation is:

$$\begin{aligned}\Delta\Phi_{\text{osc}} &= \int \phi' \rho_0 dV + \int \phi_0 \rho' dV \\ &= \sum_{n\ell m} \int_0^{R_\odot} G_{n\ell}(r) \phi_0 \rho_0 r^2 dr \times \int_0^\pi P_\ell^m(\cos\theta) \sin\theta d\theta \times \int_0^{2\pi} \text{Re}[\exp(-i[\omega_{n\ell m}t + \delta_{n\ell m}] - m\psi)] d\psi \\ &= 0 \text{ (except for } (\ell, m) = (0, 0) \text{). But no } \ell = 0 \text{ g modes exist.)}\end{aligned}$$

$$\phi'(r, \theta, \psi, t) = \sum_{n\ell m} \text{Re}[\phi'_{n\ell}(r) Y_\ell^m(\theta, \psi) e^{-i(\omega_{n\ell m}t + \delta_{n\ell m})}]$$

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What are other possible effects...?

Current assumptions:

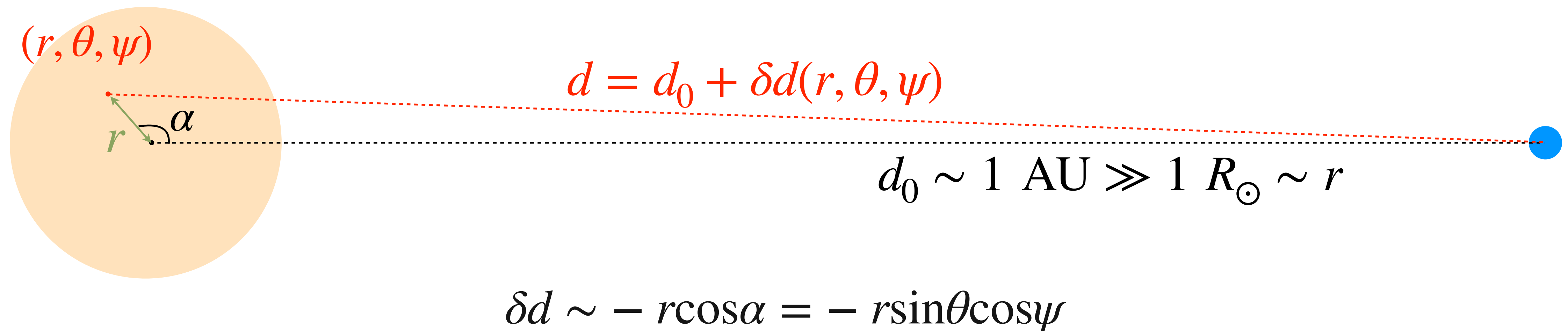
- linear adiabatic oscillation
- $\varepsilon \propto \rho^\beta T^\eta$
- $\delta \ln \phi = \delta \ln \varepsilon = c_1 \delta \ln T$
- Negligence of time-delay effect:
 - All the neutrinos from the solar interior are assumed to arrive at an observational point at the same time

Taking into account the time-delay effects

- The neutrino flux fluctuation is:

$$\Delta\Phi_{\text{osc}}(t) = \int \phi' \rho_0 dV + \int \phi_0 \rho' dV \quad (C: \text{speed of light})$$

$$= \int \phi'(r, \theta, \psi, t - d/c) \rho_0(r) dV + \int \phi_0(r) \rho'(r, \theta, \psi, t - d/c) dV$$



Taking into account the time-delay effects

$$\Delta\Phi_{\text{osc}}(t) = \frac{2\pi}{\sqrt{3}} \sum_n \frac{\omega_{n11}}{c} \int G_{n1}(r) r dr \times \sin(\omega_{n11}t + \delta'_{n11}) - \frac{2\pi}{\sqrt{3}} \sum_n \frac{\omega_{n1,-1}}{c} \int G_{n1}(r) r dr \times \sin(\omega_{n1,-1}t + \delta'_{n1,-1})$$

$$\delta d \sim -r \cos \alpha = -r \sin \theta \cos \psi$$

- Only $(\ell, m) = (1, \pm 1)$ g modes can be non-zero
 - We can evaluate $G_{n\ell}$ with
 - a reasonable **solar model** (we used solar models in Kunitomo+2022)
 - **eigenfunctions** that can be obtained by numerically computing linear adiabatic oscillations of the reference model (we used GYRE; Townsend & Teitler 2013)
 - and the **g-mode amplitude parameter** $A_{n\ell} = |\delta \ln T|_{\text{max}}$ (we assumed $A_{n\ell} = 10^{-5}$)
- The relative fluctuation in the neutrino flux thus evaluated is $< 10^{-10}$
 (for ^8B , ^7Be , ^{13}N , ^{15}O , and ^{17}F neutrinos)

What are other possible effects...?

