

# The Initial Mass Function of stars

1st lecture

17th Russbach School  
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# Pavel Kroupa

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*Astronomical Institute,  
Charles University in Prague*

*c/o Argelander-Institut für Astronomie  
University of Bonn*

Pavel Kroupa: University of Bonn / Charles University

One of the most important distribution functions in astrophysics

## Important for :

- the luminosity of a stellar system
  - its mass locked up in long lived late-type stars
  - the rate of CCSN per late-type star
  - the rate of merging white dwarf binaries (SN1a) per late-type star
  - the rate of merging neutron star binaries per late-type star
  - the rate of merging BHs per late-type star

The IMF :  $\xi(m)$

$$dN = \xi(m) dm$$

= number of stars with initial  
 $m \in [m, m + dm]$

One of the most important distribution functions in astrophysics

How do we constrain the shape of the IMF ?  
Does the shape vary ?

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## Prelude

# Galaxies



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# Clusters



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# Star Formation



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# Star Formation



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The outcome of the star-formation process appears to be messy and very complicated.

It is easy to argue that star formation is *stochastic*, but the challenge is to find *patterns* and *rules*.

One such effort is to constrain the shape of the initial distribution function of stellar masses, the **IMF**.

Related to the above issue is the question whether the IMF is a *probability density* distribution function  
or  
an *optimally sampled* distribution function ?

(Kroupa et al., 2013, arXiv:1112.3340)

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**Finding the shape of the IMF**

**The Galactic-field LF**

We have  $dN = \Psi(M_V) dM_V$  = # of stars with  
 $M_V \in [M_V, M_V + dM_V]$

$dN = \xi(m) dm$  = # of stars with  
 $m \in [m, m + dm]$

since

$$\frac{dN}{dM_V} = -\frac{dm}{dM_V} \frac{dN}{dm}$$

follows

$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$

the observable

the obstacle

the target

the stellar luminosity function (LF)

the stellar mass function  
(IMF if initial)

II

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There are *two luminosity functions*  
for the solar neighbourhood

## I. Count stars nearby to Sun

Obtain  $M_V$  and  $d$  from trigonometric parallax



Well observed individual stars but  
*small numbers at faint end* ( $\Psi_{\text{near}}$ )

# There are *two luminosity functions* for the solar neighbourhood

## I. Count stars nearby to Sun

Obtain  $M_V$  and  $d$  from trigonometric parallax



Well observed individual stars but  
*small numbers at faint end* ( $\Psi_{\text{near}}$ )

## II. Deep (100 - 300 pc) pencil-beam photographic/CCD surveys

Formidable data reduction ( $10^5$  images  $\rightarrow \approx 100$  stars)

Obtain  $M_V$  and  $d$  from photometric parallax



Large # of stars but *poor resolution* (2"-3")  
and *Malmquist bias* ( $\Psi_{\text{phot}}$ )

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The possibility of *dark matter* in the *Galactic disk*

(Bahcall 1984)



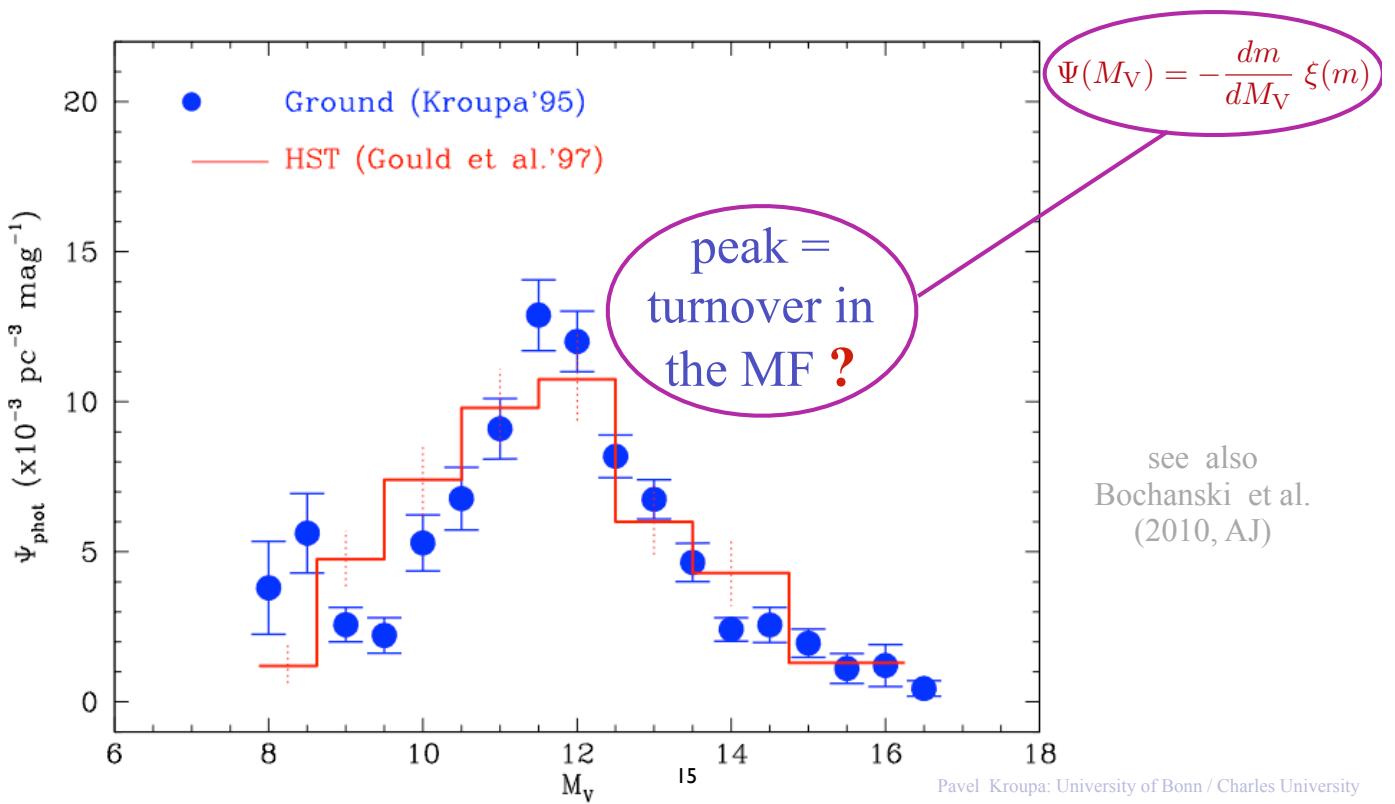
Many surveys of type II (pencil-beams)  
to constrain the LF :

ground	Reid & Gilmore	1982
	Gilmore, Reid & Hewett	1985
	Hawkins & Bessell	1988
	Leggett & Hawkins	1988
	Stobie, Ishida & Peacock	1989
	Tinney, Reid & Mould	1993
HST	Kirkpatrick et al.	1994
	Gould, Bahcall & Flynn	1997
	Zheng, Flynn, Gould et al.	2001

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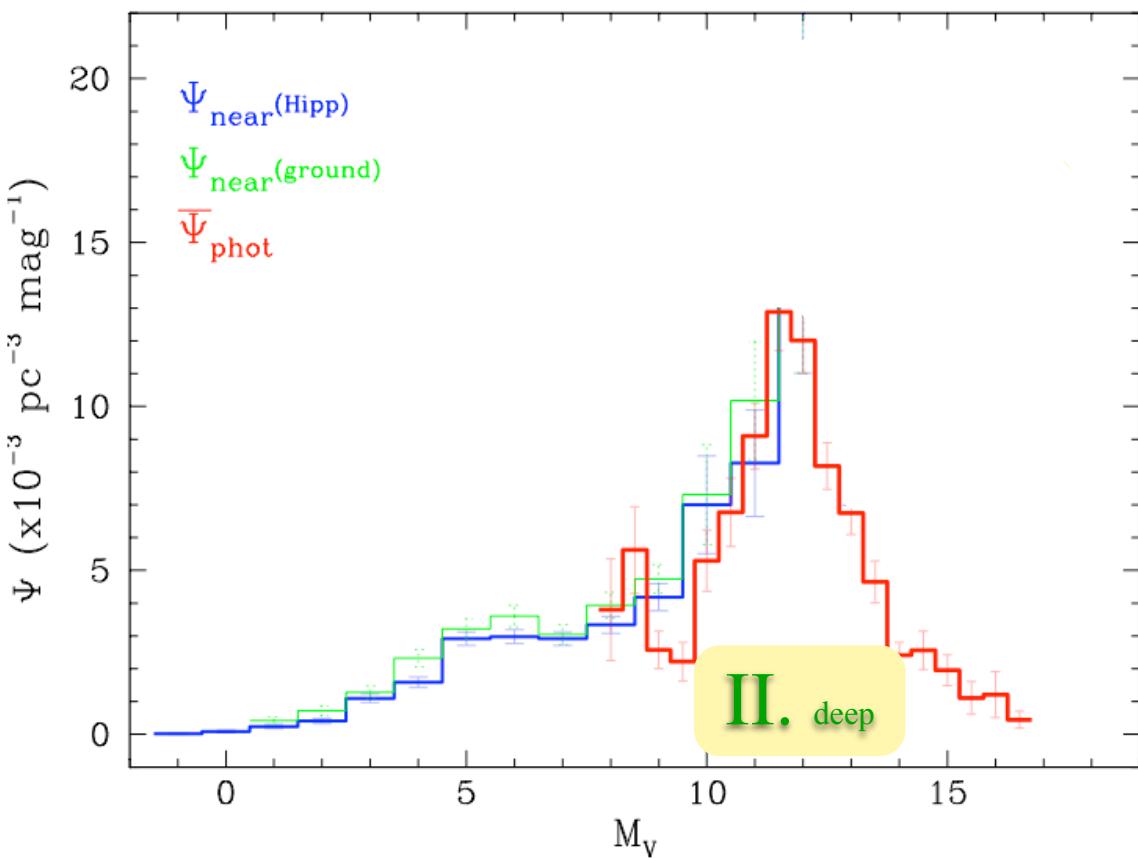
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↓  
 $\Psi_{\text{phot}}$  { - independent of direction  
 - maximum (peak) at  $M_V \approx 12$



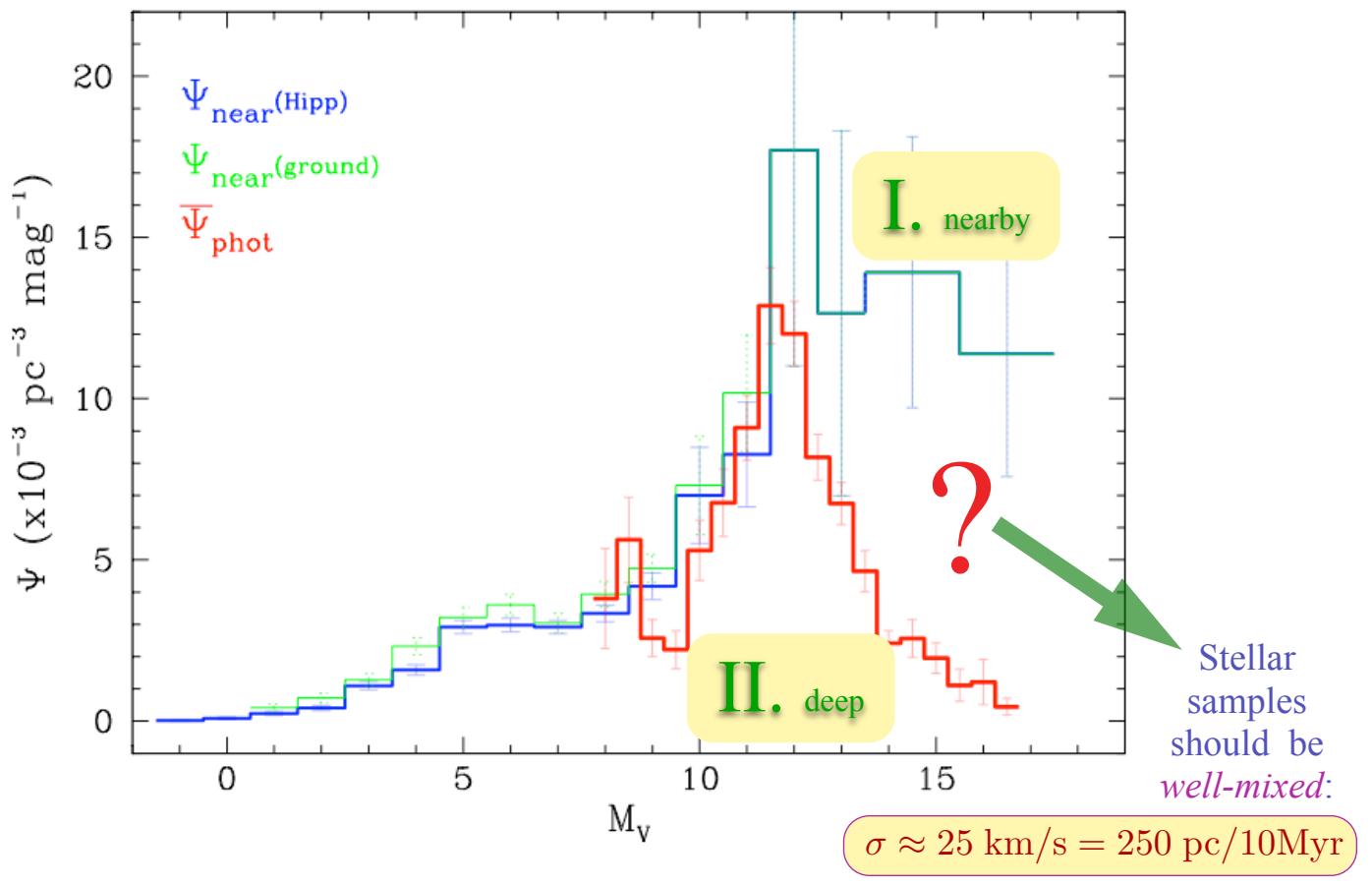
## Two solar-neighbourhood samples:

$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$



## BUT: two solar-neighbourhood samples:

$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$



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## Problem :

The nearby and deep LFs are not equal.

