The formation of super-massive black holes and on the (probable?) very rapid initial metal & r-process enrichment

2nd lecture

17th Russbach School 14-19 March 2022

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For recent progress on cosmology and the existence of dark matter particles (for those interested)

"Dark Matter Kindergarten Stuff"

https://www.youtube.com/watch?v=O5MbOCgCMDs

Comment

This project was published *prior* to the black-hole Nobel Prize of 2020.

That is, there is no influence / curruption whatsoever of this project by that event.

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Monthly Notices

ROYAL ASTRONOMICAL SOCIETY

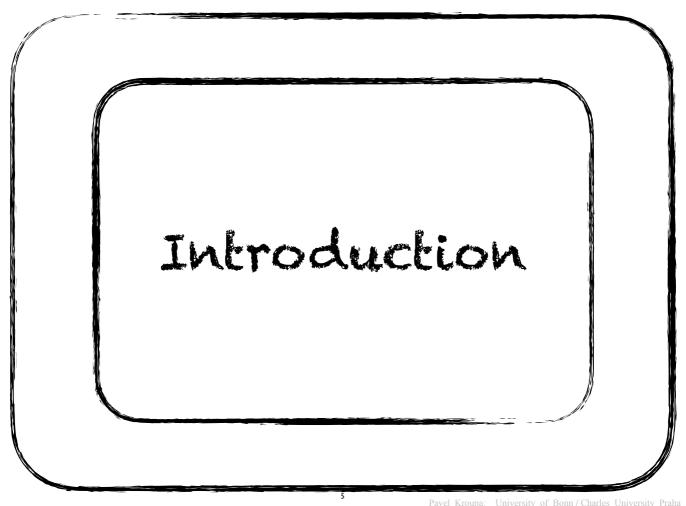
MNRAS **498**, 5652–5683 (2020) Advance Access publication 2020 August 5 doi:10.1093/mnras/staa2276

Very high redshift quasars and the rapid emergence of supermassive black holes

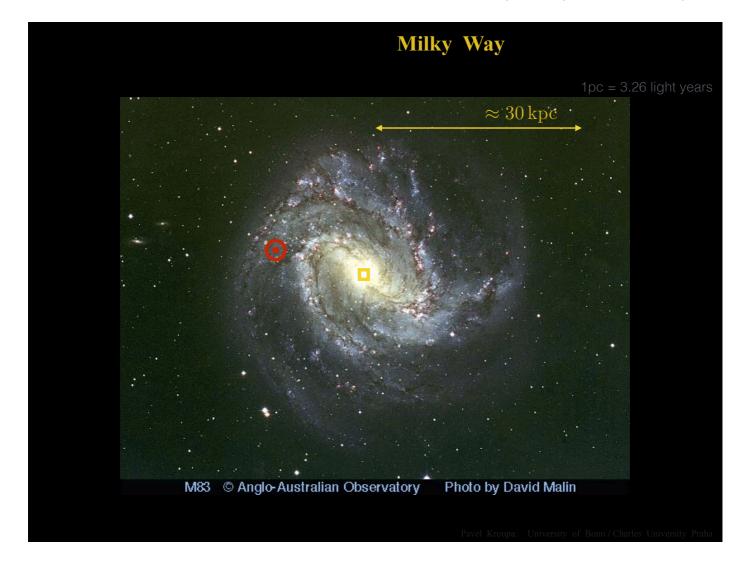
Pavel Kroupa ^⑤, ^{1,2,3}* Ladislav Subr, ² Tereza Jerabkova ^⑥1,2,3,4,5,6</sup>† and Long Wang ^⑥7

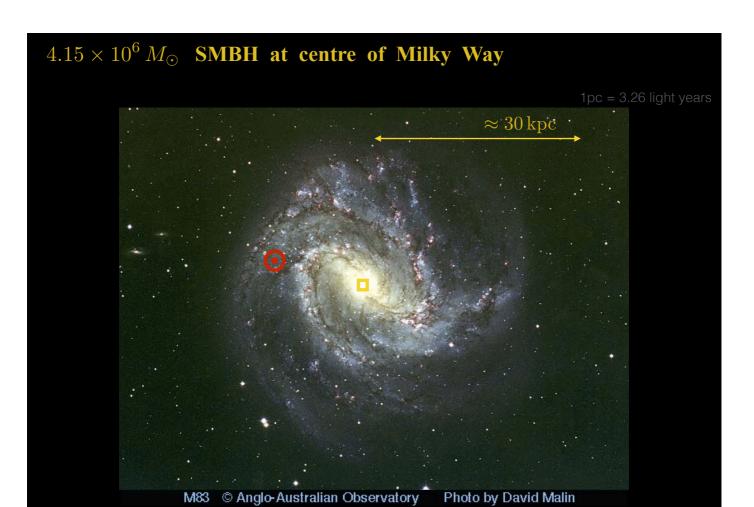
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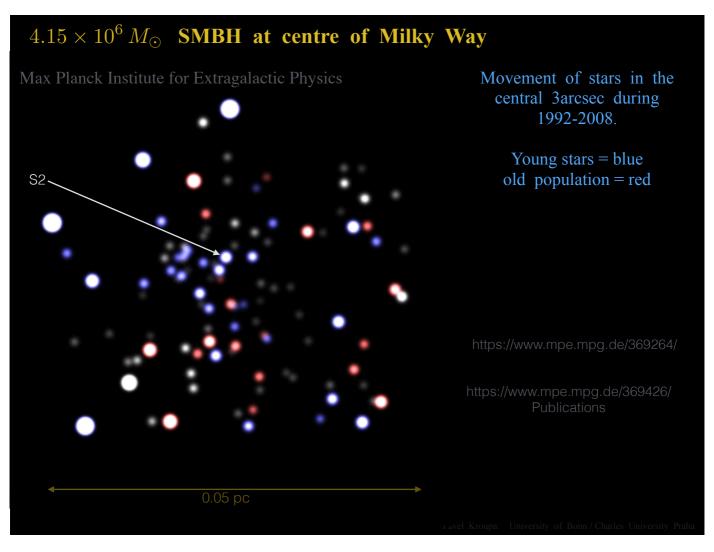


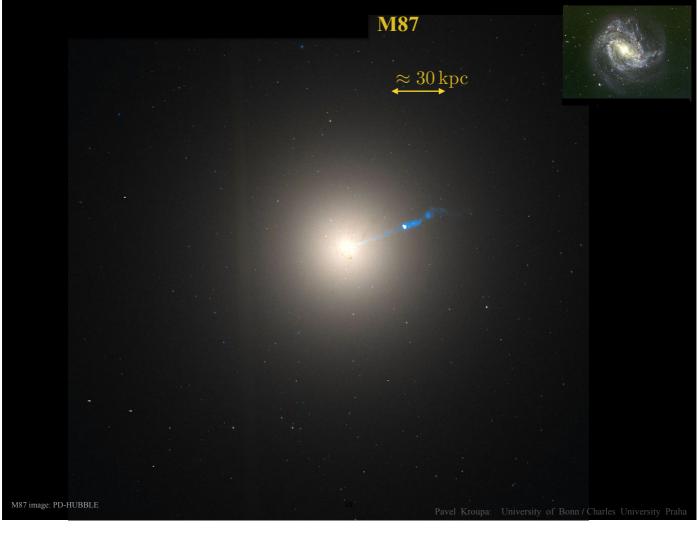
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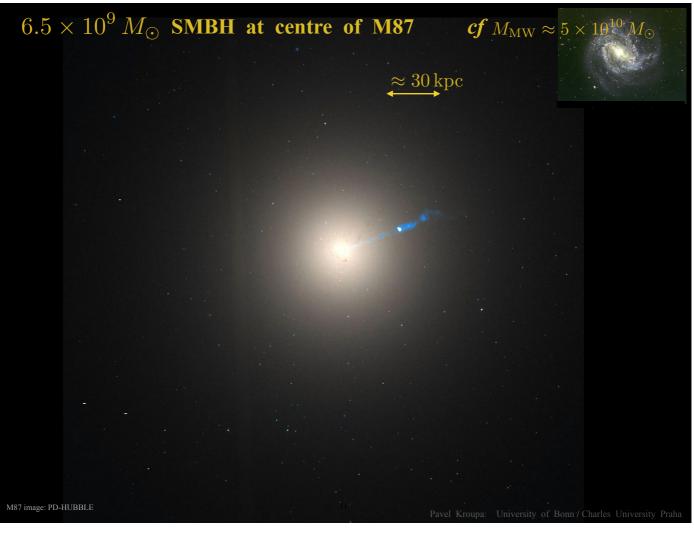