Current assumptions:

- linear adiabatic oscillation

- $\varepsilon \propto \rho^\beta T^\eta$

- $\delta \ln \phi = \delta \ln \varepsilon = c_1 \delta \ln T \leftarrow$ can be expanded up to the second-order...??

- Negligence of time-delay effect:

- All the neutrinos from the solar interior are assumed to arrive at an observational point at the same time

$$\Delta \Phi_{\text{osc}} = \int \phi' \rho_0 dV + \int \phi_0 \rho' dV + \int \phi' \rho' dV$$

$$\delta \ln \phi = c_1 \delta \ln T + c_2 (\delta \ln T)^2$$

Second-order effect might be non-negligible...??

$$\Delta\Phi_{\text{osc}}(t) =$$

$$\sum_{(n,\ell,m \neq 0)} \sum_{(n'=n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr$$

CONSTANT

$$+ \sum_{(n,\ell,m \neq 0)} \sum_{(n' \neq n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o - o')$$

$$+ \sum_{(n,\ell,m \neq 0)} \sum_{(n',\ell'=\ell,m'=-m)} \pi (-1)^m \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o + o')$$

$$+ \sum_{(n,\ell,m=0)} \sum_{(n',\ell'=\ell,m'=0)} 2\pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o)\cos(o'),$$

(MOSTLY) PERIODIC

where $o(t) = \omega_{n\ell m}t + \delta_{n\ell m}$, **and** $o'(t) = \omega_{n'\ell' m'}t + \delta_{n'\ell' m'}$

Periodic components are still small...

$$\Delta\Phi_{\text{osc}}(t) = \sum_{(n,\ell,m \neq 0)} \sum_{(n'=n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr$$

CONSTANT $10^{-10} \dots$

$$+ \sum_{(n,\ell,m \neq 0)} \sum_{(n' \neq n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o - o')$$

Difficulty in frequency analysis??
(MOSTLY) PERIODIC

$$+ \sum_{(n,\ell,m \neq 0)} \sum_{(n',\ell'=\ell,m'=-m)} \pi(-1)^m \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o + o')$$

$$+ \sum_{(n,\ell,m=0)} \sum_{(n',\ell'=\ell,m'=0)} 2\pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o)\cos(o'),$$

$A_{n\ell} = 10^{-5}$

where $o(t) = \omega_{n\ell m}t + \delta_{n\ell m}$, and $o'(t) = \omega_{n'\ell' m'}t + \delta_{n'\ell' m'}$

Speculation...

$$10^{-10} \times 10^5 \sim 10^{-5}$$

CONSTANT

These terms are just sums of all g modes

$$\begin{aligned} \Delta\Phi_{\text{osc}}(t) = & \sum_{(n,\ell,m \neq 0)} \sum_{(n'=n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \\ & + \sum_{(n,\ell,m \neq 0)} \sum_{(n' \neq n,\ell'=\ell,m'=m)} \pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o - o') \\ & + \sum_{(n,\ell,m \neq 0)} \sum_{(n',\ell'=\ell,m'=-m)} \pi(-1)^m \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o + o') \\ & + \sum_{(n,\ell,m=0)} \sum_{(n',\ell'=\ell,m'=0)} 2\pi \int_0^{R_\star} Q_{(n\ell),(n'\ell')} dr \times \cos(o)\cos(o'), \end{aligned}$$

(MOSTLY) PERIODIC

$$A_{n\ell} = 10^{-5}$$

where $o(t) = \omega_{n\ell m}t + \delta_{n\ell m}$, and $o'(t) = \omega_{n'\ell'm'}t + \delta_{n'\ell'm'}$

Summary and conclusions

- Solar g modes are important
 - Solar neutrino flux measurements could help us to find the solar g modes
 - The first-order fluctuation in the neutrino flux is **zero**
 - We see non-zero first-order fluctuations if we take into account the time-delay effect
 - but the relative fluctuation in the neutrino flux is $< 10^{-10}$ when we assume $A_{n\ell} = 10^{-5}$ (for ${}^8\text{B}$, ${}^7\text{Be}$, ${}^{13}\text{N}$, ${}^{15}\text{O}$, and ${}^{17}\text{F}$ neutrinos)
 - We are evaluating the second-order fluctuation in the neutrino flux
 - again, the relative fluctuation caused by a single g mode is too small to detect $\sim 10^{-10}$ when we assume $A_{n\ell} = 10^{-5}$ (for ${}^8\text{B}$, ${}^7\text{Be}$, ${}^{13}\text{N}$, ${}^{15}\text{O}$, and ${}^{17}\text{F}$ neutrinos)
- it is almost impossible to detect individual g modes via neutrino flux measurements...
- **But, the non-time varying component can be non-negligible ($\sim 10^{-5}$??)**
(when we assume $A_{n\ell} = 10^{-5}$ and that the number of g modes that exist inside the Sun is $\sim 10^5$)