→ **Which LF** do we use to calculate the **MF** ?

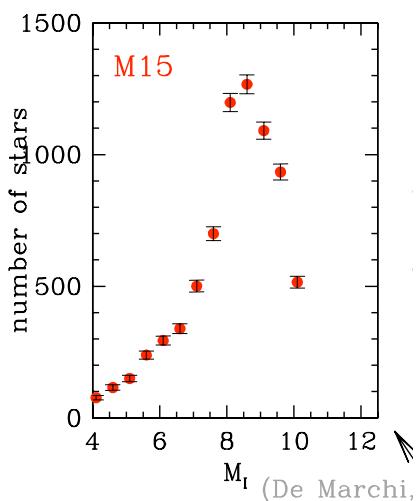
$$\xi(m) = - \left( \frac{dm}{dM_V} \right)^{-1} \Psi(M_V)$$

$\Psi_I$  OR  $\Psi_{II}$  ??

# The LF of stars in star clusters coeval, equidistant stars!

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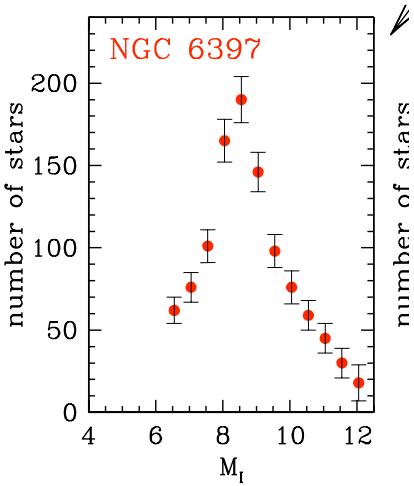
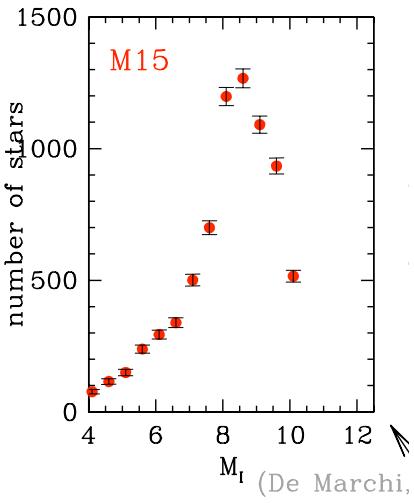
Kroupa 2002  
Science

$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$

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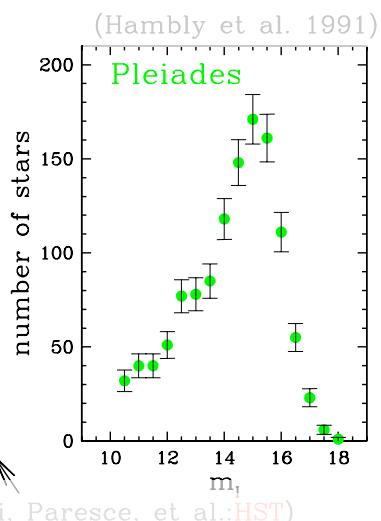
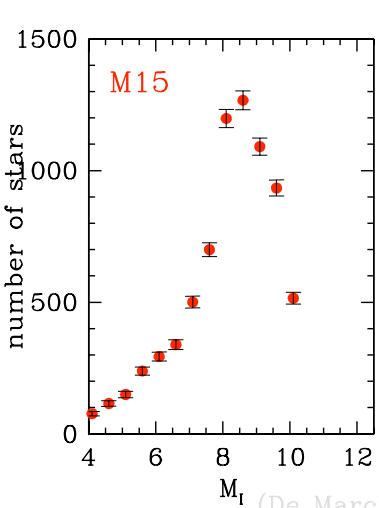
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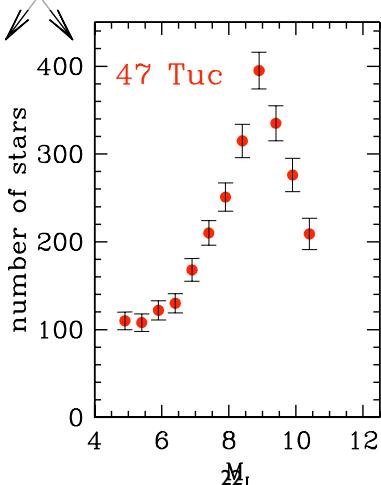
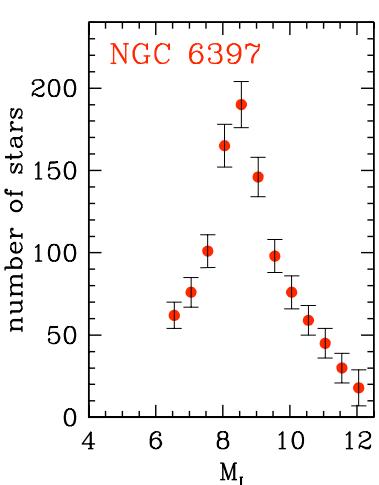


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$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$



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I. nearby

II. deep

Looking  
into  
the stars  
by  
just counting them

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*The mass-  
luminosity relation  
of low-mass stars*

Kroupa, Tout & Gilmore 1990;  
Kroupa, 2002, *Science*

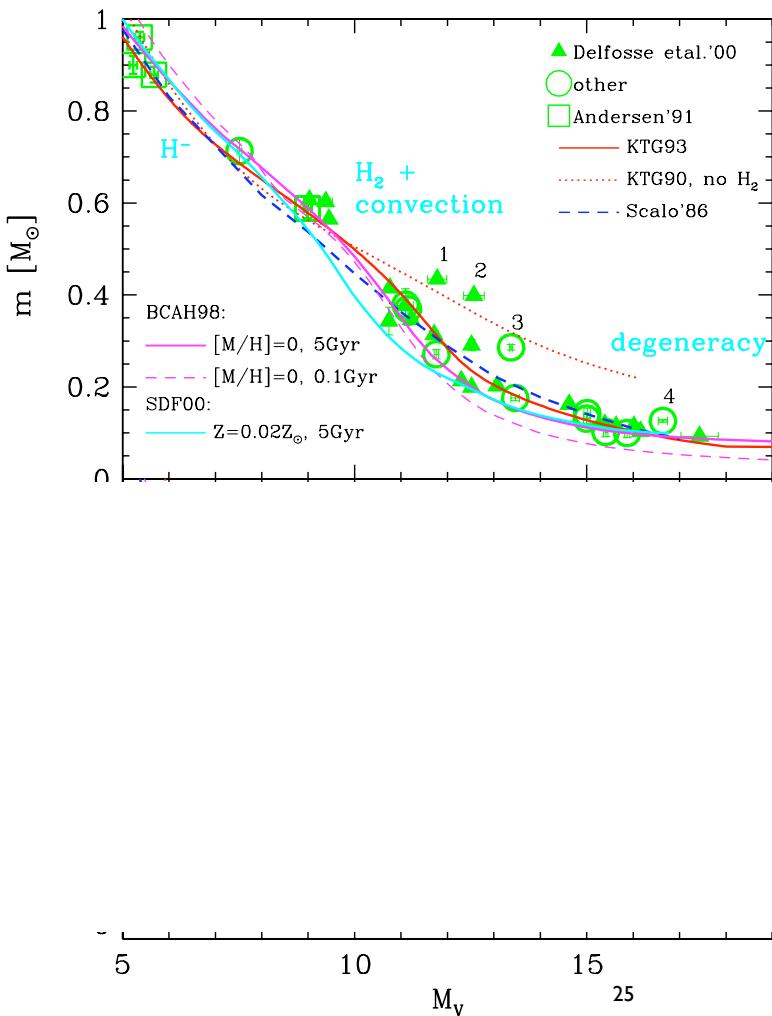
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# The mass-luminosity relation of low-mass stars

Kroupa, Tout & Gilmore 1990;  
Kroupa, 2002, *Science*

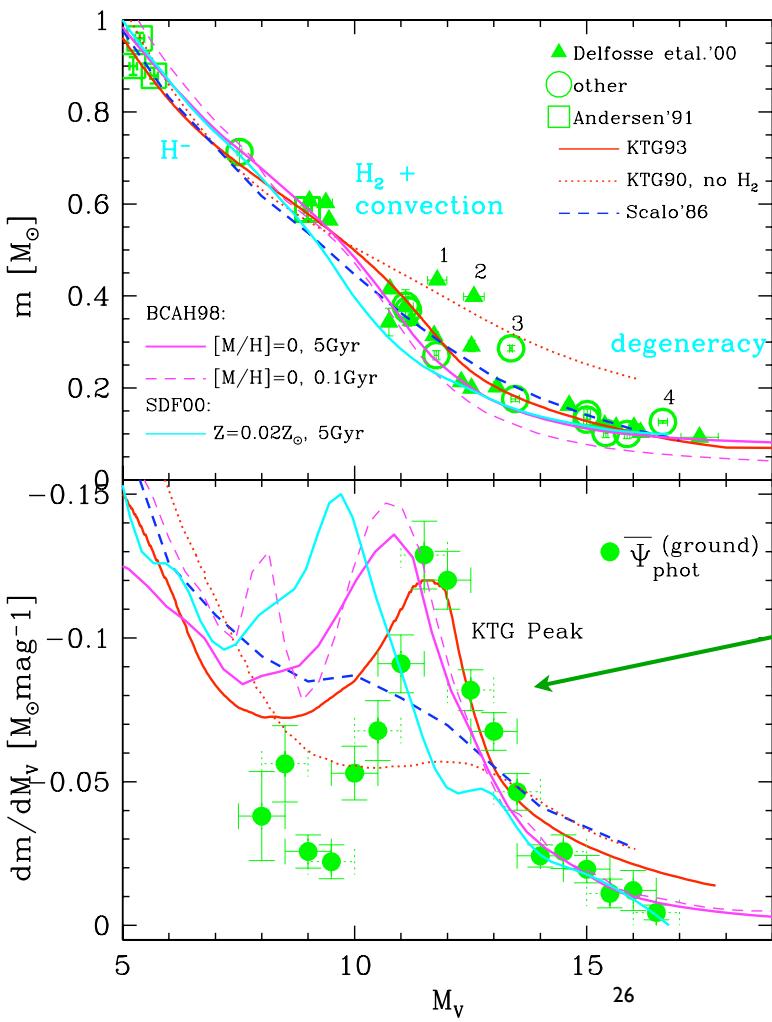


$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$

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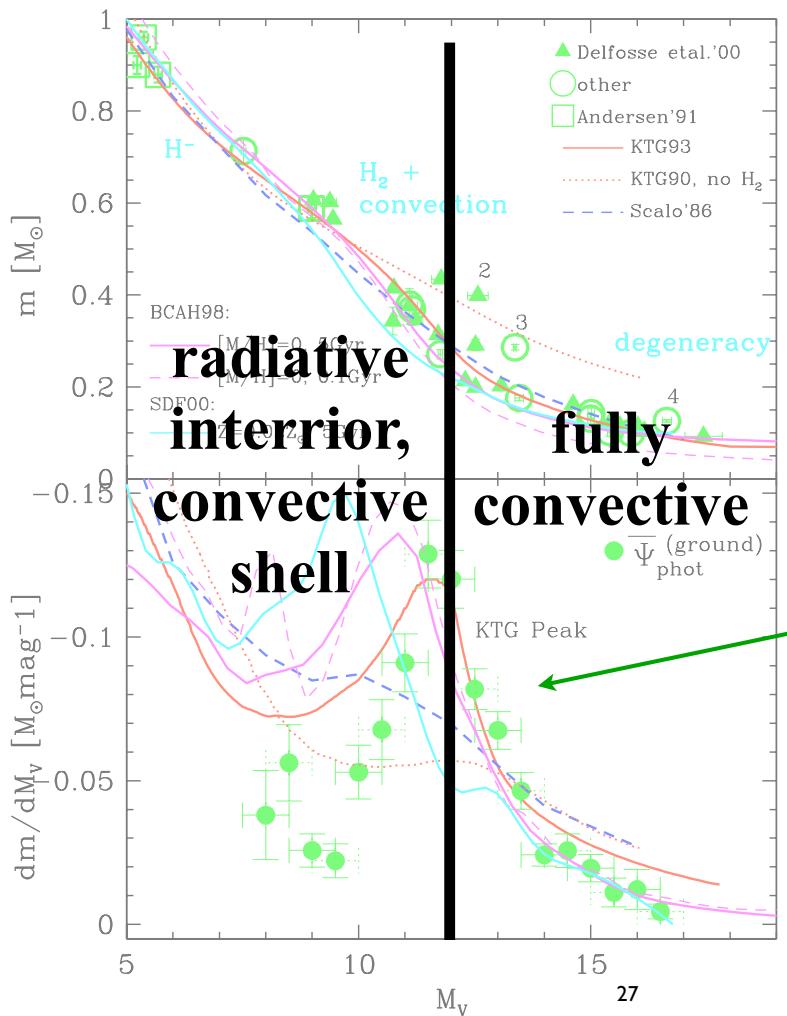
$$\Psi(M_V) = -\frac{dm}{dM_V} \xi(m)$$