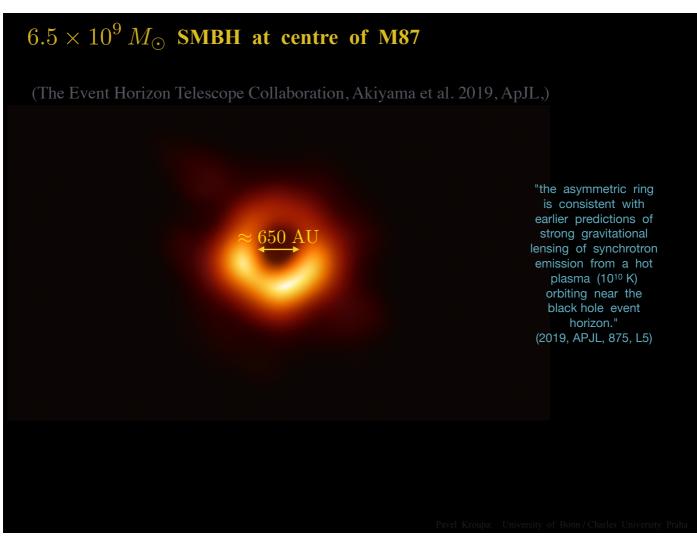


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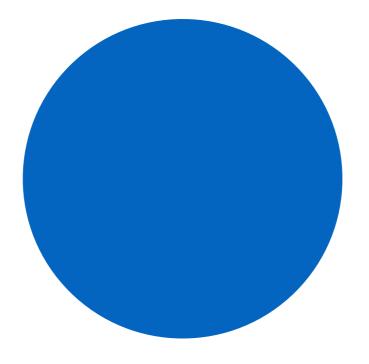








Why "shadow"?



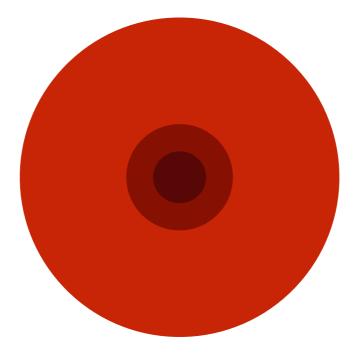
Extremely hot (blue emitting) accretion disk

The SMBH in it redshifts the light

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Why "shadow"?

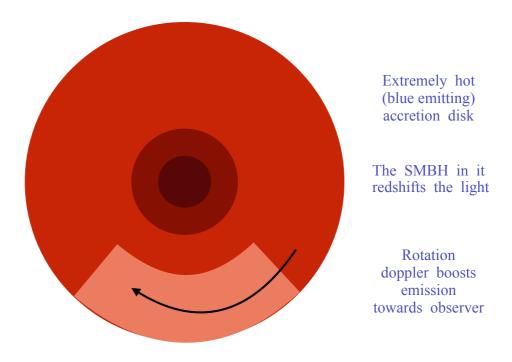


Extremely hot (blue emitting) accretion disk

The SMBH in it redshifts the light

Rotation doppler boosts emission towards observer (here to the right)

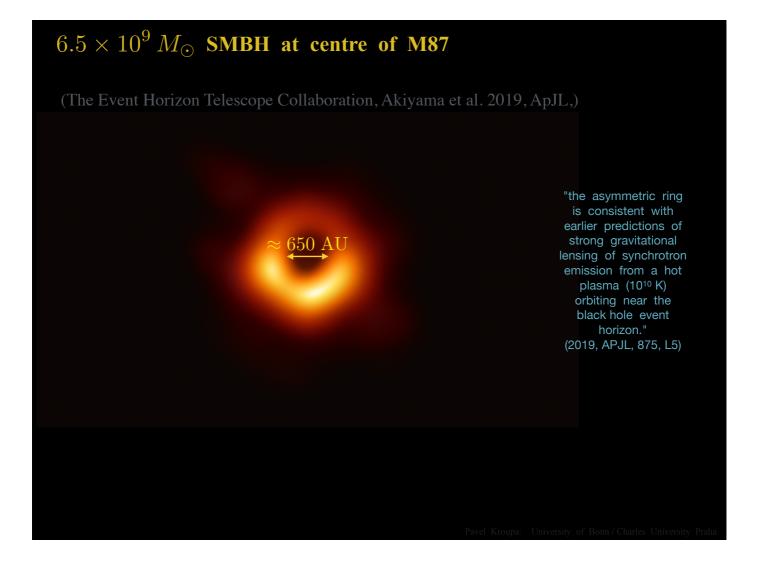
Why "shadow"?

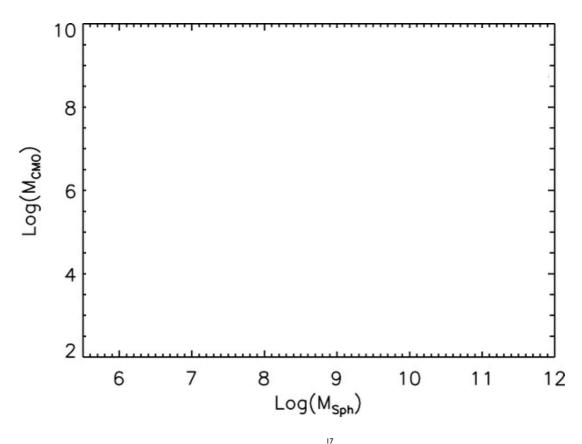


(The Event Horizon Telescope Collaboration, Akiyama et al. 2019, ApJL 875)

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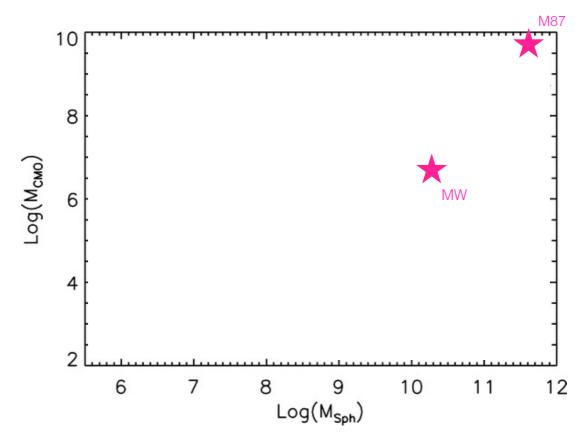
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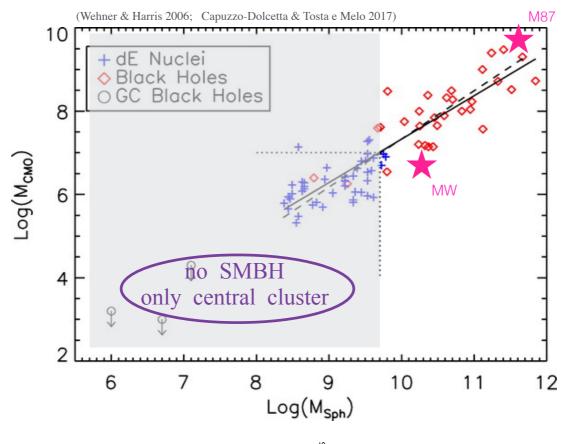


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SMBH mass vs host galaxy mass (spheroid component: elliptical galaxy = classical bulge) (Gadotti & Kaufmann 2009)



SMBH mass vs host galaxy mass (spheroid component: elliptical galaxy = classical bulge) (Gadotti & Kaufmann 2009)



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Spatial scales

(measurement of SMBH masses see review Peterson, B.M., 2014, Space Science Reviews)

$$R_{
m horizon} \equiv rac{2 \, G \, M_{
m SMBH}}{c^2}$$

event horizon = Schwarzschild radius (exact solution of GR for spherical non-rotating mass)

$$M_{\mathrm{SMBH}} = 5.15 \times 10^6 \, M_{\odot}$$
 $R_{\mathrm{horizon}} = 0.106 \, \mathrm{AU}$
 $M_{\mathrm{SMBH}} = 5.15 \times 10^9 \, M_{\odot}$ $R_{\mathrm{horizon}} = 106 \, \mathrm{AU}$

$$R_{\rm BHinfl} \equiv \frac{G M_{\rm SMBH}}{\sigma_*^2}$$

black-hole influence radius (σ_* stellar velocity disp. of stars in bulge)

e.g. globular cluster:
$$\sigma_*=10~{\rm km/s}$$
, $M_{\rm BH}=10^4~M_{\odot}$
$$R_{\rm BHinfl}=0.45~{\rm pc}$$