Vielen Dank!
Thank you very much!
ご清聴いただきありがとうございました！

おまけ

In the case of the second-order fluctuation

- What we have to compute:

$$\Delta\phi_{\text{osc}} = \int \phi' \rho_0 dV + \int \phi_0 \rho' dV + \int \phi' \rho' dV$$

Eulerian perturbation:

$$q' = q(\mathbf{r}, t) - q_0(\mathbf{r}, t)$$

- But we have a few differences as below:

- Lagrangian perturbation $\delta\phi$ is related to Eulerian perturbation ϕ' to second-order,

$$\text{namely, } \delta\phi = \phi' + (\boldsymbol{\xi} \cdot \nabla)\phi_0 + \frac{1}{2}\boldsymbol{\xi} \cdot \{H(\phi_0)\boldsymbol{\xi}\}$$

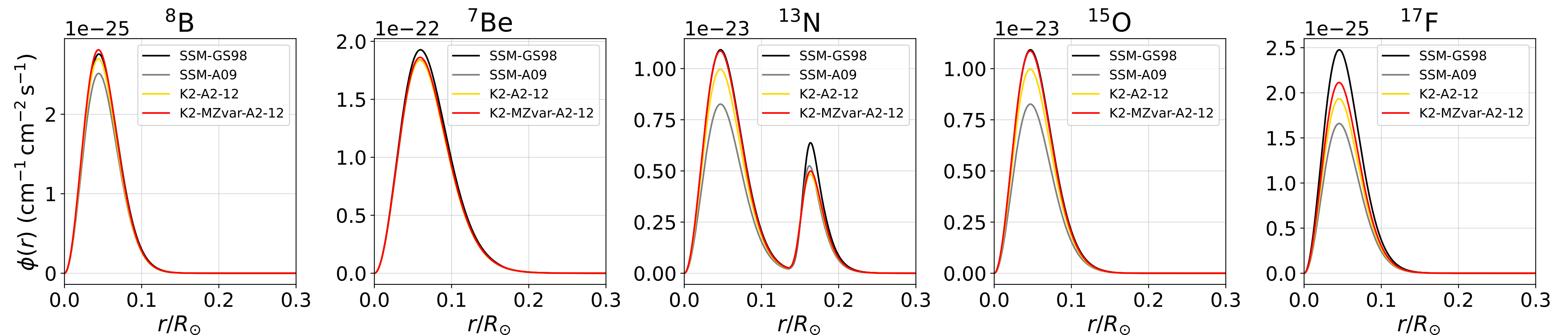
- we take into account the second-order terms in $\delta\phi$:

$$\delta \ln \phi = c_1(\delta \ln T) + c_2(\delta \ln T)^2$$

- and, we have to consider couplings among modes with different mode indices

Evaluation with solar models

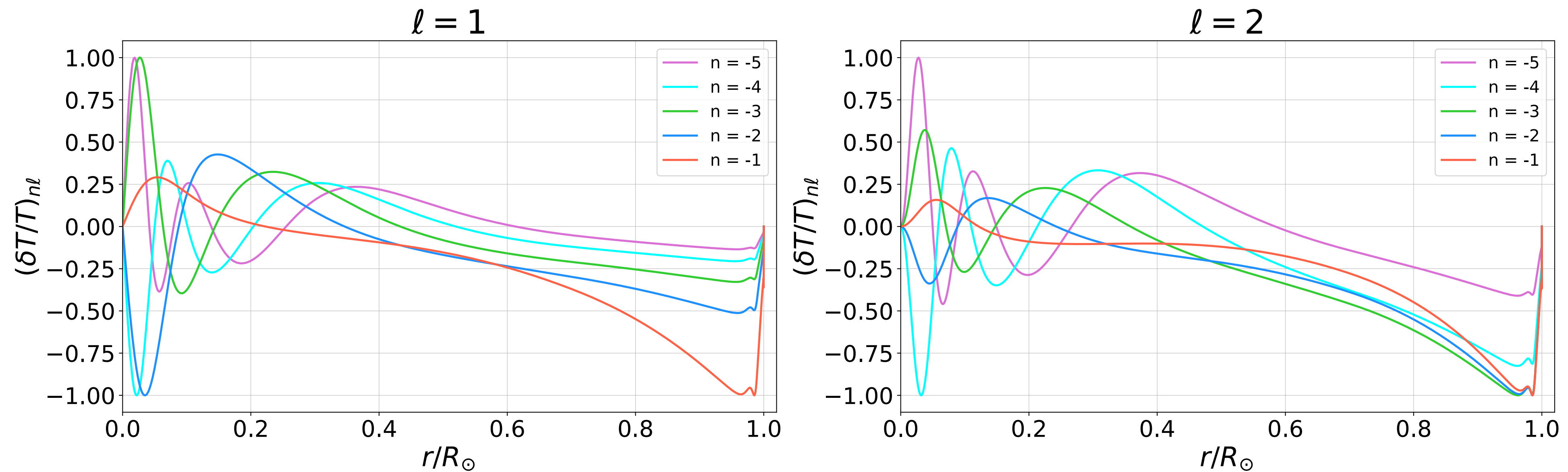
- Four 1-dimensional solar equilibrium models (Kunitomo+2022):
 - SSM-GS98: standard solar model (SSM) with the GS98 composition (high- $Z \sim 0.02$)
 - SSM-A09: SSM with the A09 composition (low- $Z \sim 0.01$)
 - K2-A2-12: a model where accretion in pre-main-sequence evolution is taken into account
 - K2-MZvar-A2-12: the same as K2-A2-12 except that chemical composition of accreted material is not constant



Flux densities inside the Sun for different neutrinos

Eigenfunctions

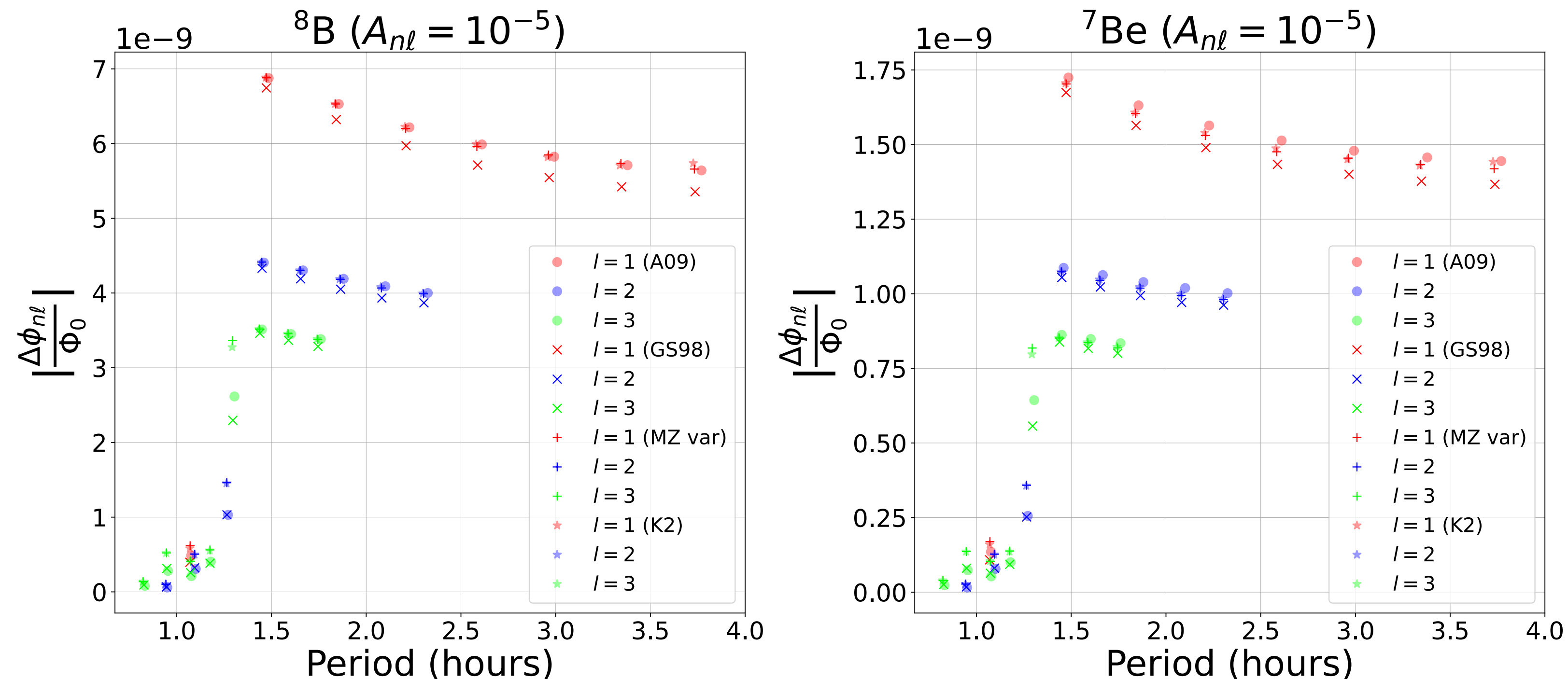
- For each solar model, we compute linear adiabatic oscillations via GYRE (Townsend+2013)
- Eigenfunctions are normalized so that the absolute value of maximum $(\delta \ln T)_{n\ell}$ is unity
- We will multiply them later by $A_{n\ell}$ to (somehow) quantify the g-mode amplitudes