- 1. position
- 2. width
- 3. amplitude
- all agree!*

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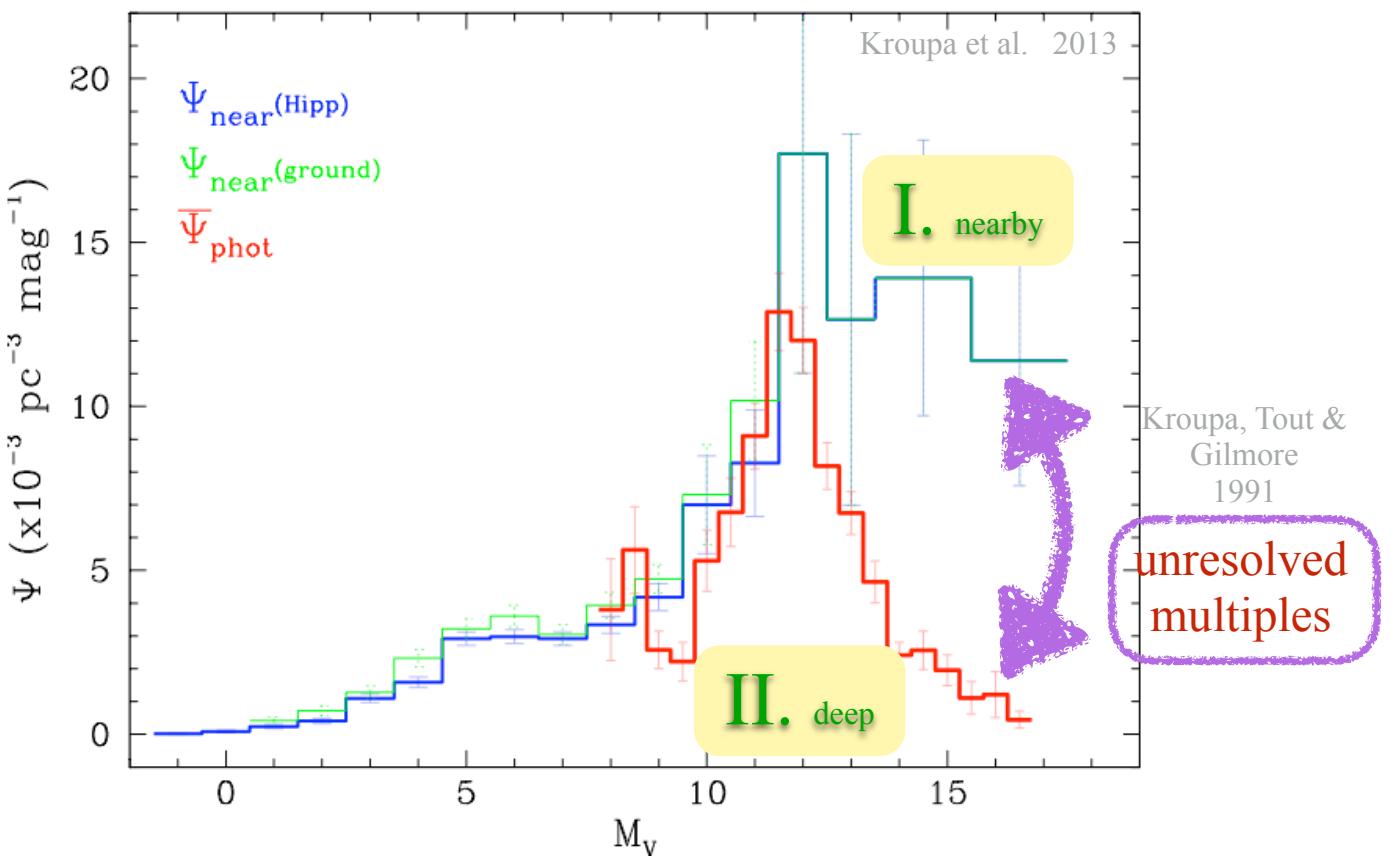
$$\Psi(M_V) = - \frac{dm}{dM_V} \xi(m)$$

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## Two solar-neighbourhood samples:

$$\Psi(M_V) = - \frac{dm}{dM_V} \xi(m)$$



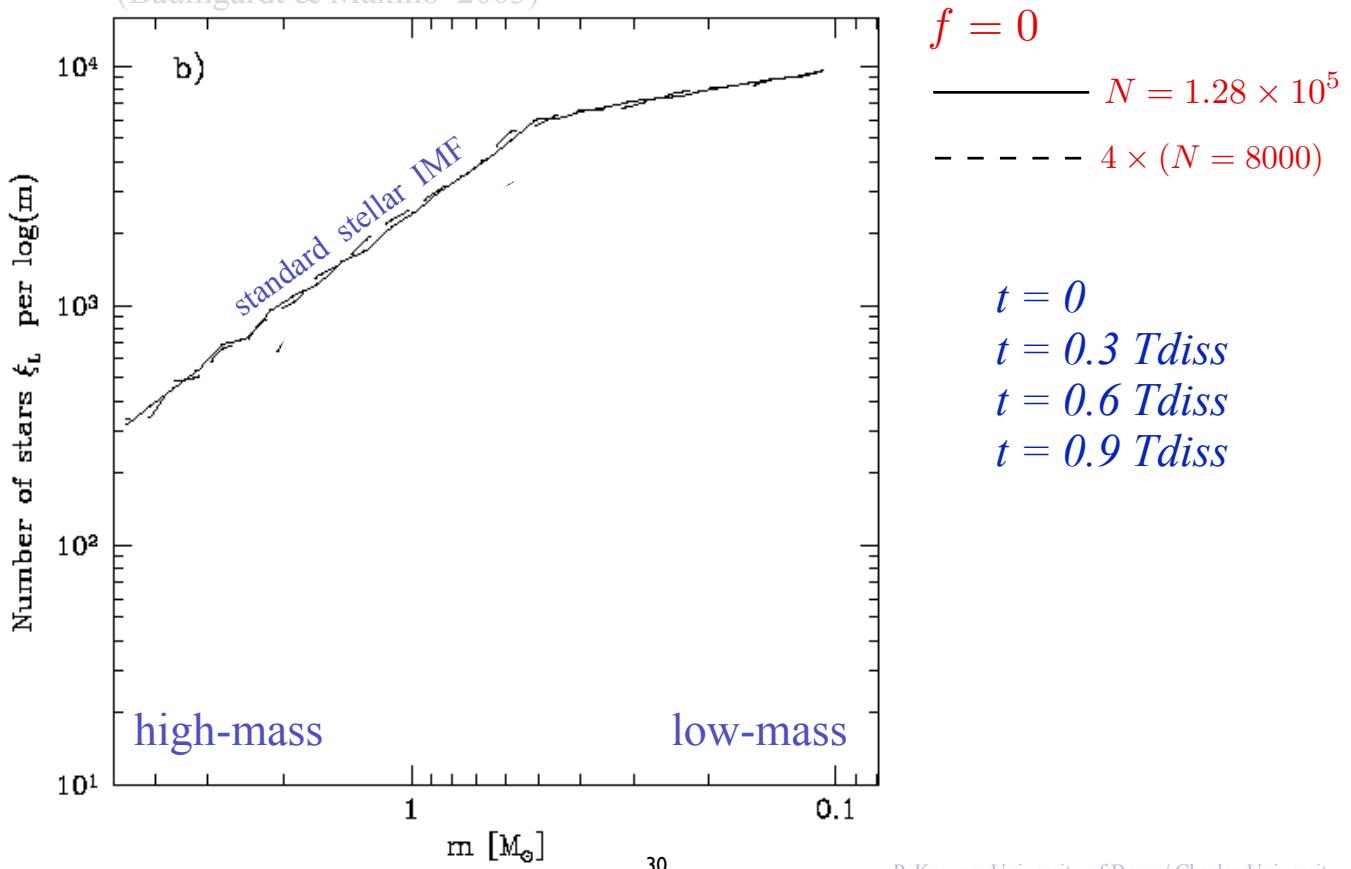
# Inferring the IMF in star clusters

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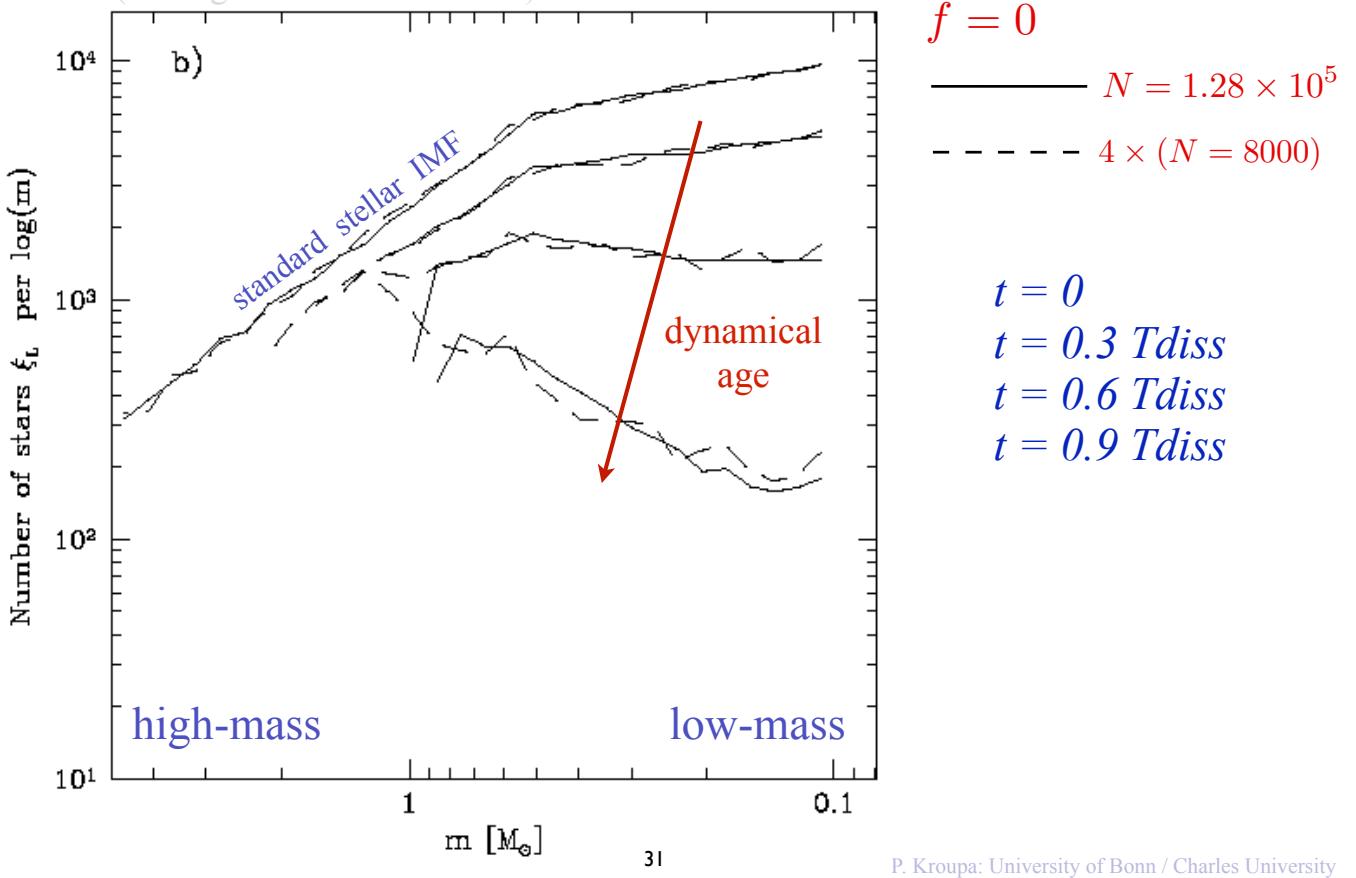
## *MF( $t$ ) due to cluster evolution*

(Baumgardt & Makino 2003)