SMBHs appear at very high redshift

THE ASTROPHYSICAL JOURNAL, 833:222 (17pp), 2016 December 20

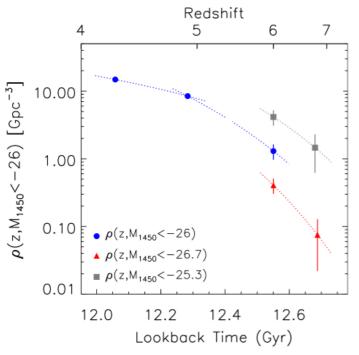


Figure 10. Density evolution of luminous quasars at z > 4. The blue circles, red triangles, and gray squares are the cumulative densities down to $M_{1450} = -26.0$, -26.7, and -25.3 mag, respectively. The data points at $z \sim 4$ and 5 are from McGreer et al. (2013), the data points at z = 6 are calculated from our new sample. The gray data point at $z \sim 7$ is taken from Venemans et al. (2013), and the red data point at $z \sim 7$ is estimated from two z > 6.5 quasars in the UKIDSS area (see details in Section 5.1). The dotted lines are the power-law (Equation (4)) fits to the data points. The figure shows a rapid decline of the quasar spatial density from $z \sim 5$ toward higher redshifts.

(Jiang, McGreer et al. 2016)

13.6 13.8 first stars thought to form

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... and at even higher redshift

THE ASTROPHYSICAL JOURNAL LETTERS

 $pprox 700~{
m Myr}$ $1.5 imes 10^9~{
m M}_{\odot}$ Pōniuā'ena: A Luminous z = 7.5 Quasar Hosting a 1.5 Billion Solar Mass Black Hole

Jinyi Yang¹ D, Feige Wang^{1,13} D, Xiaohui Fan¹ D, Joseph F. Hennawi^{2,3} D, Frederick B. Davies^{2,4} D, Minghao Yue¹, Eduardo Banados³, Xue-Bing Wu^{5,6}, Bram Venemans³, Aaron J. Barth⁷ + Show full author list

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The Astrophysical Journal Letters, Volume 897, Number 1

Citation Jinyi Yang et al 2020 ApJL 897 L14

Need $10^4~M_{\odot}$ seed to grow continuously at superEddington rate since z=30(100Myr after Big Bang).

+ Article information

Abstract

We report the discovery of a luminous quasar, J1007+2115 at z = 7.515 ("Pōniuā'ena"), from our wide-field reionization-era quasar survey. J1007+2115 is the second quasar now known at z > 7.5, deep into the reionization epoch. The quasar is powered by a $(1.5 \pm 0.2) \times 10^9 M_{\odot}$ supermassive black hole (SMBH), based on its broad Mg II emission-line profile from Gemini and Keck near-IR spectroscopy. The SMBH in J1007+2115 is twice as massive as that in quasar J1342+0928 at z = 7.54, the current quasar redshift record holder. The existence of such a massive SMBH just 700 million years after the Big Bang

The Problem(s)

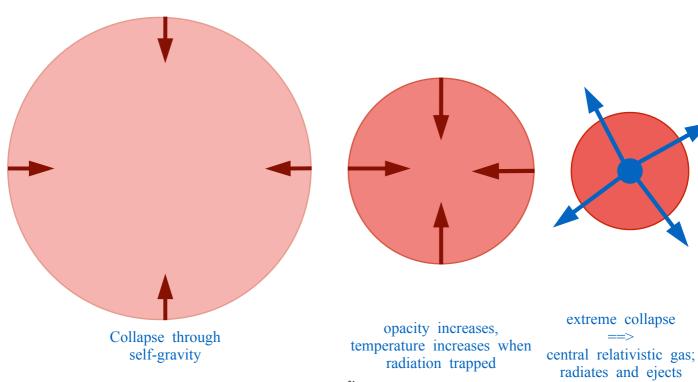
Three problems:

- I) How can SMBHs even form ? (need to squeeze $> 10^6 M_{\odot}$ of matter into extremely small space)
- II) Why do SMBH masses correlate with the masses of their hosting galaxies?
- III) How can SMBHs form within a few 100 Myr after the Big Bang?

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Why can a SMBH seed not form from direct collapse of a post-Big-Bang gas cloud?



This leads to self-radiation-limited growth

$$L_{\rm acc} = \epsilon_{\rm r} \, \dot{M} \, c^2$$

 $\epsilon_{\rm r} \approx 0.1-0.5$

fraction of rest-mass radiated

plus radiation pressure leads to accretion shutoff — difficult to form BH

$$M_{\mathrm{BH}}(t) = M_{\mathrm{BH,0}} e^{\frac{f_{\mathrm{Edd}}(t-t_0)}{\tau_{\mathrm{Edd}}}}$$

$$f_{\mathrm{Edd}} = \frac{1 - \epsilon_{\mathrm{r}}}{\epsilon_{\mathrm{r}}}$$

$$\tau_{\mathrm{Edd}} = \frac{c \, \sigma_{\mathrm{T}}}{4\pi \, G \, m_{\mathrm{p}}} \approx 452 \, \mathrm{Myr}$$

c speed of light

 $\sigma_{\rm T}$ Thompson scattering cross section

 $m_{\rm P}$ proton mass

Thus, to grow a

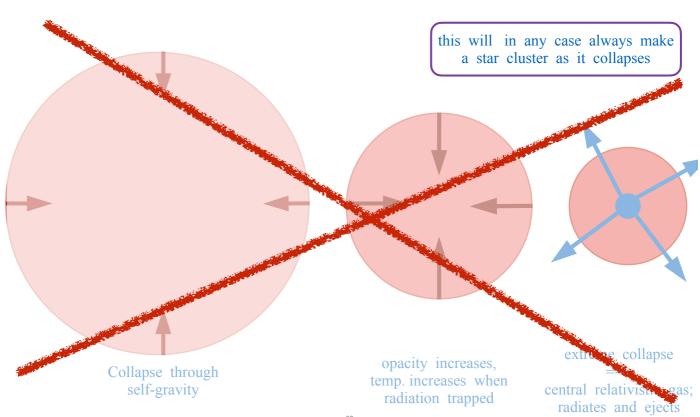
 $M_{\rm BH,0}=10\,M_{\odot}$ seed to

 $M_{\rm SMBH} = 10^{10} \, M_{\odot}$ need sustained continuous accretion for $\approx 1 \, {\rm Gyr}$ at the superEddington rate ($\epsilon_{\rm r} = 0.1$).

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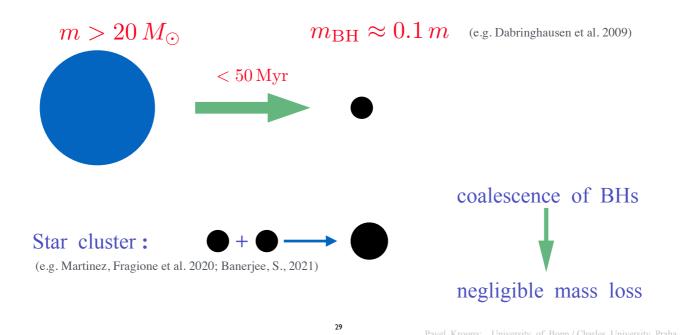
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Why not form an SMBH seed from direct collapse of post-Big-Bang gas cloud?