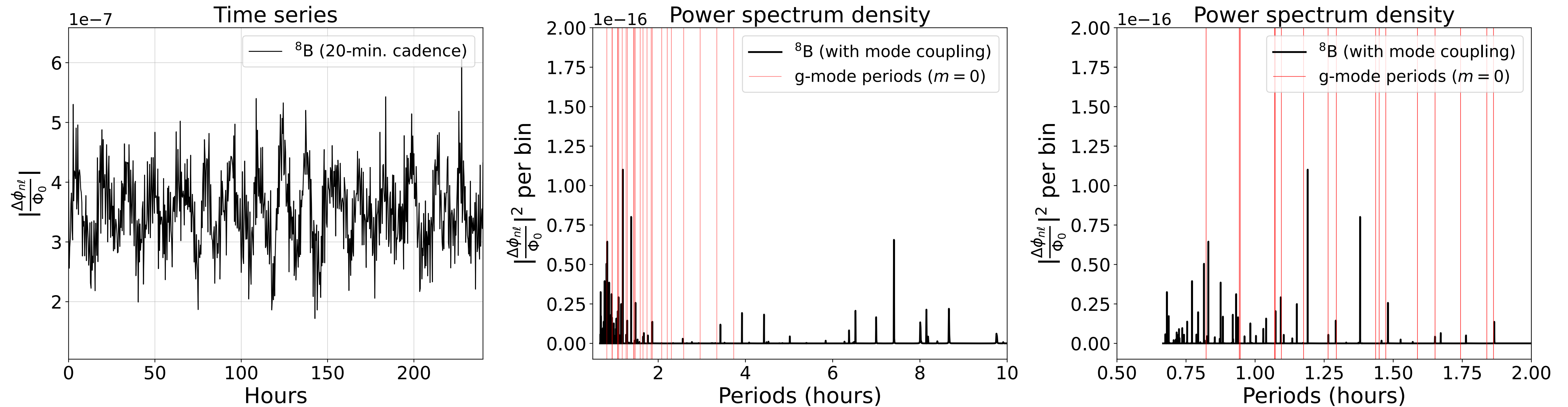
Temperature eigenfunctions $\delta T_{n\ell}(r)$ for the model K2-MZvar-A2-12

Fluctuation evaluation in the case of a single g mode

- In the case of a single g mode: $\Delta\phi_{\text{osc}} = \pi \int_0^{R_\odot} Q_{(n\ell),(n\ell)} dr + \pi \int_0^{R_\odot} Q_{(n\ell),(n\ell)} dr \times \cos(2o)$.
- $A_{n\ell} = 10^{-5}$ is assumed
- Neutrino flux fluctuations caused by this g mode are of the order of 10^{-9} in relative difference



Fluctuation evaluation in the case of a few dozens of g modes



If $A_{n\ell}$ changes in accordance with the solar cycle...

- In our analysis, we assume that $A_{n\ell}$ is a constant for all the g modes
 - What determines $A_{n\ell}$, the g-mode amplitude, is turbulent convection (e.g. Belkacem+2022)
- The constancy in $A_{n\ell}$ guarantees the constancy of the non-time-varying component, but...
- If $A_{n\ell}$ changes with some periodicities, **the non-time-varying component could vary** with the same periods
 - One promising candidate for such periodicities is the 11-year solar cycle (??)
 - Magnetic fields in the convection zone changes $\rightarrow$ convection changes $\rightarrow$ g-mode amplitudes ($A_{n\ell}$) change $\rightarrow$ non-time-varying component in the neutrino flux fluctuation changes
- Such long-period (~ 11 years) variations in the solar neutrino flux are not confirmed yet...