# $MF(t)$ due to cluster evolution

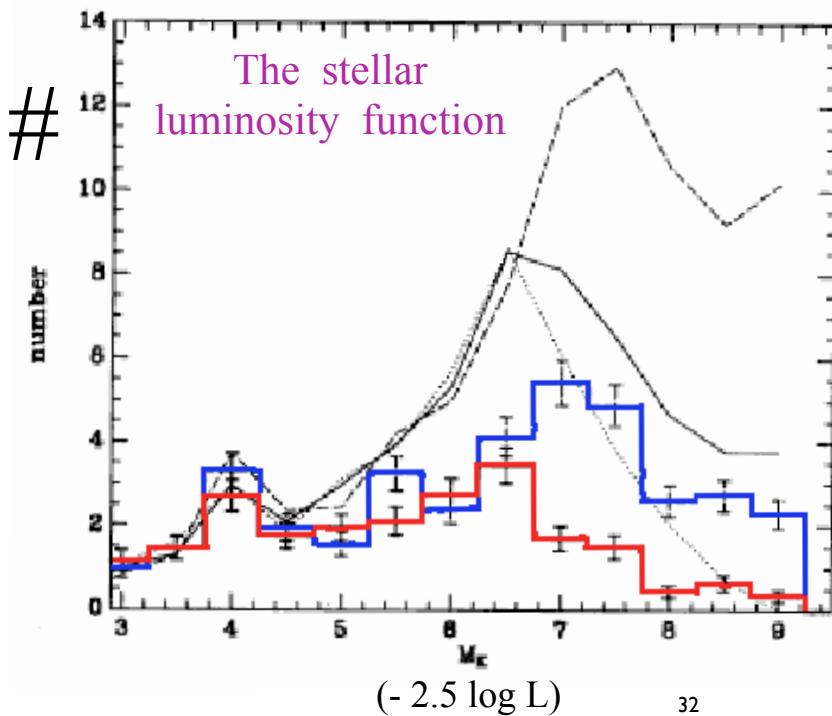
(Baumgardt & Makino 2003)



## *N-body Models of Binary-Rich Clusters*

(Kroupa 1995)

$20 \times (N = 400 \text{ stars})$   
 $f = 1$



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# Massive stars in very young clusters



## *OB stars in clusters / HII regions*

Two competing processes:

### *Mass segregation*

$$t_{\text{msgr}} \approx 2 \left( \frac{m_{\text{av}}}{m_{\text{massive}}} \right) t_{\text{relax}}$$

$$t_{\text{relax}} = \frac{21}{\ln(0.4N)} \left( \frac{M_{\text{ecl}}}{100 M_{\odot}} \right)^{\frac{1}{2}} \left( \frac{1 M_{\odot}}{m_{\text{av}}} \right) \left( \frac{R_{0.5}}{1 \text{ pc}} \right)^{\frac{3}{2}}$$

e.g.  $t_{\text{relax}} \approx 0.6 \text{ Myr}$   
for pre-exposed ONC



$$t_{\text{msgr}} \approx 0.12 \text{ Myr} \ll \text{age of ONC}$$

### *Core decay*

$$t_{\text{decay}} \approx N_{\text{m}} \times t_{\text{core,cross}}$$

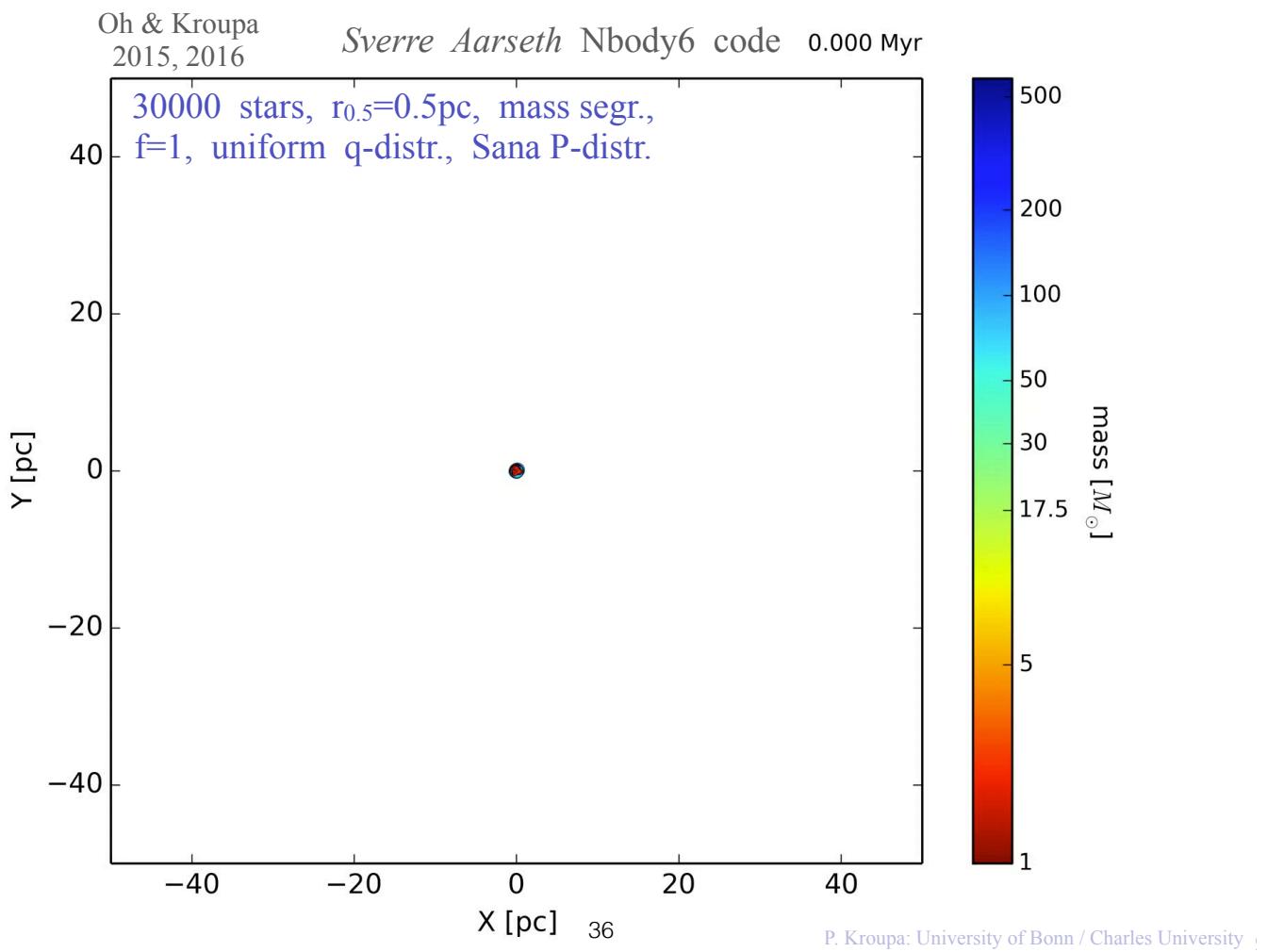
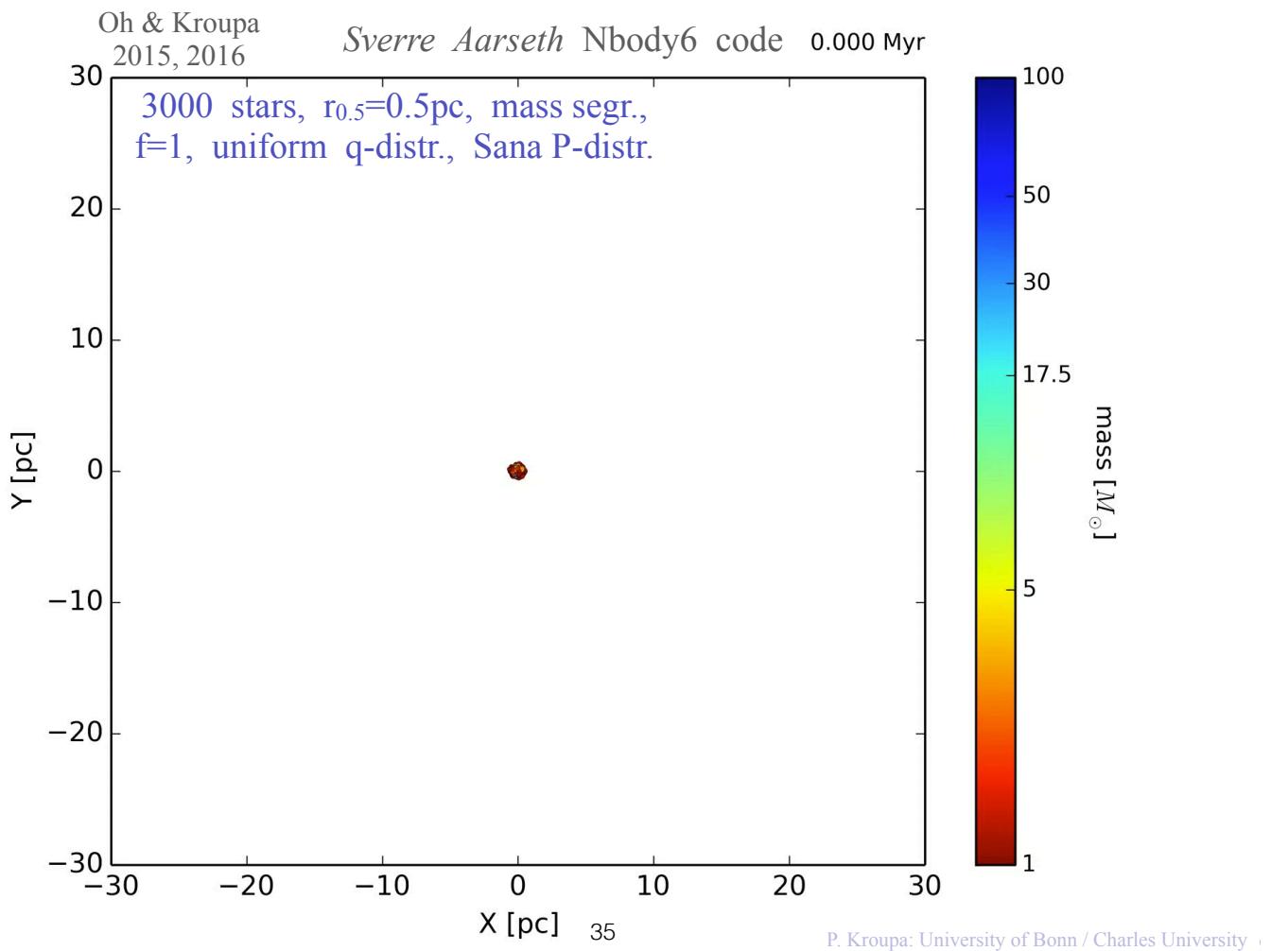
$$t_{\text{cross}}^{\text{core}} \approx 5 \left( \frac{M_{\text{core}}}{100 M_{\odot}} \right)^{-\frac{1}{2}} \left( \frac{R_{0.5}}{1 \text{ pc}} \right)^{\frac{3}{2}}$$

e.g.  $R_{\text{core}} \approx 0.02 \text{ pc}$ ,  $M_{\text{core}} \approx 150 M_{\odot}$

$$t_{\text{cross}}^{\text{core}} \approx 1.2 \times 10^4 \text{ yr}$$



$$t_{\text{decay}} \approx 10^4 - 10^5 \text{ yr} \ll \text{age of ONC}$$



Clusters depopulate themselves  
off low-mass stars and high mass stars.



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Thus, stellar-dynamical processes  
are  
extremely important  
when determining the IMF shape!!

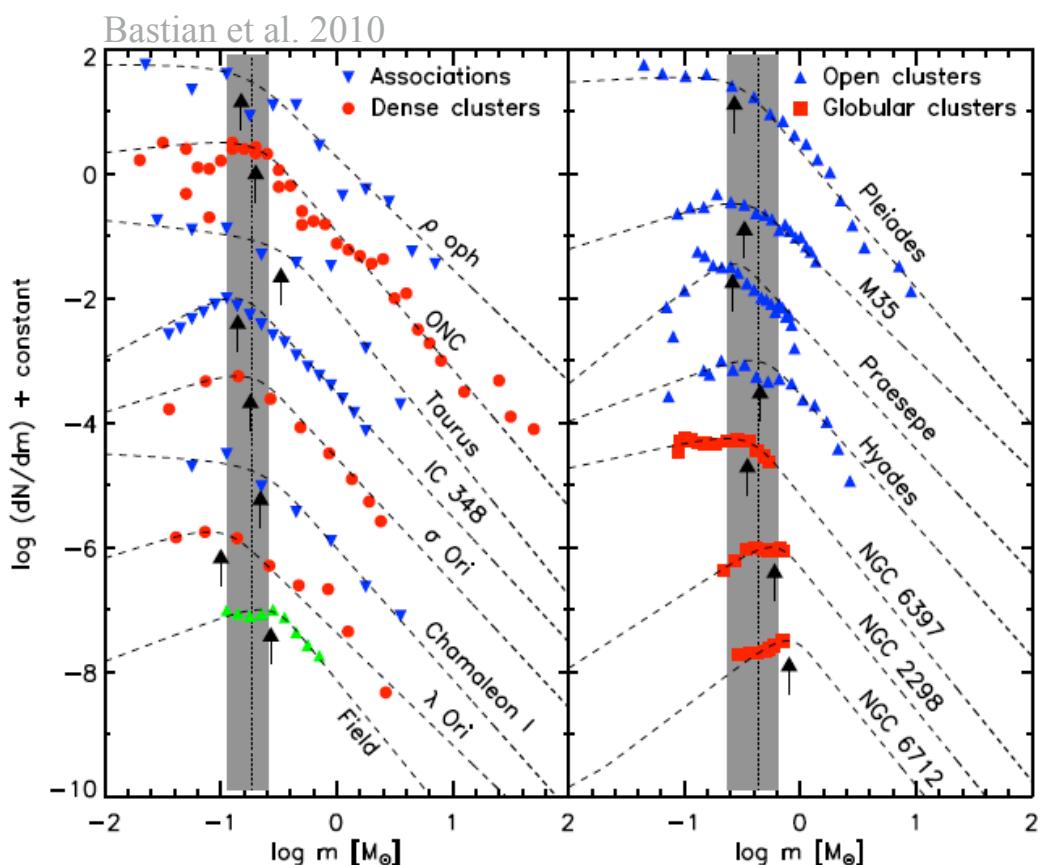
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# An invariant canonical IMF