Need to efficiently compactify matter to a BH seed

Start with first massive stars:



coalescence of BHs ===> negligible mass loss:

coalescence of BHs ===> negligible mass loss:

Table III in Abbott et al. 2019; LIGO Scientific Collaboration and Virgo Collaboration)

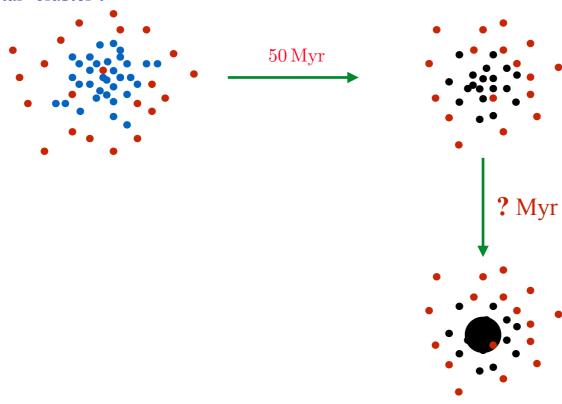
Event	m_1/M	m_2/M_0	$_{\odot}$ \mathcal{M}/M_{\odot}	$\chi_{ m eff}$	M_f/M_{\odot}	a_f	$E_{\rm rad}/(M_{\odot}c^2$) <i>e</i> _r
GW150914	35.6+4	$^{7}_{1}$ 30.6 $^{+3}_{-4}$	$^{0}_{4}$ 28.6 $^{+1.7}_{-1.5}$	$-0.01^{+0.12}_{-0.13}$	$63.1_{-3.0}^{+3.4}$	$0.69^{+0.05}_{-0.04}$	$3.1^{+0.4}_{-0.4}$	3.
GW151012	$23.2^{+14}_{-5.5}$	$^{1.9}_{5}$ 13.6 $^{+4}_{-4}$	$^{1}_{8}$ 15.2 $^{+2.1}_{-1.2}$	$0.05^{+0.31}_{-0.20}$	$35.6^{+10.8}_{-3.8}$	$0.67^{+0.13}_{-0.11}$	$1.6^{+0.6}_{-0.5}$	3.
GW151226	13.7^{+8}_{-3}	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$8.9^{+0.3}_{-0.3}$	$0.18^{+0.20}_{-0.12}$	$20.5^{+6.4}_{-1.5}$	$0.74^{+0.07}_{-0.05}$	$1.0^{+0.1}_{-0.2}$	3.
GW170104	30.8^{+7}_{-5}	$^{3}_{6}$ 20.0 $^{+4}_{-4}$	$^{9}_{6}$ 21.4 $^{+2.2}_{-1.8}$	$-0.04^{+0.17}_{-0.21}$	$48.9^{+5.1}_{-4.0}$	$0.66^{+0.08}_{-0.11}$	$2.2^{+0.5}_{-0.5}$	3.
GW170608	11.0^{+5}_{-1}	$7.6^{+1.4}_{-2.2}$	$7.9^{+0.2}_{-0.2}$	$0.03^{+0.19}_{-0.07}$	$17.8^{+3.4}_{-0.7}$	$0.69^{+0.04}_{-0.04}$	$0.9^{+0.0}_{-0.1}$	3.
GW170729	50.2^{+16}_{-10}	$\frac{0.2}{0.2}$ 34.0 $^{+9}_{-10}$	$^{1}_{0.1}$ 35.4 $^{+6.5}_{-4.8}$	$0.37^{+0.21}_{-0.25}$	79.5 ^{+14.7} _{-10.2}	$0.81^{+0.07}_{-0.13}$	$4.8^{+1.7}_{-1.7}$	4.
GW170809	35.0^{+8}_{-5}	$^{3}_{9}$ 23.8 $^{+5}_{-5}$	$^{1}_{2}$ 24.9 $^{+2.1}_{-1.7}$	$0.08^{+0.17}_{-0.17}$	56.3 ^{+5.2} _{-3.8}	$0.70^{+0.08}_{-0.09}$	$2.7^{+0.6}_{-0.6}$	3.
GW170814	30.6^{+5}_{-3}	$^{6}_{0}$ 25.2 $^{+2}_{-4}$	$^{8}_{0}$ 24.1 $^{+1.4}_{-1.1}$	$0.07^{+0.12}_{-0.12}$	$53.2^{+3.2}_{-2.4}$	$0.72^{+0.07}_{-0.05}$	$2.7^{+0.4}_{-0.3}$	3.
GW170817	$1.46^{+0.}_{-0.}$	$^{12}_{10}$ $1.27^{+0.0}_{-0.0}$	$^{09}_{09} 1.186^{+0.001}_{-0.001}$	$0.00^{+0.02}_{-0.01}$	≤ 2.8	≤ 0.89	≥ 0.04	≥
GW170818	35.4+7	$^{5}_{7}$ 26.7 $^{+4}_{-5}$	$^{3}_{2}$ $26.5^{+2.1}_{-1.7}$	$-0.09^{+0.18}_{-0.21}$	59.4 ^{+4.9} _{-3.8}	$0.67^{+0.07}_{-0.08}$	$2.7^{+0.5}_{-0.5}$	3.
GW170823	$39.5^{+11}_{-6.}$	$\frac{12}{7}$ 29.0 $^{+6}_{-7}$	$\frac{7}{8}$ 29.2 $^{+4.6}_{-3.6}$	$0.09^{+0.22}_{-0.26}$	$65.4^{+10.1}_{-7.4}$	$0.72^{+0.09}_{-0.12}$	$3.3^{+1.0}_{-0.9}$	3.

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This is promising, but how to merge $> 10^4$ BHs sufficiently quickly?

Star cluster:



This is promising, but how to merge $> 10^4$ BHs sufficiently quickly?

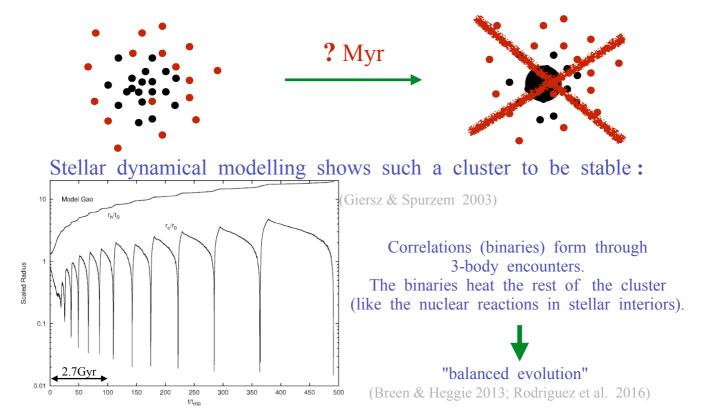
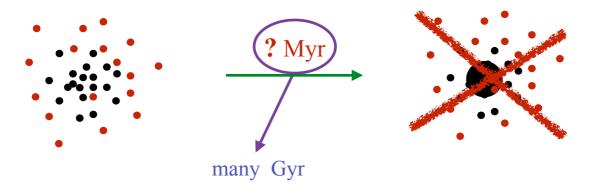


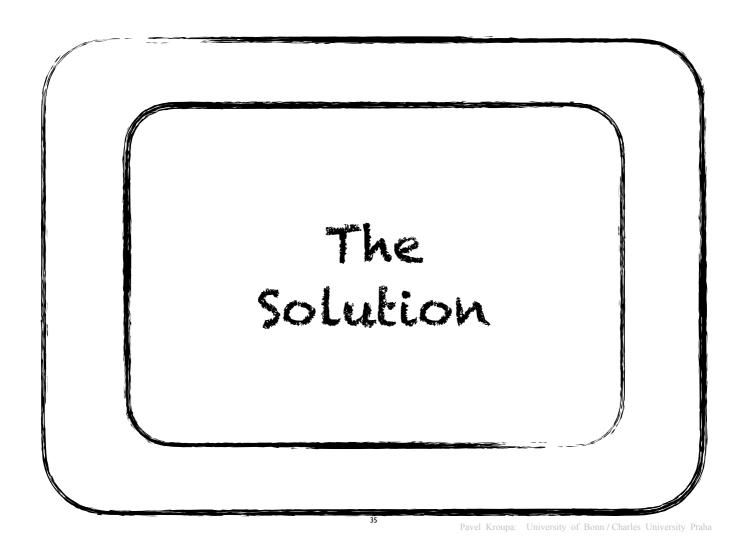
Figure 4. Evolution of the core and half-mass radii (scaled by the scale-length of the Plummer model) as a function of time (scaled by the initial half-mass relaxation time).

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This is promising, but how to merge $> 10^4$ BHs sufficiently quickly?



