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3D hydrodynamics simulations of massive stars with the PROMPI code

Federico Rizzuti (Keele University)

in collaboration with: C. Georgy, R. Hirschi, D. Arnett, E. Kaiser, C. Meakin, M. Mocak, L. Scott

Outline

- Modelling stars: 1D vs 3D
- Setup choices and the PROMPI code
- A new set of Ne-shell simulations
- Kinematics and abundance profiles
- Evolution of the convective boundaries
- From 3D to 1D : parametrizing the entrainment law
- Conclusions

Stellar evolution models: the limits of 1D

Advantages:

- can model the full star
- can cover the entire lifetime
- easily compared to obsevations
- can explore mass and metallicity
- used for progenitor models

What's missing?

 self-consistent physical descriptions of mass loss, convection, rotation, magnetic fields, opacity, binarity (and their interplay)

Disadvantages:

- spherical symmetry assumed
- parametrized physics for multi-D processes
- cannot model turbulence



3D hydrodynamics models

Modelling a 3D box enclosed in / enclosing a star

Advantages:

- deviations from spherical symmetry
- can model fluid instabilities
- can include naturally 3D processes (convection, turbulence) without assuming any prescription

Disadvantages:

- high computational cost
- limited by fluid dynamical timescales
- cannot simulate full star or entire lifetime
- more difficult to compare results to observations



Why employ hydrodynamics models?

Multi-D processes can be reproduced:

- Convection, rotation, magnetic fields
- No need to assume prescriptions as in 1D e.g. mixing length theory (MLT), convective boundary mixing (CBM)
 - \rightarrow possible to use 3D data to constrain 1D parametrization
- Turbulent mixing leading to convection
- Turbulent entrainment at convective boundaries
- Internal gravity waves

321D: the link between 1D and multi-D



Possible choices for a setup

First of all, the physics of the problem:

→ stellar mass, age, core or burning layers...

Then:

- Initial conditions from a 1D stellar evolution model
- Problem geometry and resolution: plane-parallel, spherical...
 - \rightarrow be careful with singularities
- Boundary conditions: periodic, reflective...
- Gravity: constant, polynomial...



Muller (2020)

The PROMPI 3D Hydrodynamics Code

PROMPI solves the Euler equations (inviscid approximation) by:

 $\frac{\partial \rho}{\partial t} + \boldsymbol{\nabla} \cdot (\rho \, \boldsymbol{v}) = 0;$

$$\rho \, \frac{\partial \boldsymbol{v}}{\partial t} + \rho \, \boldsymbol{v} \cdot \boldsymbol{\nabla} \boldsymbol{v} = -\boldsymbol{\nabla} p + \rho \, \mathbf{g};$$

$$\rho \, \frac{\partial E_{\rm t}}{\partial t} + \rho \, \boldsymbol{v} \cdot \boldsymbol{\nabla} E_{\rm t} + \boldsymbol{\nabla} \cdot (p \, \boldsymbol{v}) = \rho \, \boldsymbol{v} \cdot \mathbf{g} + \rho (\epsilon_{\rm nuc} + \epsilon_{\nu});$$

 $\rho \, \frac{\partial X_{\mathrm{i}}}{\partial t} + \rho \, \boldsymbol{v} \cdot \boldsymbol{\nabla} X_{\mathrm{i}} = R_{\mathrm{i}},$

PROMPI : Meakin, Arnett+ 2007-onwards PROMETHEUS : Fryxell, Mueller, Arnett 1989 PPM method : Colella & Woodward 1984

• PROmetheus MPI

- finite-volume, time explicit, Eulerian, PPM implementation
- domain decomposition for parallel computing (MPI)
- Cartesian, spherical or cylindrical geometry
- reflective or periodic boundary conditions, velocity damping

Simulations of a neon-burning shell

Modelling a 3D cell in the Ne-shell of 15 M_{\odot} star with PROMPI:

- Plane-parallel "box-in-a-star" of $(0.64 \times 10^8 \text{ cm})^3$
- Multiple simulations with different resolutions (mesh size) and nuclear energy generation rates ("boosting factors")

We focus on reproducing/studying:

- Turbulent convection
- Convective boundary mixing
- Turbulent entrainment

| resolution → boosting ↓ | 128 ³ | 256 ³ | 512 ³ | 1024 ³ |
|----------------------------------|------------------|------------------|------------------|--------------------------|
| x 1 | | | Ex1 | |
| x 10 | lrez | mrez | Ex10 | vhrez |
| x 100 | | | Ex100 | |
| x 1000 | | | Ex1000 | |

Kinematic study: velocity movie



Chemical study: abundance movies



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Abundance profiles

- We study the mass fraction of the most abundant isotopes
- The profiles are used to define the convective boundaries



3.8

3.9

4.0

Evolution of the boundary locations



- The location of the boundaries moves with time, because the convective zone is growing: this is entrainment
- The entrainment rate is strongly dependent on the boosting factor
- The entrainment rate is not dependent on the resolution

Computing the entrainment law



$$E = \frac{v_{\rm e}}{v_{\rm rms}} = A \cdot Ri_B^{-n}$$

- Entrainment rate can be parametrized with a simple law using the "bulk Richardson number", representing the "stiffness" of the boundary
- Then the law can be used to improve convection in 1D models
 1D → 3D → 1D

Conclusions

- 3D hydrodynamics codes like PROMPI reproduce turbulent flow for short timescales but with great accuracy
- The interaction between nuclear burning and turbulent flow can be studied in unprecedented detail
- We completed the first detailed 3D simulations of the Ne-shell: different resolutions and luminosity boosting factors
- The entrainment rate is correlated to the boosting factor and can be parametrized with a law useful for 1D

For the future: • build a library of burning shells with the PROMPI code (C- and O-shells already present in the literature)

• complete the loop $1D \rightarrow 3D \rightarrow 1D$ and continue $\rightarrow ...$