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The IMF appears largely invariant (in MW CSFEs / embedded clusters)

Kroupa 2001, 2002

Bastian, Covey,  
Meyer,  
2010, ARAA

Kroupa et al.,  
2013

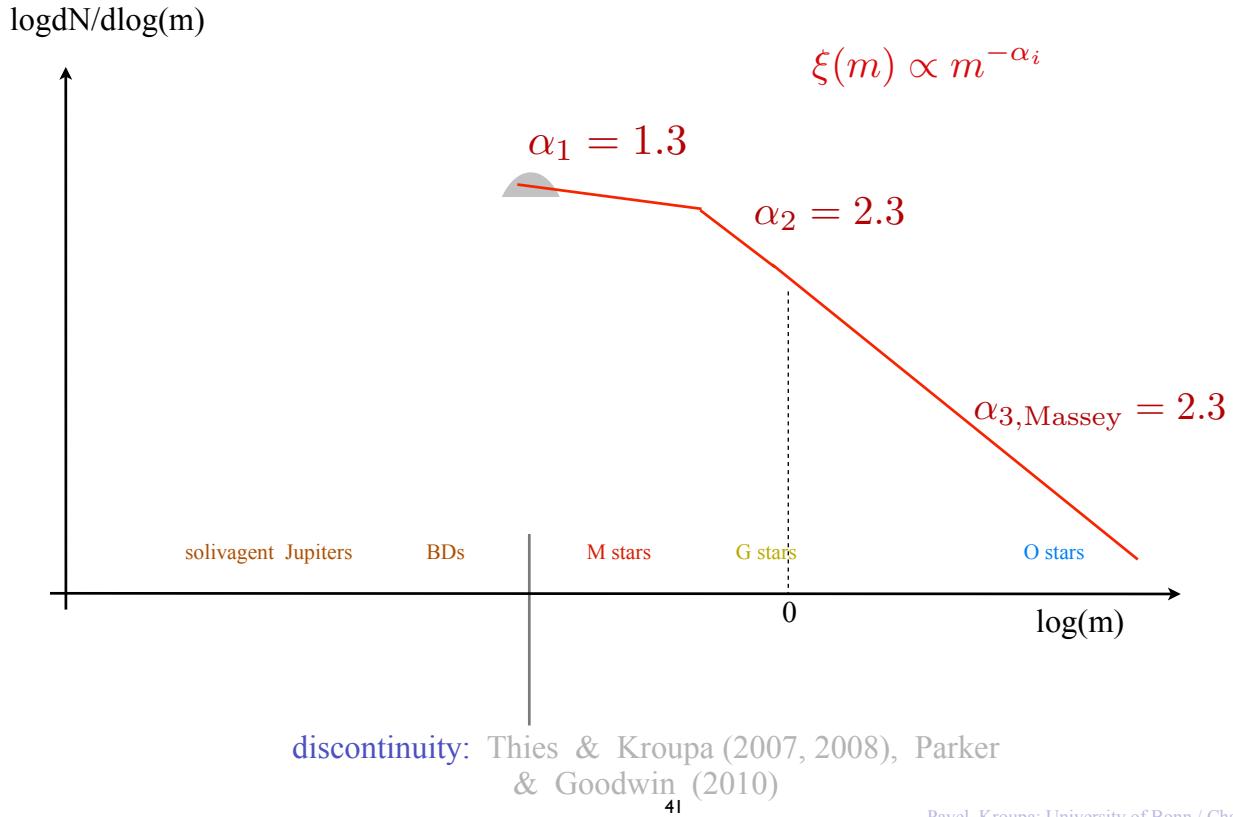
Offner et al.,  
2014, PPVI

40

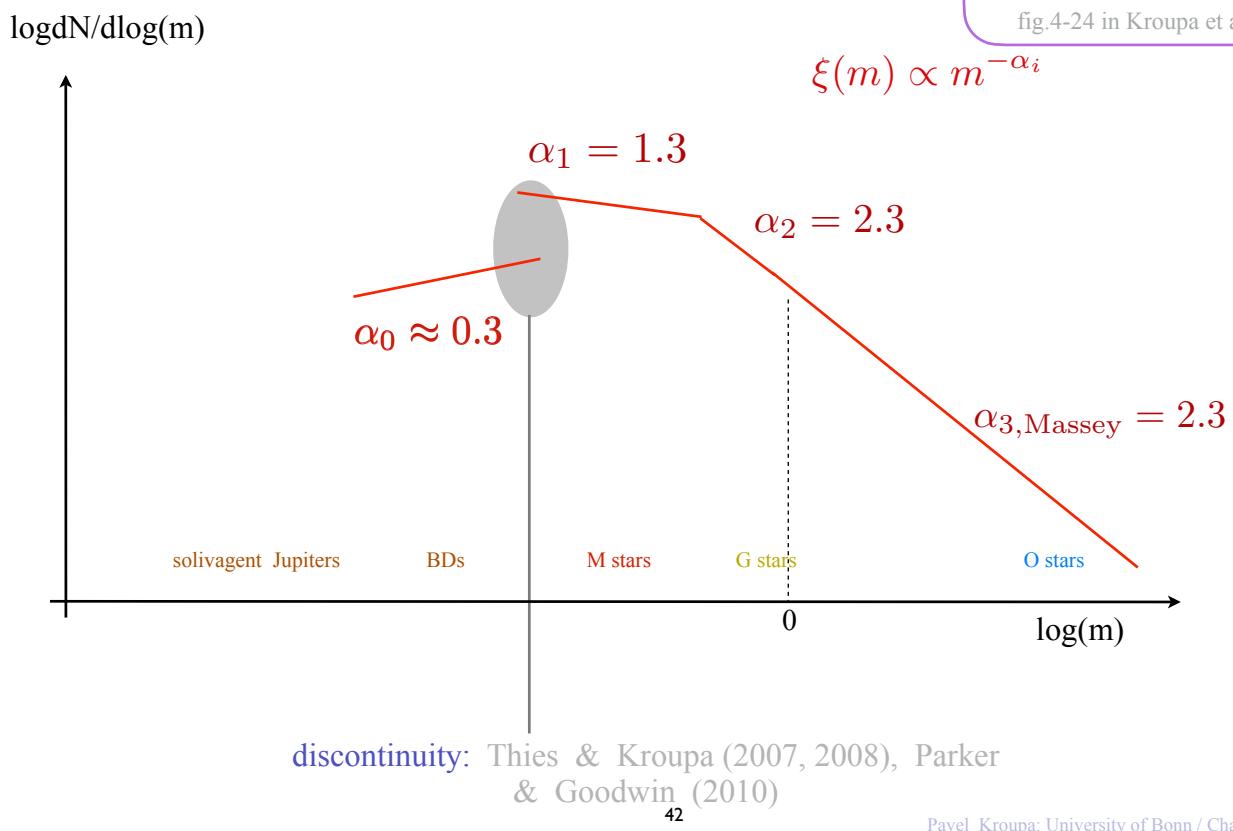
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### *the canonical IMF :*



### *the canonical IMF :*

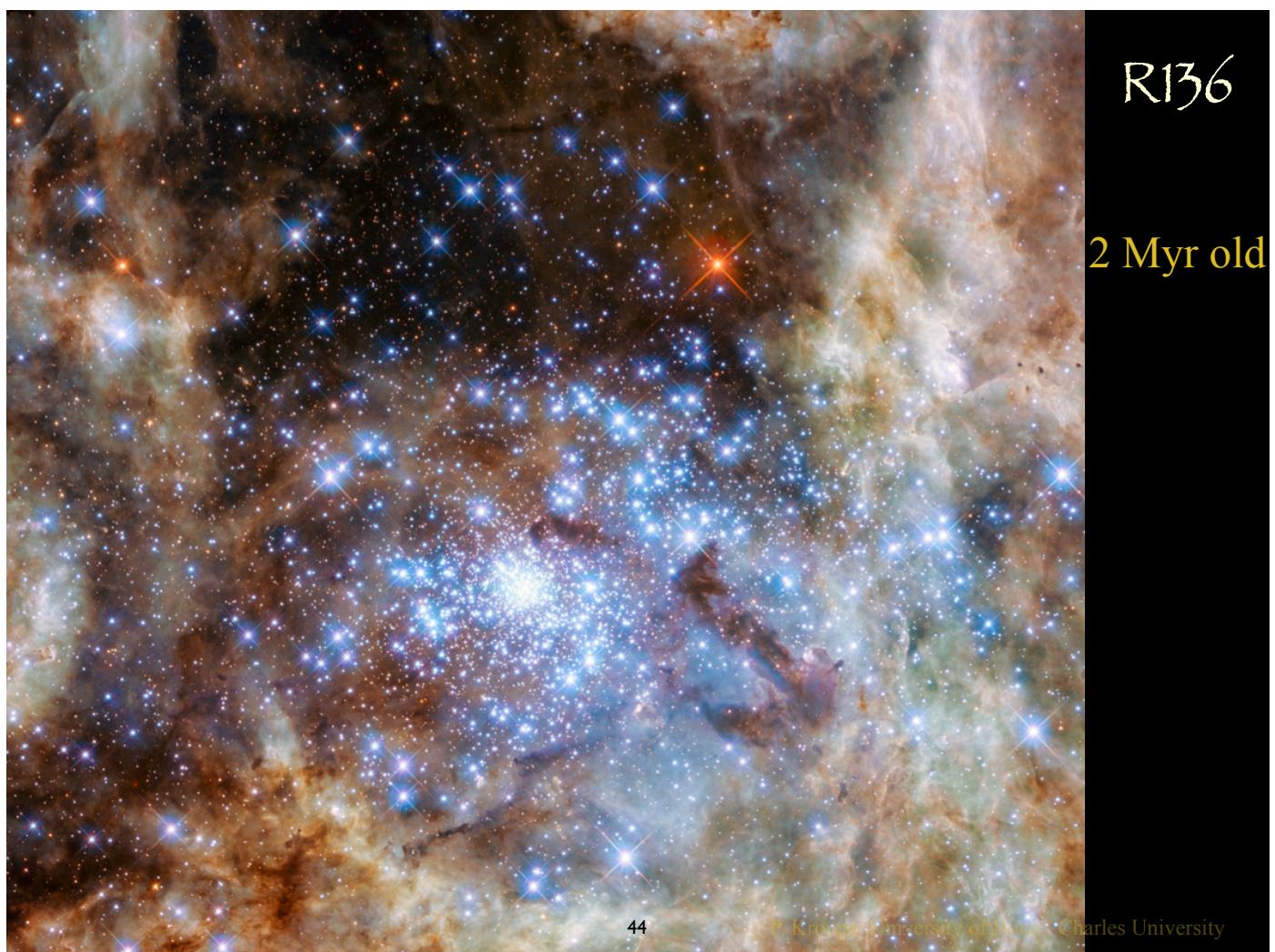


# Extreme environments: evidence for a systematically-varying IMF

evidence began emerging  
since about 2010

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## Some subtle hints for a systematically varying IMF are available at high masses

**Star-counts:** Correct star-counts in R136 for ejected stars  
→ IMF in R136 top-heavy (Banerjee & Kroupa 2012)

Excess of massive stars in whole 30Dor region  
(Schneider et al. 2018)

Ultra-violett spectroscopy of seven nearby (< 120 Mpc) star-bursting regions imply the same (Senchyna et al. 2021).