Clusters of BHs cannot evolve to a SMBH seed within $\approx 200 \,\mathrm{Myr}$



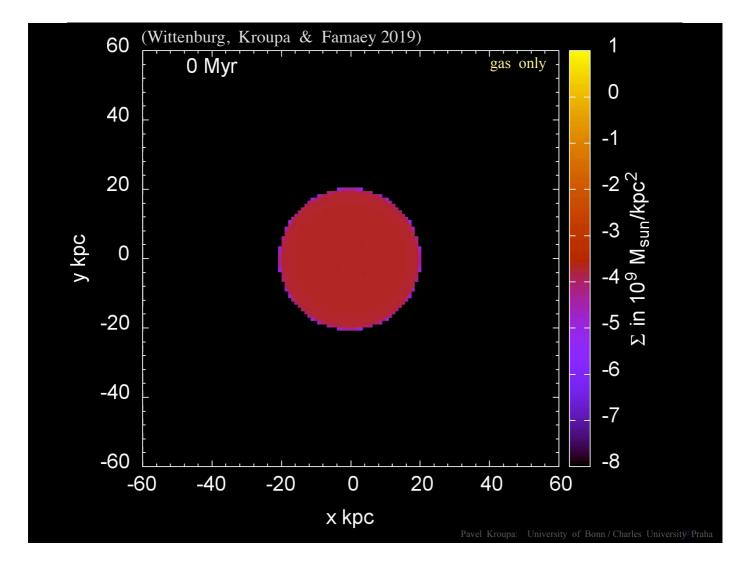
Put this into the context of galaxy formation and evolution: (Kroupa, Subr, Jerabkova, Wang 2020)

First-formed very metal-poor extremely massive star-burst star cluster forms at centre of *future* spheroid: (confirmed by Martin-Navarro et al. 2018 = Vazdekis IACanary team)

Galaxy formation and evolution: (Wittenburg, Kroupa & Famaey 2019)

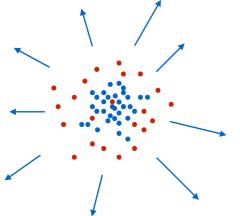
The evolution over 10 Gyr of a spherical gas cloud of mass $M_{gas} = 6.4 \times 10^9 \, M_{\odot}$ and $r_{sph} = 20 \, kpc$ and with an initial cylindrical rotational law with $\eta = 0.025 \, Myr^{-1}$

Actual E galaxy-formation simulations are underway (Eappen, Kroupa, Wittenburg & Famaey, in prep.)



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extreme radiation field and outwards wind stops gas infall from surrounding forming spheroid (Ploeckinger et al. 2019)

First-formed hyper-massive star-burst cluster:

(Jerabkova, Kroupa, Dabringhausen, Hilker & Bekki 2017)

T. Jeřábková et al.: The formation of UCDs: Testing for a variable stellar IMF

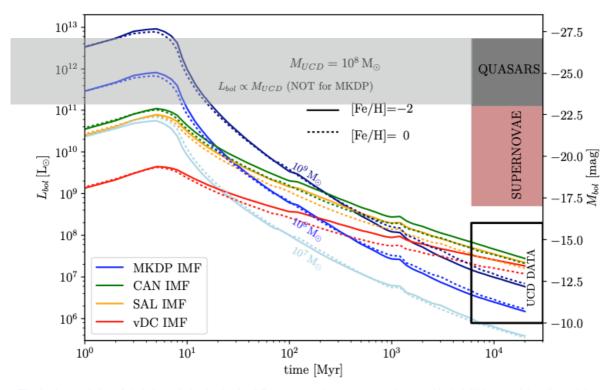


Fig. 3. Time evolution of the bolometric luminosity for different IMFs. The MKDP IMF changes with the initial mass of the UCD and does not scale linearly with its mass. In contrast, the vDC, SAL, and CAN IMFs have UCD-mass-independent slopes and scale proportionally with M_{UCD} . The grey panel shows the typical luminosities of quasars (Dunlop et al. 1993, 2003; Souchay et al. 2015) and the brown panel shows the luminosity span for the peak luminosities supernovae (Gal-Yam 2012; Lyman et al. 2016), which might cause luminosity variations in L_{bol} of UCDs younger than about 50 Myr according to SNe rates.

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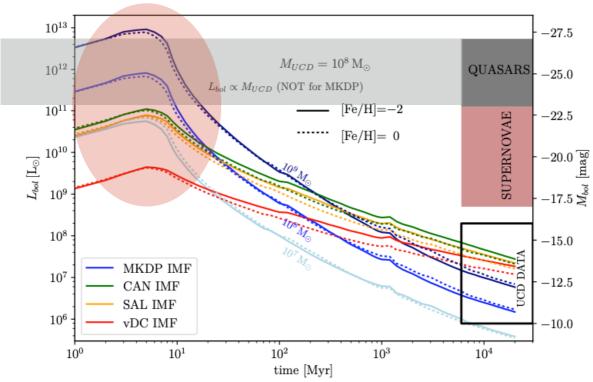
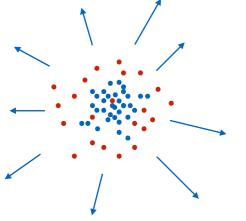


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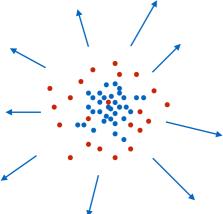
e.g. for 50 Myr, cluster of $10^8\,M_\odot$ luminosity of $10^{12}\,L_\odot$ drives metal-rich outflow at $2\,M_\odot/{\rm yr}$ at few $10^3\,{\rm km/s}$

Might well be mistaken for a z>6 quasar.

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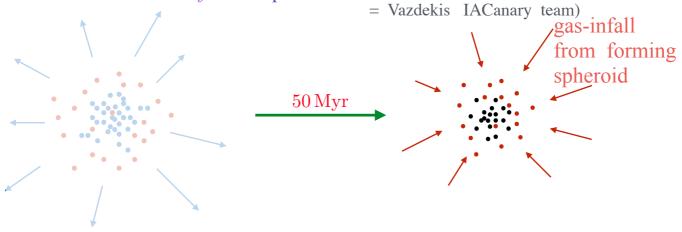
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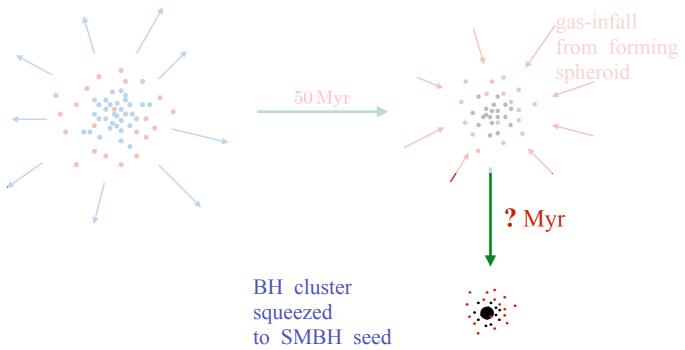
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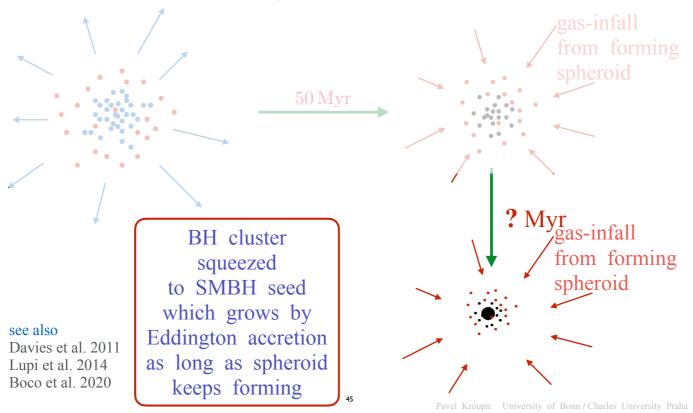
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First-formed very metal-poor extremely massive star-burst star cluster forms at centre of *future* spheroid:



Evolution of BH cluster as gas in-falls

(Kroupa, Subr, Jerabkova, Wang 2020)

$$E(t) = \frac{G (M_{\rm gas}(t) + N m_{\rm BH}(t))^2}{R(t)} = N m_{\rm BH}(t) \sigma(t)^2$$

potential energy; cluster in virial equilibrium at any time.

$$\sigma(t) = \left(\frac{G \left(M_{\text{gas}}(t) + N m_{\text{BH}}(t)\right)^2}{R N m_{\text{BH}}(t)}\right)^{\frac{1}{2}}$$

velocity dispersion of BH cluster

$$\dot{R} = \left(\frac{\dot{E}}{G} - \frac{2}{R}(M_{\text{gas}} + N \, m_{\text{BH}})(\dot{M}_{\text{gas}} + N \, \dot{m}_{\text{BH}})\right) \frac{R^2}{(M_{\text{gas}} + N \, m_{\text{BH}})^2}$$

with
$$\dot{E} = \dot{E}_{\rm diss} + \dot{E}_{\rm bin,heat}$$

binding energy change through gas drag and binary heating through BH-BH binaries.

(Kroupa, Subr, Jerabkova, Wang 2020)



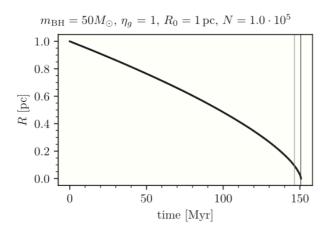


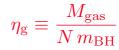
Figure 5. The evolution of the cluster radius, R(t), for $\eta_{\rm g}=1.0$, $m_{\rm BH}=50\,{\rm M}_\odot$, $M_{\rm BH,0}=5\times10^6\,{\rm M}_\odot$, and $R_0=1$ pc. The cluster of BHs contracts because gas drag outweighs binary heating. See Fig. 6 for the ratio of the rates of energy change due to gas drag and binary heating. The vertical thin grey line represents the time, $t_{\rm crit}$, when binary heating ceases to be relevant (equation 17) and the vertical thick line shows the time, $t_{\rm rel}$, when the BH cluster reaches the relativistic state (equation 20).

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Evolution of BH cluster as gas infalls

(Kroupa, Subr, Jerabkova, Wang 2020)



Quasars and SMBHs 5663

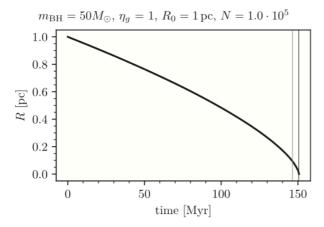


Figure 5. The evolution of the cluster radius, R(t), for $\eta_{\rm g}=1.0$, $m_{\rm BH}=50\,{\rm M}_{\odot}$, $M_{\rm BH,0}=5\times10^6\,{\rm M}_{\odot}$, and $R_0=1$ pc. The cluster of BHs contracts because gas drag outweighs binary heating. See Fig. 6 for the ratio of the rates of energy change due to gas drag and binary heating. The vertical thin grey line represents the time, $t_{\rm crit}$, when binary heating ceases to be relevant (equation 17) and the vertical thick line shows the time, $t_{\rm rel}$, when the BH cluster reaches the relativistic state (equation 20).

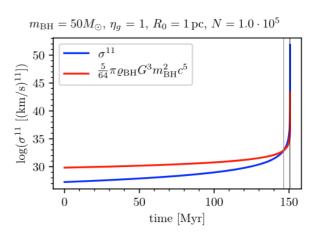
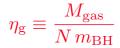
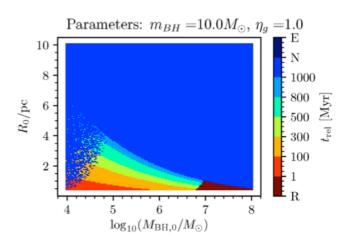


Figure 7. The evolution of the velocity dispersion (equation 24, blue line) and of the condition when binaries cease to be a heating source (equation 17, red line) for the parameters as in Fig. 5. The velocity dispersion increases as the cluster plus gas cloud shrinks. When the two lines intersect (the left vertical line), binary heating ceases to be active. The vertical lines have the same meaning as in Fig. 5. Note that $\log = \log_{10}$.

(Kroupa, Subr, Jerabkova, Wang 2020)

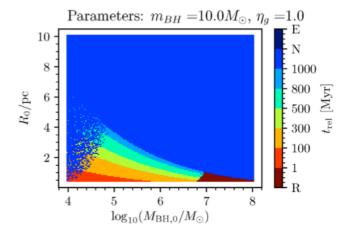


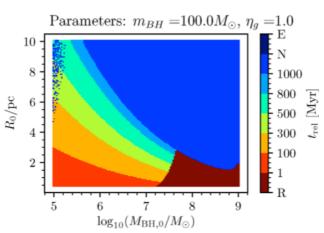


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Evolution of BH cluster as gas infalls (Kroupa, Subr, Jerabkova, Wang 2020)

$$\eta_{
m g} \equiv rac{M_{
m gas}}{N\,m_{
m BH}}$$

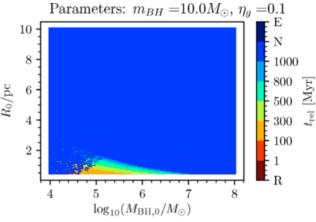


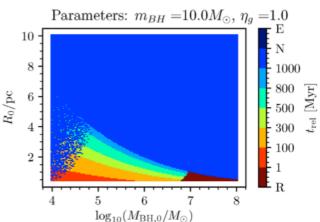


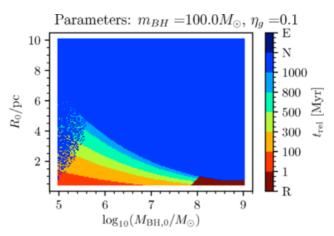
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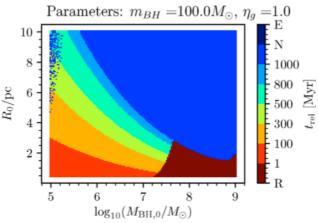
 $\eta_{
m g} \equiv rac{M_{
m gas}}{N \, m_{
m BH}}$

(Kroupa, Subr, Jerabkova, Wang 2020)









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Evolution of BH cluster as gas infalls

(Kroupa, Subr, Jerabkova, Wang 2020)

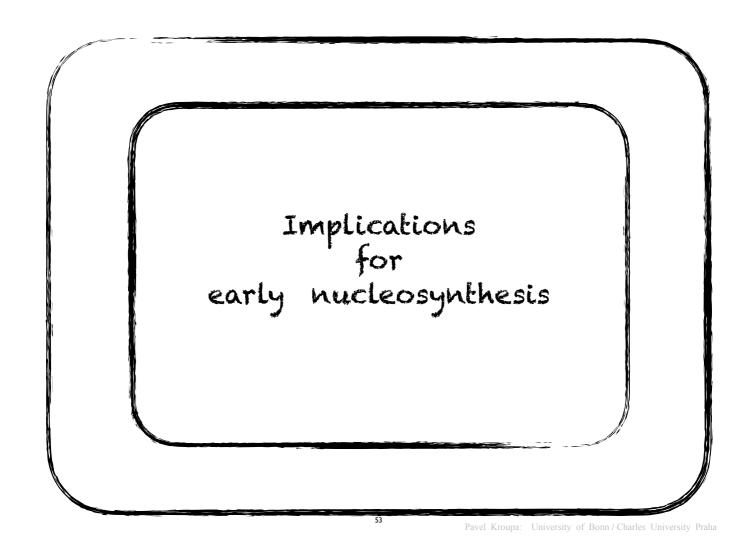


shrinkage to relativistic regime ($\sigma_{\rm rel} \equiv 3000 \ {\rm km/s}$) within 200 Myr of BH cluster formation (250 Myr after 1st star forms).

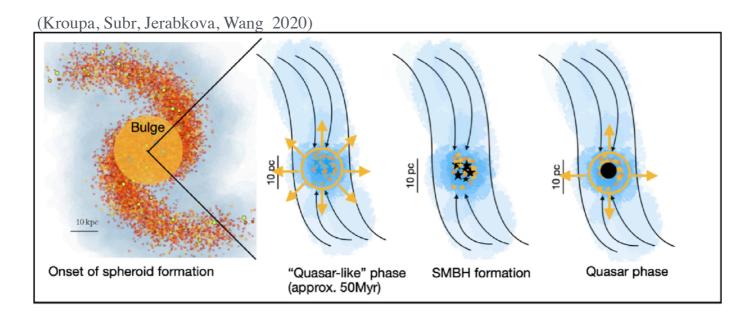
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then *collapse of >5%* of BH cluster to SMBH seed through gravitational wave energy loss essentially *instantaneous* (see also Lee 1993; Kupi et al. 2006 for Nbody).

A SMBH seed weighing >5% of the BH cluster is possible within 200-300 Myr of the formation of the first stars.



Remember/Summary



(Kroupa, Subr, Jerabkova, Wang 2020)

During the squeeze millions of neutron stars are squashed together



Strong flux of elements before any CCSN?

Strong flux of neutrons?

To intergalactic distances ?

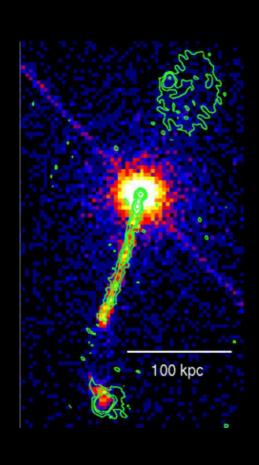
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Within 11 Mpc there are about 140 galaxies more massive than 109 Msun i.e.
140 SMBHs

First star formation about 200Myr after Big Bang
i.e. at redshift z=20
the Universe was 1/21 its present size
i.e. there were 140 forming SMBHs within 500kpc of "us"

Jets from quasars can reach lengths of 100s kpc



A Chandra X-ray image of the quasar 3C19.44. The overlayed contours show the radio emission (the dimension 100kpc corresponds to 329,000 light-years; the extremely bright core produces a line of bright pixels as an artifact).