Top-heavy IMF in Magellanic Bridge cluster NGC796  
(Kalari et al. 2018)

GCs in M31: more top-heavy IMF at lower metallicity  
(Zonoozi et sl. 2016; Haggi et al. 2017)

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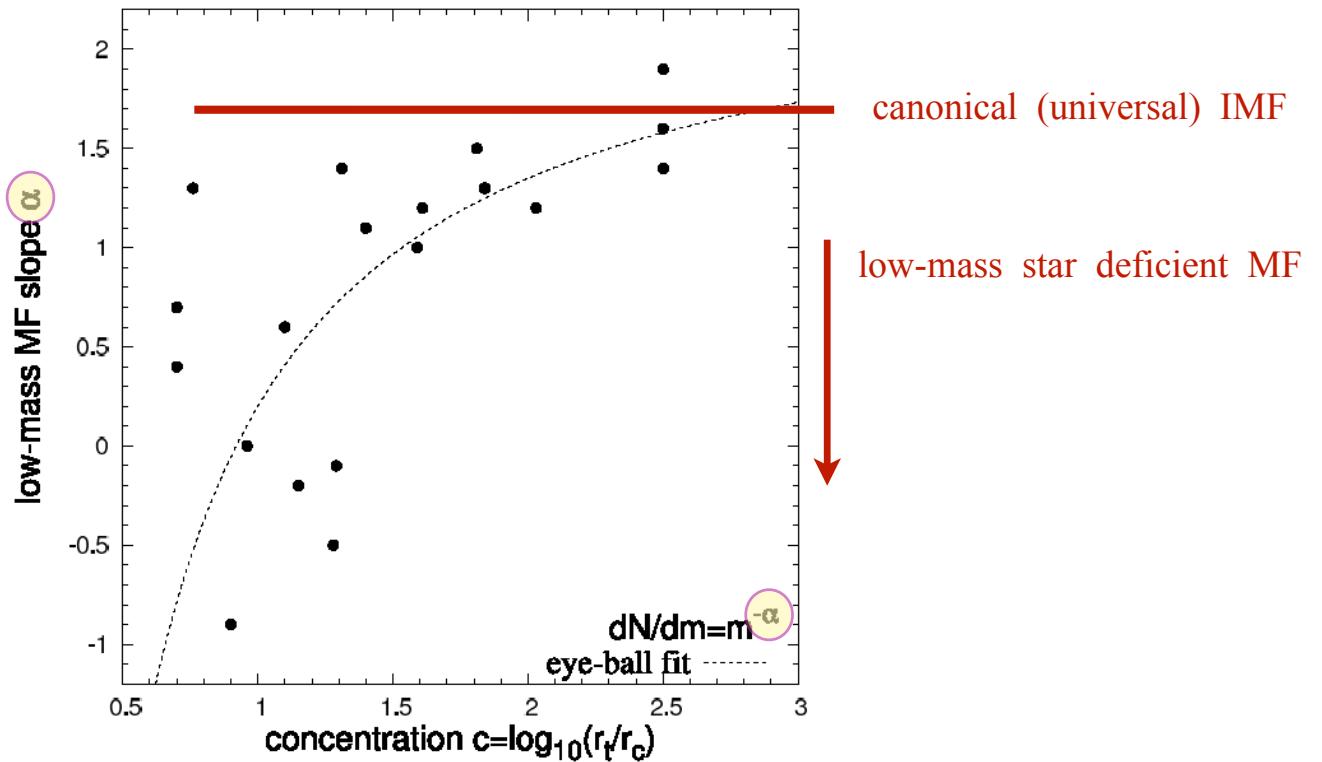
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## *What we know from observation :*

Globular clusters : deficit of low-mass stars increases with decreasing concentration  
→ disagrees with dynamical evolution (Marks et al. 2012)

*A sample of 20 Galactic GCs  
with solid global MF measurements from  
deep HST or VLT data.*

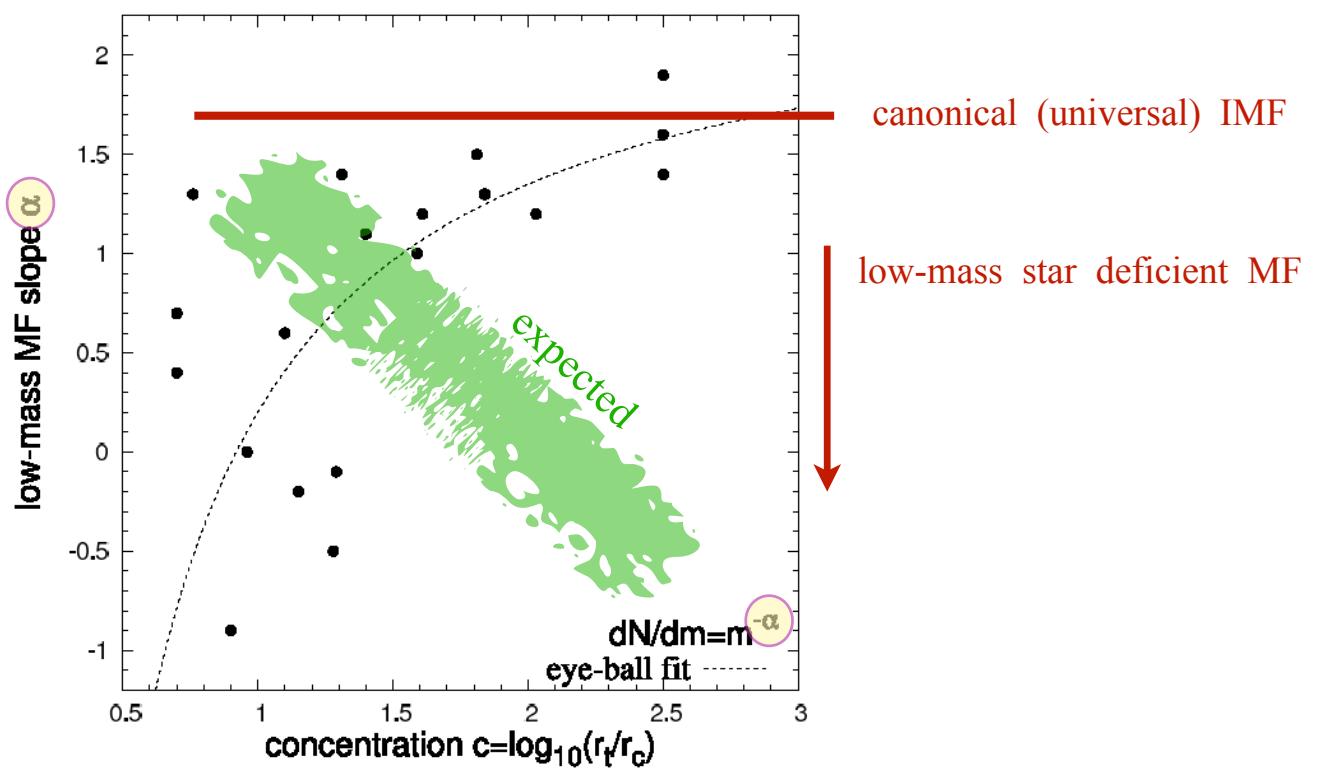
(de Marchi, Paresce & Pulone 2007)



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deep HST or VLT data.*

(de Marchi, Paresce & Pulone 2007)



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# *GCs* *(extreme star burst "clusters")*

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gas expulsion  
+  
mass segregation  
!

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# *Cluster reaction to thermal gas removal :*

Time = 0.0 Myr  
Gas content: 100%

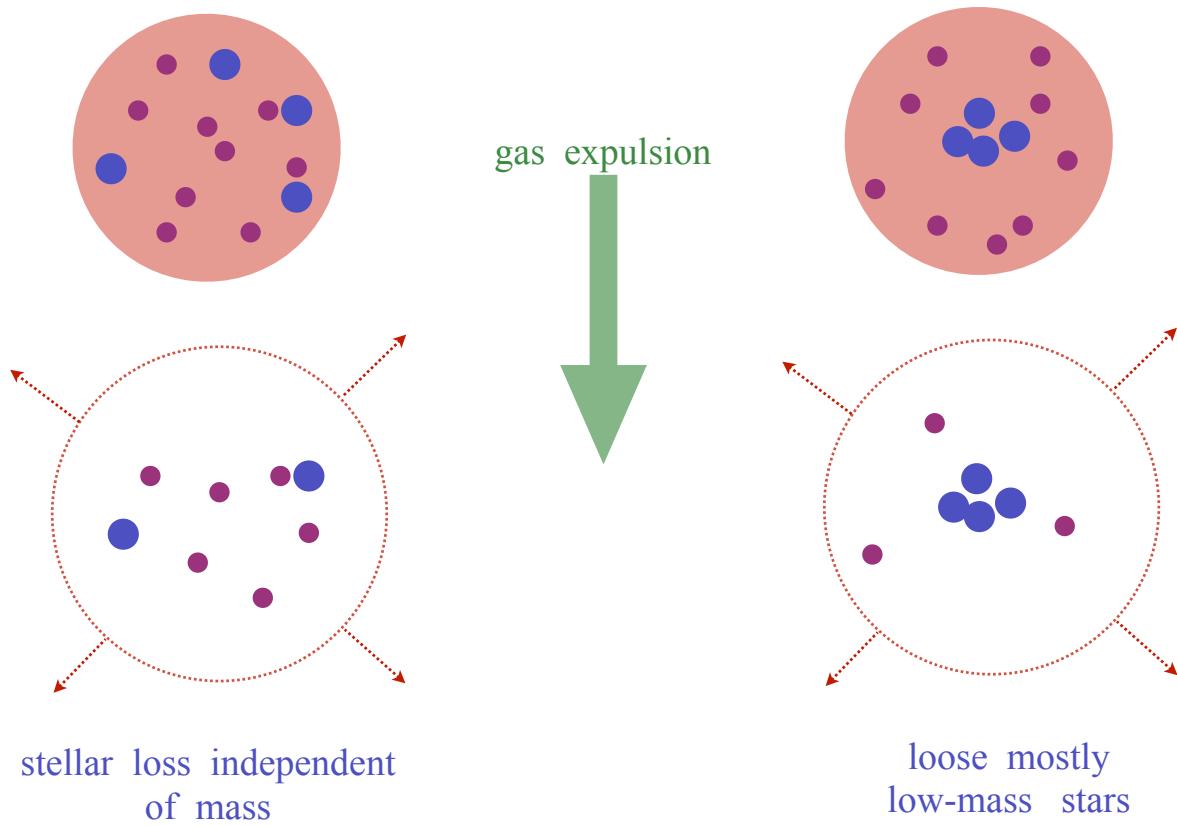
(movie by Baumgardt)

Baumgardt & Kroupa 2007

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## *Nbody models of binary rich initially mass segregated clusters with residual gas expulsion after birth*

(Marks, Kroupa & Baumgardt 2008)



## What we know from observation :

Globular clusters : deficit of low-mass stars increases with decreasing concentration

→ disagrees with dynamical evolution (Marks et al. 2012)  
(need gas expulsion from mass-segregated clusters)

UCDs : higher dynamical M/L ratios

(Dabringhausen et al. 2009)

→ cannot be exotic dark matter

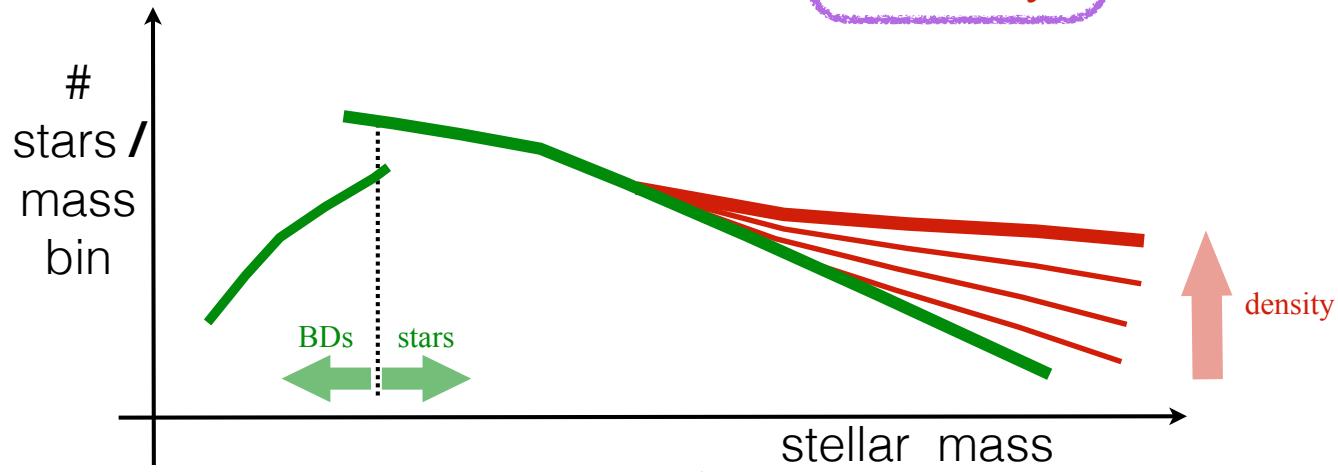
UCDs : larger fraction of X-ray sources than expected

(Dabringhausen et al. 2012)

→ no explanation other than many remnants

## What this implies :

mutual  
consistency !!



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$$\text{IMF} = \text{IMF}(Z, SFRD)$$

$Z$ =metallicity,  $SFRD$ =star-formation rate density

Variation of IMF in pc-scale, 1-Myr time-scale  
(molecular cloud core = embedded cluster).