Credit: NASA/Chandra VLA and Harris et al

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Enrichment by heavy elements

(even r-process?)

within a Myr after Big Bang

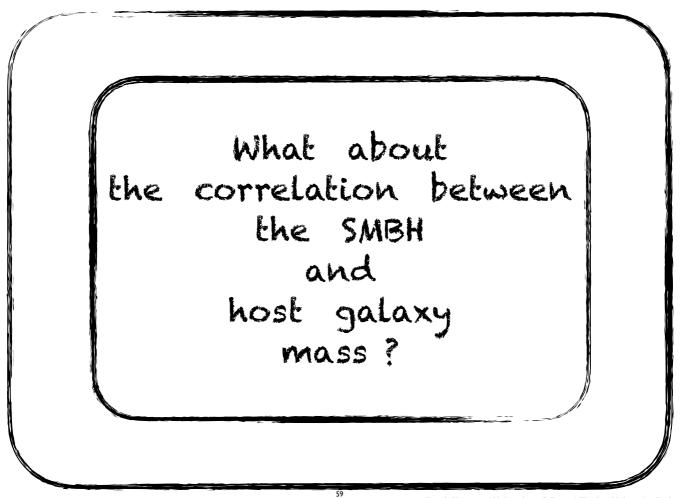
to distances of 100s of kpc,

i.e.

all of space

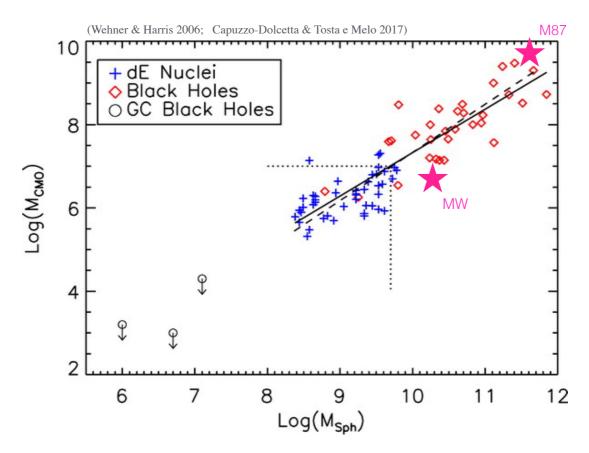
?

this is likely to be non-homogeneous i.e. more r-process "here" vs more s-process "there" etc.



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Remember: SMBH mass vs host galaxy mass (spheroid component)



Need to ensure that enough gas falls into the central region of a forming galaxy.

This is not possible in disk galaxies, they form too slowly.



EWASS 2019

Halve of all disk galaxies have bulges, and a few percent of all galaxies are huge bulges = elliptical galaxies

These galaxy spheroids form rapidly after Big Bang, vast quantities of gas collapsed rapidly.



https://ned.ipac.caltech.edu/level5/Sept14/Kormendy/Kormendy7.html

Need to calculate the most massive forming star cluster ($M_{ecl,max}$) in a galaxy with a global star-formation rate (SFR)

$$SFR = \frac{\text{mass of galaxy spheroid}}{\text{time to form galaxy}}$$

$$SFR = \frac{M_{\rm i,gal}}{\Delta \tau} \approx \frac{M_{\rm p,gal}}{\Delta \tau}$$

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What about the correlation between the SMBH and host galaxy mass?

Easy: use IGIMF theory (Kroupa & Weidner 2003; Kroupa et al. 2013; Yan, Jerabkova et al. 2017; Jerabkova, Yan et al. 2018)

The total mass in stars formed in a galaxy over time δt is $M_{\rm tot} = SFR \times \delta t$

$${\rm But} \qquad M_{\rm tot} = \int_{M_{\rm ecl,min}}^{M_{\rm ecl,max}} \, \xi_{\rm ecl}(M_{\rm ecl}) \; M_{\rm ecl} \; dM_{\rm ecl}$$

For
$$M_{
m ecl,min} = 5~M_{\odot}$$
 and with $1 = \int_{M_{
m ecl,max}}^{M_{
m ecl,max}} \xi_{
m ecl}(M_{
m ecl})~dM_{
m ecl}$

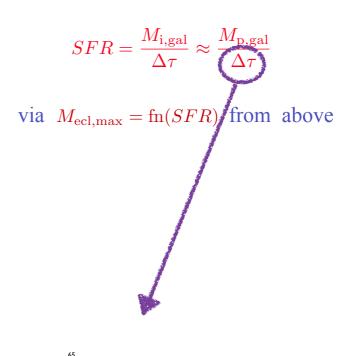
where
$$M_{\rm ecl,max*} = 10^{10} \, M_{\odot}$$

Thus
$$M_{\rm ecl,max} = {\rm fn}(SFR)$$
 as observed (Weidner et al. 2004; Randriamanakoto et al. 2013)

What is δt ?

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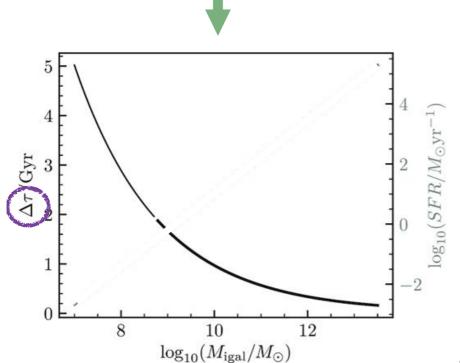
The mass of this cluster, $M_{\rm ecl,max}$, strongly correlates with later host galaxy mass, $M_{p,gal}$, through



The "downsizing" or "Thomas time" $\Delta \tau$

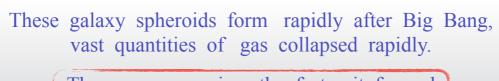


ages, Z, $\left[\frac{\alpha}{\text{Fe}}\right]$ values for 124 early-type galaxies (Thomas et al. 2005; Recchi et al. 2009)

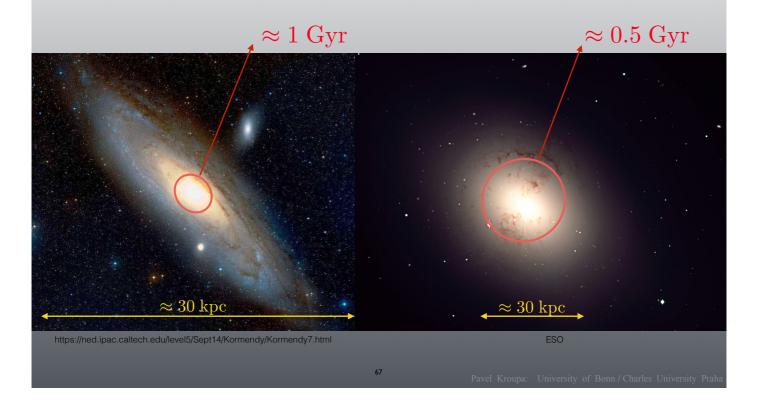


confirmed by de la Rosa et al. 2011; McDermid et al. 2015; Martin-Navarro et al. 2018; Salvador-Rusinol et al. 2020 Yan, Jerabkova et al. 2021

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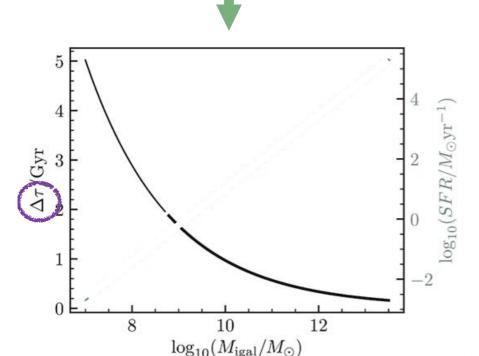
These more massive, the faster it formed = "downsizing".



The "downsizing" or "Thomas time" $\Delta \tau$



ages, Z, $\left[\frac{\alpha}{\text{Fe}}\right]$ values for 124 early-type galaxies (Thomas et al. 2005; Recchi et al. 2009)

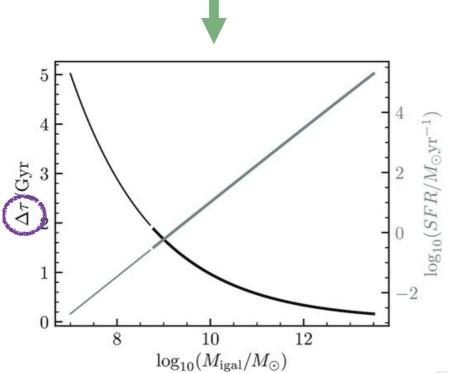


confirmed by de la Rosa et al. 2011; McDermid et al. 2015; Martin-Navarro et al. 2018; Salvador-Rusinol et al. 2020 Yan, Jerabkova et al. 2021

The "downsizing" or "Thomas time" \triangle_7



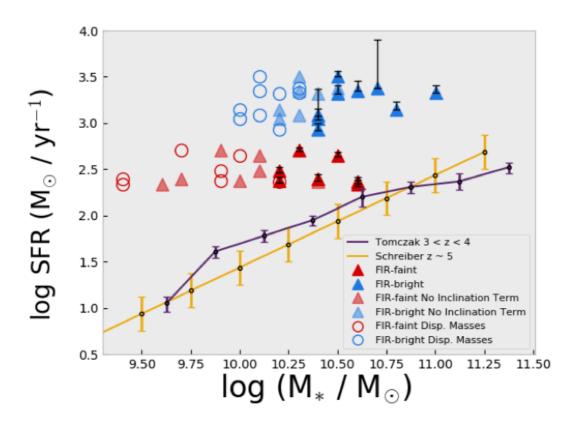
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ALMA Observations of Quasar Host Galaxies at $z \approx 4.8$ Nguyen, Lira, Trakhtenbrot et al. 2020 $t \approx 1.2 \, \mathrm{Gyr}$ after Big Bang

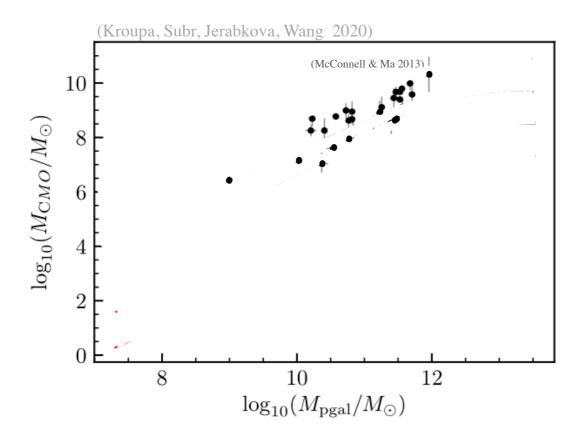


Thus we obtain

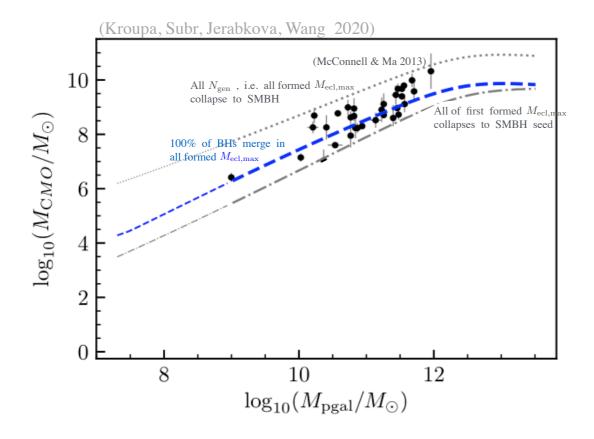
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Thus we obtain



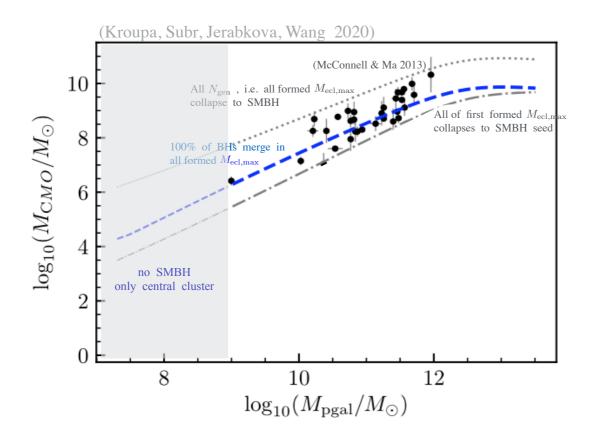
Thus we obtain



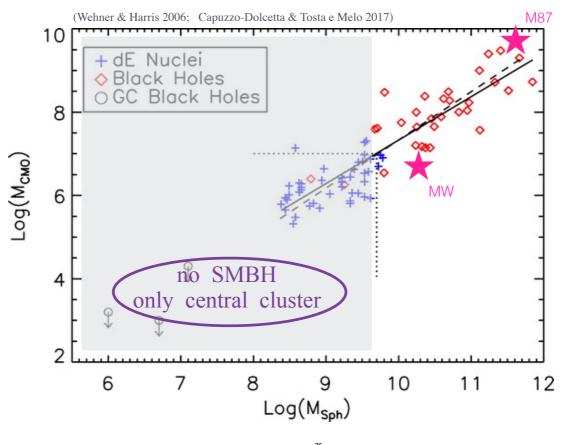
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Thus we obtain



SMBH mass vs host galaxy mass (spheroid component)

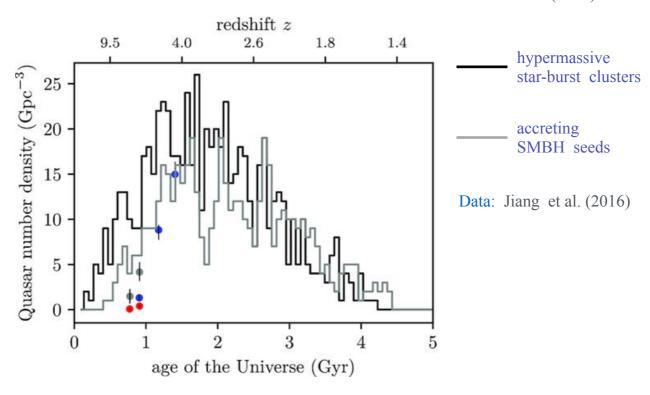


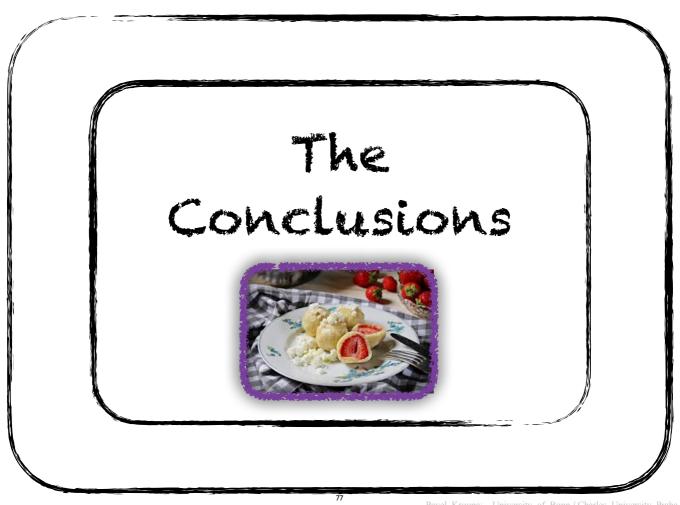
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The appearance of the first quasars

(Kroupa, Subr, Jerabkova, Wang 2020)

Based on the *empirical* galaxy-formation times scales of Thomas et al. (2005)





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Solutions to the three problems:

- I) How can SMBHs even form?

 Primarily through coalescence of $> 10^4$ stellar BHs (driven by gas-infall onto the central BH cluster)
- II) Why do SMBH masses correlate with the mass of the hosting galaxy?

 Essentially through the downsizing (Thomas) time and via the IGIMF theory.
- III) How can SMBHs form within a few 100 Myr after the Big Bang?
 - The highest-z quasars are hypermassive star-burst clusters with top-heavy IMF (Jerabkova objects).
- IV) 108 quasars / SMBHs exist per co-moving Gpc³
 - V) Thus rapid enrichment everywhere (even if patchy)?

Essentially:

Given

- (i) the standard properties of normal matter
- (ii) the observed downsizing (Thomas) time
- (iii) the correlations of the stellar population with the downsizing (Thomas) time

quasars *must* appear about 200Myr after BigBang as *hypermassive star-burst clusters* (Jerabkova objects) and SMBHs *must* correlate with host galaxy.

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