## Top-heavy IMF in extreme-density environments : $\xi(m) \propto m^{-\alpha_i}$

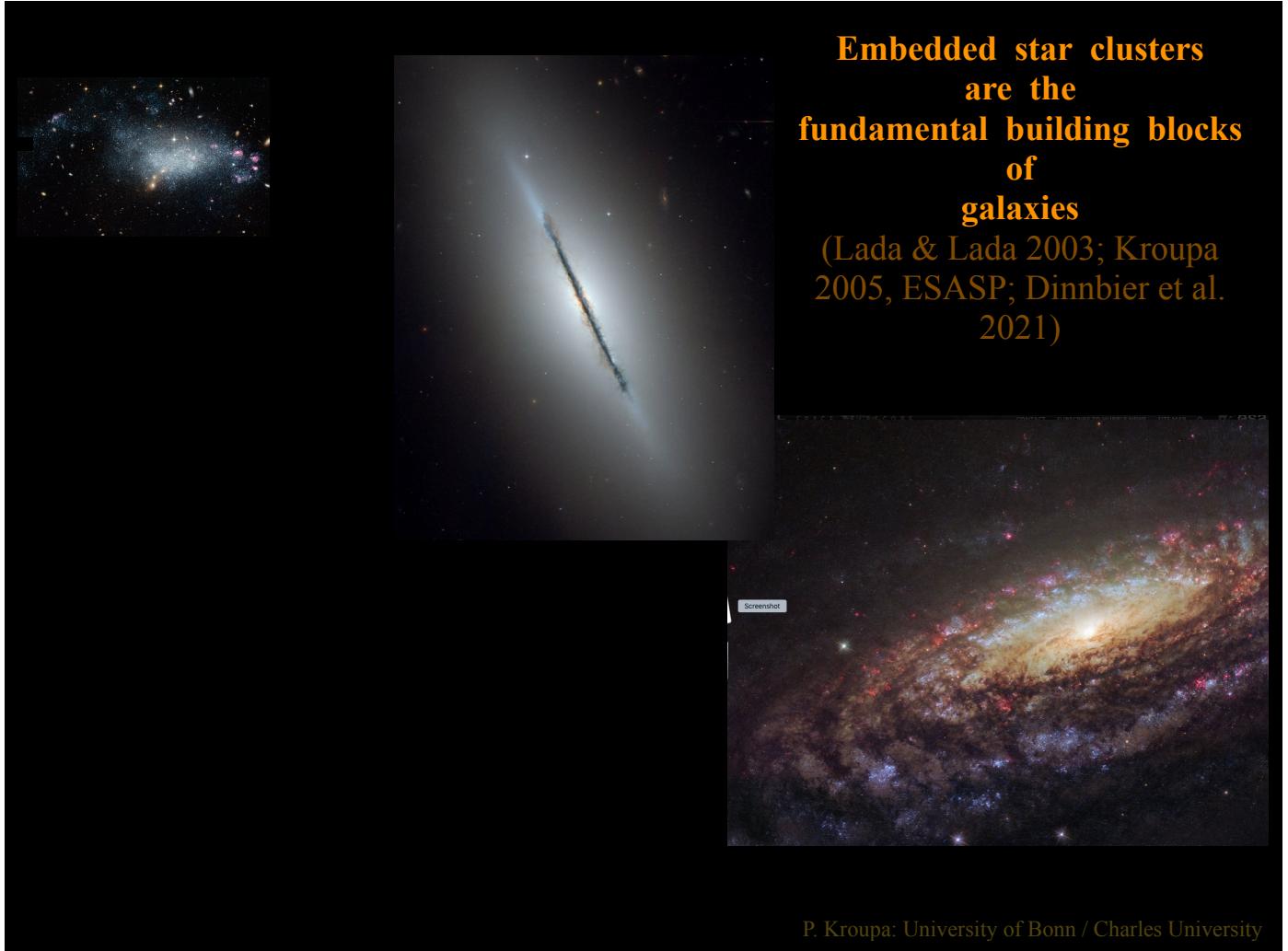
THE STELLAR IMF DEPENDENCE ON DENSITY AND METALLICITY: Resolved stellar populations show an invariant IMF (Eq. 55), but for  $SFRD \gtrsim 0.1 M_\odot/(\text{yr pc}^3)$  the IMF becomes top-heavy, as inferred from deep observations of GCs. The dependence of  $\alpha_3$  on cluster-forming cloud density,  $\rho$ , (stars plus gas) and metallicity, [Fe/H], can be parametrised as

$$\begin{aligned}\alpha_3 &= \alpha_2, & m > 1 M_\odot \wedge x < -0.89, \\ \alpha_3 &= -0.41 \times x + 1.94, & m > 1 M_\odot \wedge x \geq -0.89, \\ x &= -0.14 [\text{Fe}/\text{H}] + 0.99 \log_{10} (\rho / (10^6 M_\odot \text{pc}^{-3})) .\end{aligned}\tag{65}$$

Marks et al. 2012  
Kroupa et al. 2013  
Yan et al. 2017  
Jerabkova et al. 2018  
Yan, Jerabkova et al. 2021

newest, most  
up-to-date formulation

Whole galaxies  
are made of  
all their IMFs



## Composite Stellar Populations

Stars form in a clustered mode (Lada & Lada 2003; Megeath et al. 2016). Thus, the Integrated Galactic IMF follows from

$$\text{IGIMF} = \sum \text{ of IMFs (in all new clusters)}$$



The  
**IGIMF**  
theory

Kroupa & Weidner 2003  
Kroupa et al. 2013

$$\xi_{\text{IGIMF}}(m, t) = \int_{M_{\text{ecl,min}}}^{M_{\text{ecl,max}}(SFR(t))} \xi(m \leq m_{\max}(M_{\text{ecl}})) \xi_{\text{ecl}}(M_{\text{ecl}}) dM_{\text{ecl}}$$

Kroupa & Weidner (2003); Weidner & Kroupa (2005, 2006)  
Yan, Jerabkova & Kroupa (2017); Jerabkova et al. (2018)



Natural explanation of the mass-metallicity relation,  
*radial Halpha cutoff*  
*vs UV extended disks*  
of galaxies  
and many other problems in understanding galaxies.

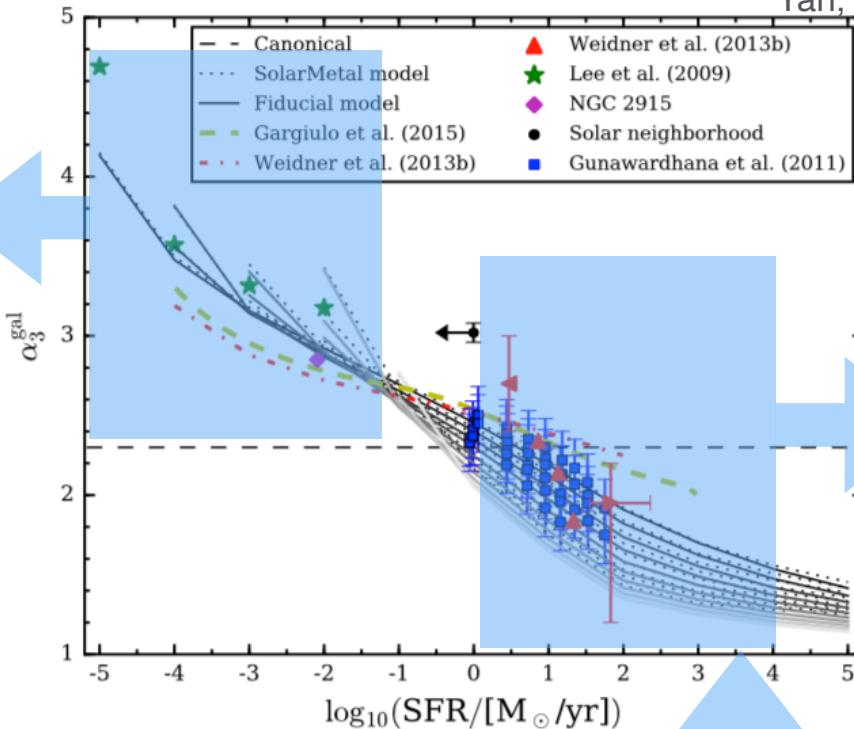
Jerabkova et al. 2018; Yan, Jerabkova et al. 2017; Fontanot et al. 2017;  
and recent papers by Simone Recchi et al.; Sylvia Ploeckinger et al.  
and previous papers by Koeppen et al., Pflamm-Altenburg et al., Weidner et al.

## Consistency Check

Calculate galaxy-wide IMF (= gwIMF) by summing all embedded clusters

Yan, Jerabkova+ 2017

2009  
*Janice Lee*:  
Star forming dwarf galaxies have top-light gwIMF

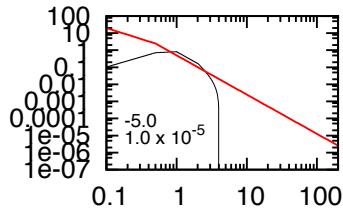


2011  
*Madusha Gunawardhana*:  
Star forming massive galaxies have top-heavy gwIMF

1994 *Francesca Matteucci*:  
formation of ellipticals --> top-heavy gwIMF

## *The IGIMF for galaxies with different SFRs & [Fe/H]*

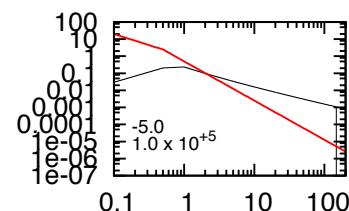
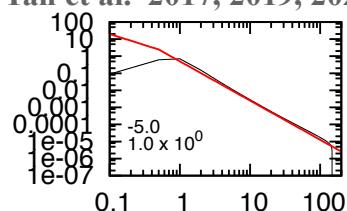
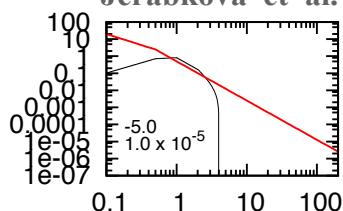
Jerabkova et al., 2018; Yan et al. 2017, 2019, 2021



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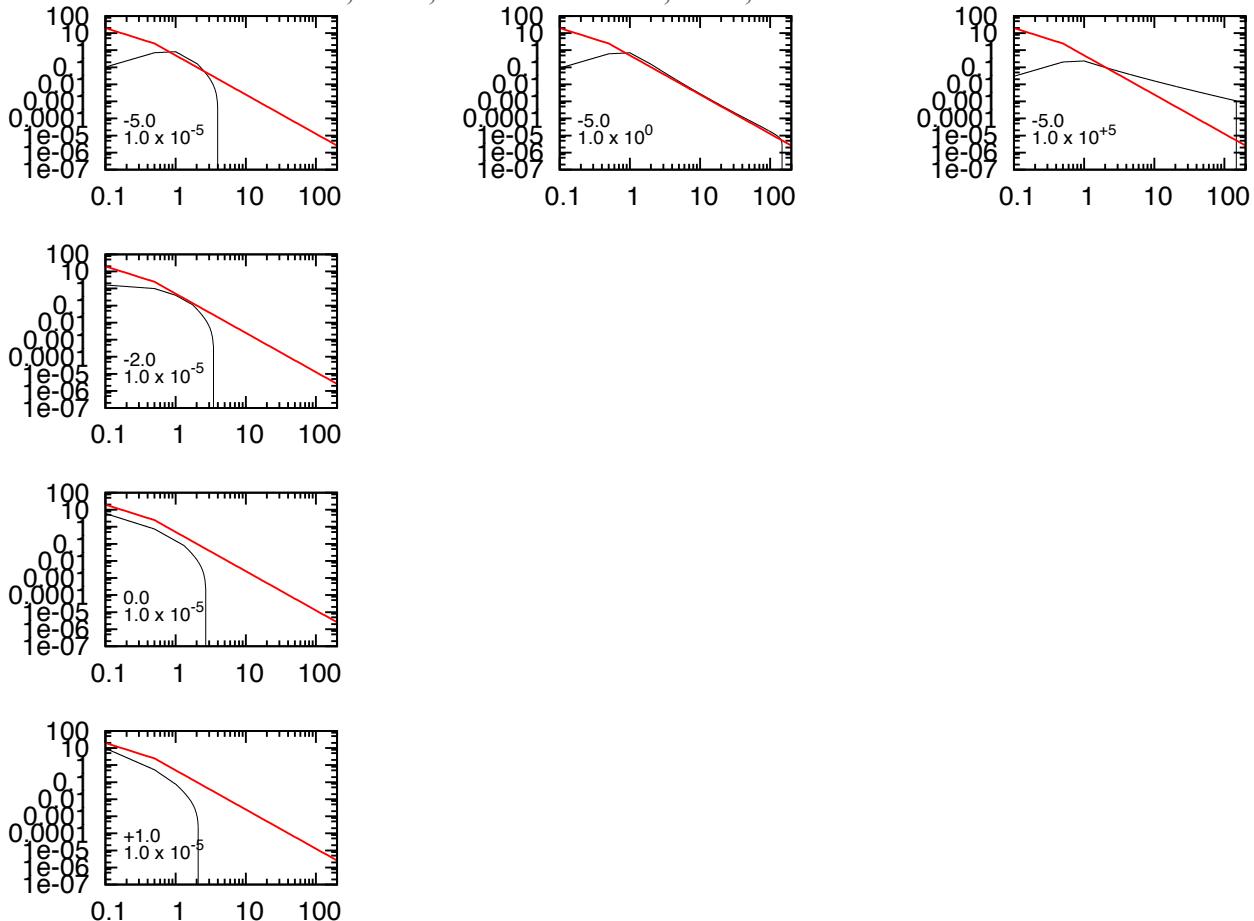
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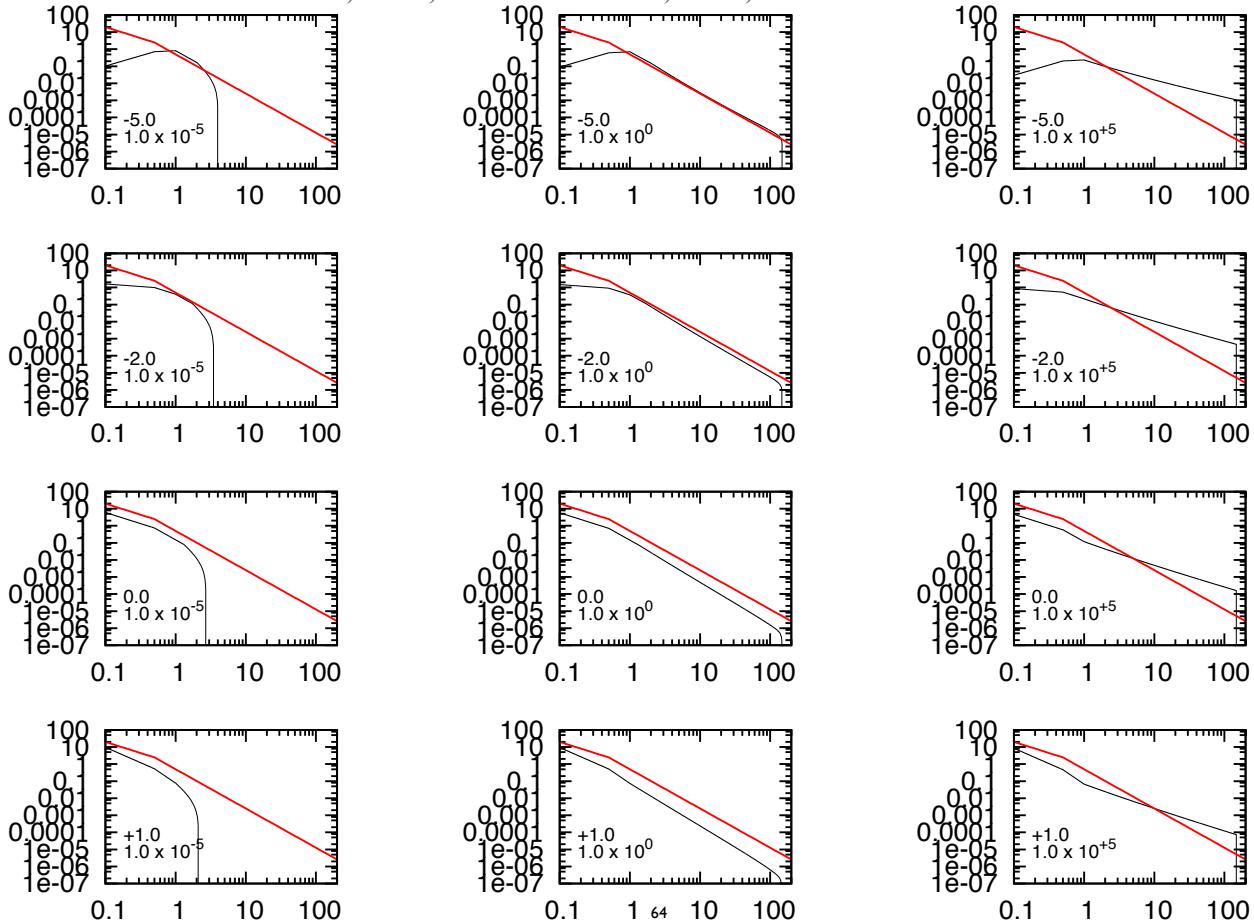
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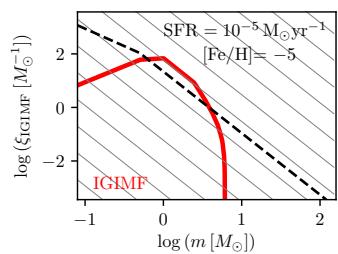
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# How would an elliptical galaxy evolve in this grid?

Jerabkova et al., A&A, 2018

1. low [Fe/H] and start of SFR burst

top-light IMF at the very beginning



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slide courtesy to Jerabkova

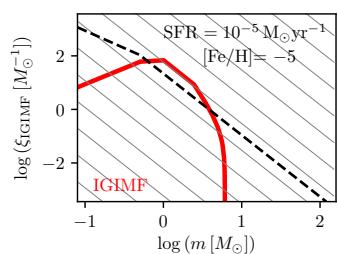
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Jerabkova et al., A&A, 2018

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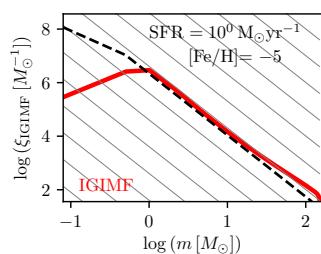
top-light IMF at the very beginning



68

2. low [Fe/H] and SFR goes up

MW IMF

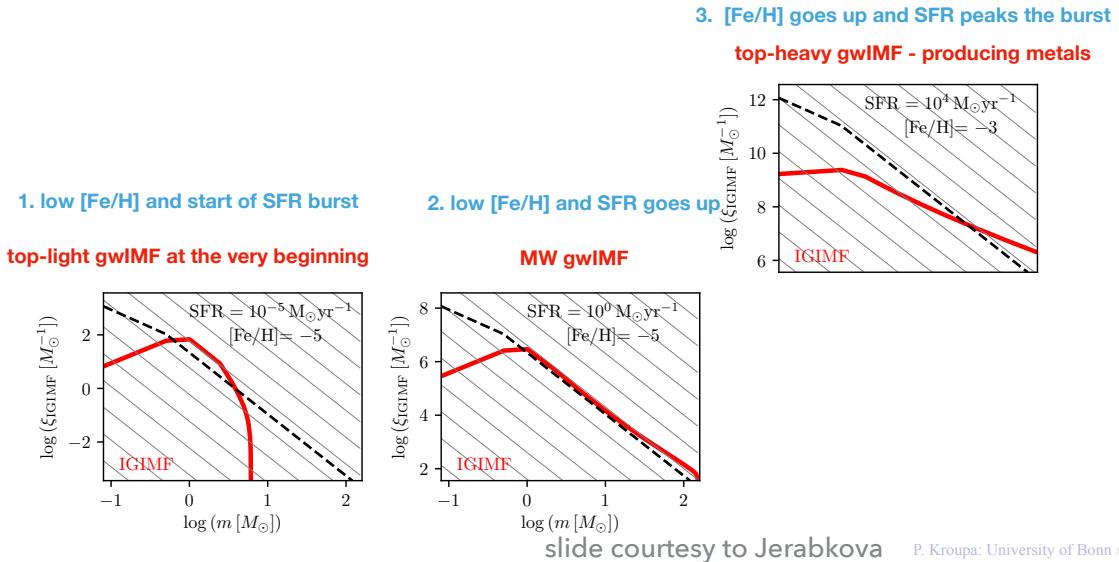


slide courtesy to Jerabkova

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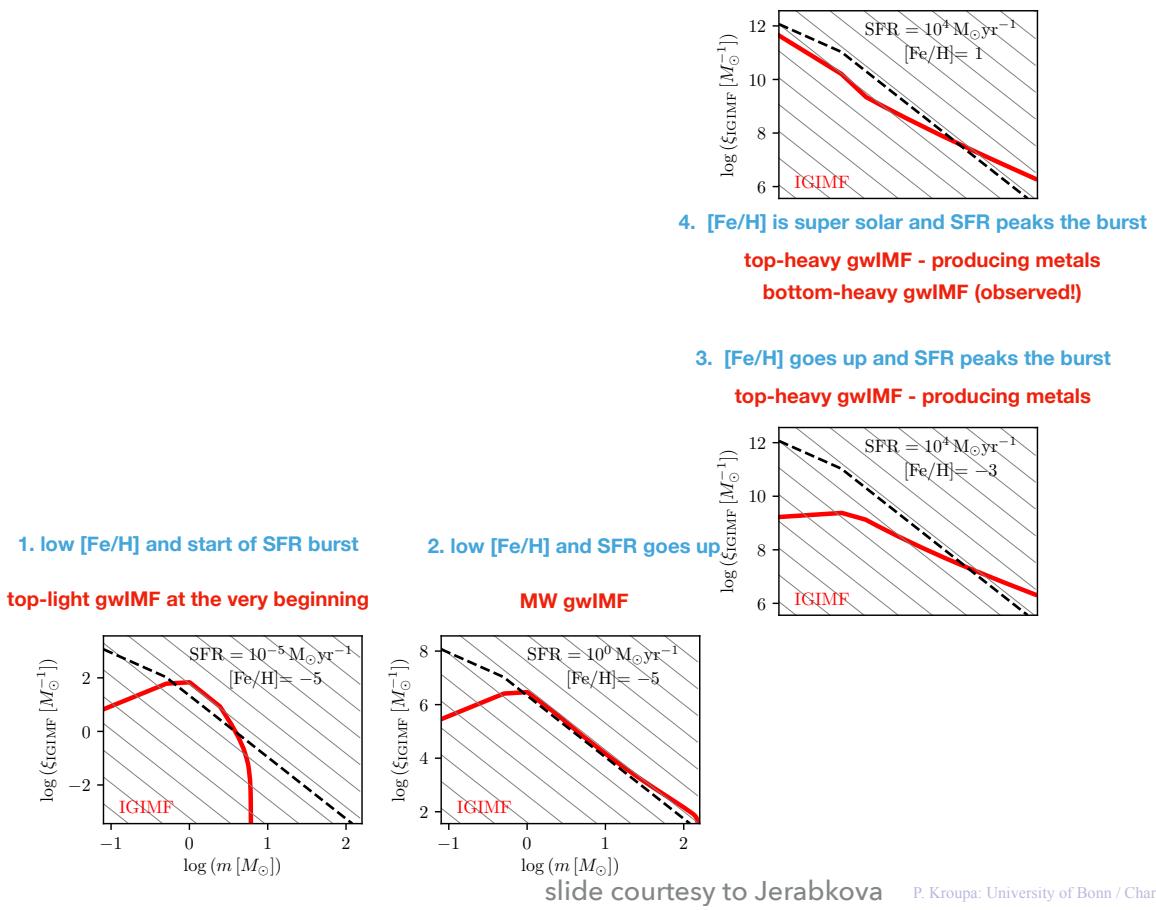
# How would an elliptical galaxy evolve in this grid?

Jerabkova et al., A&A, 2018



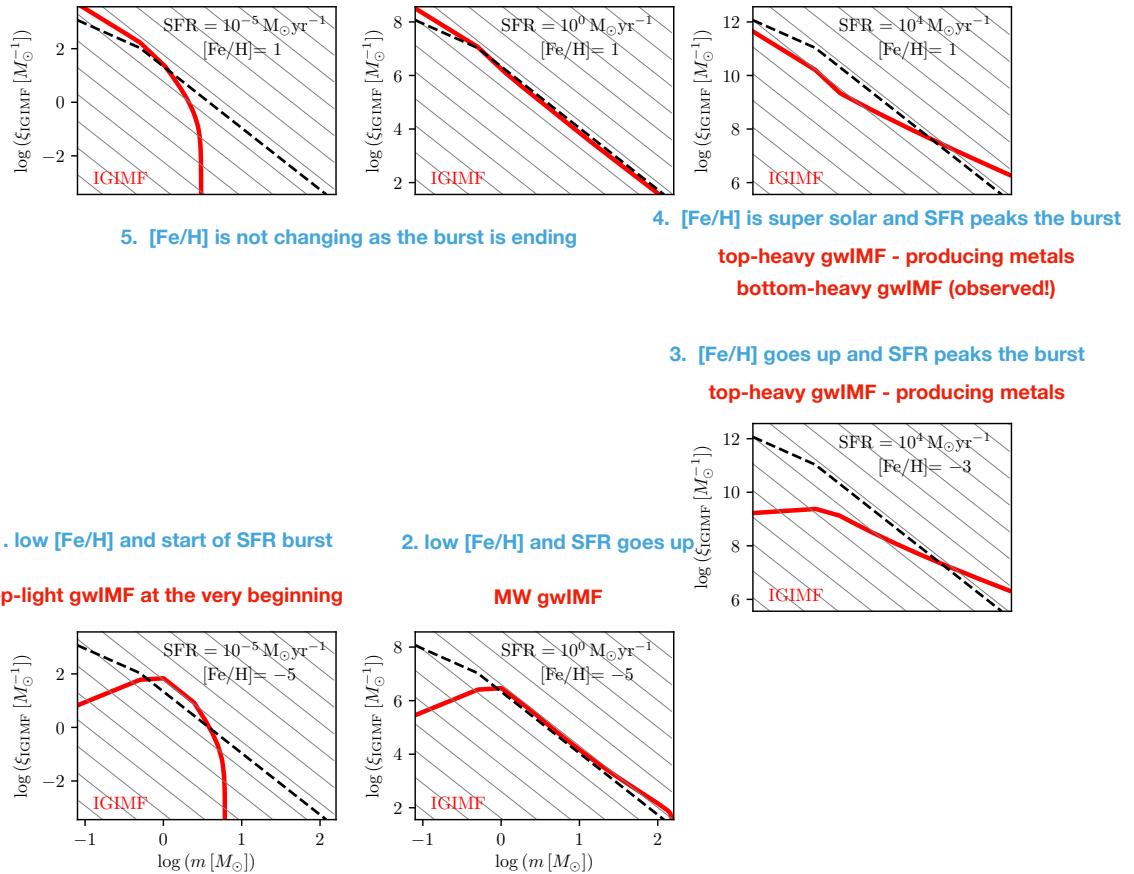
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slide courtesy to Jerabkova

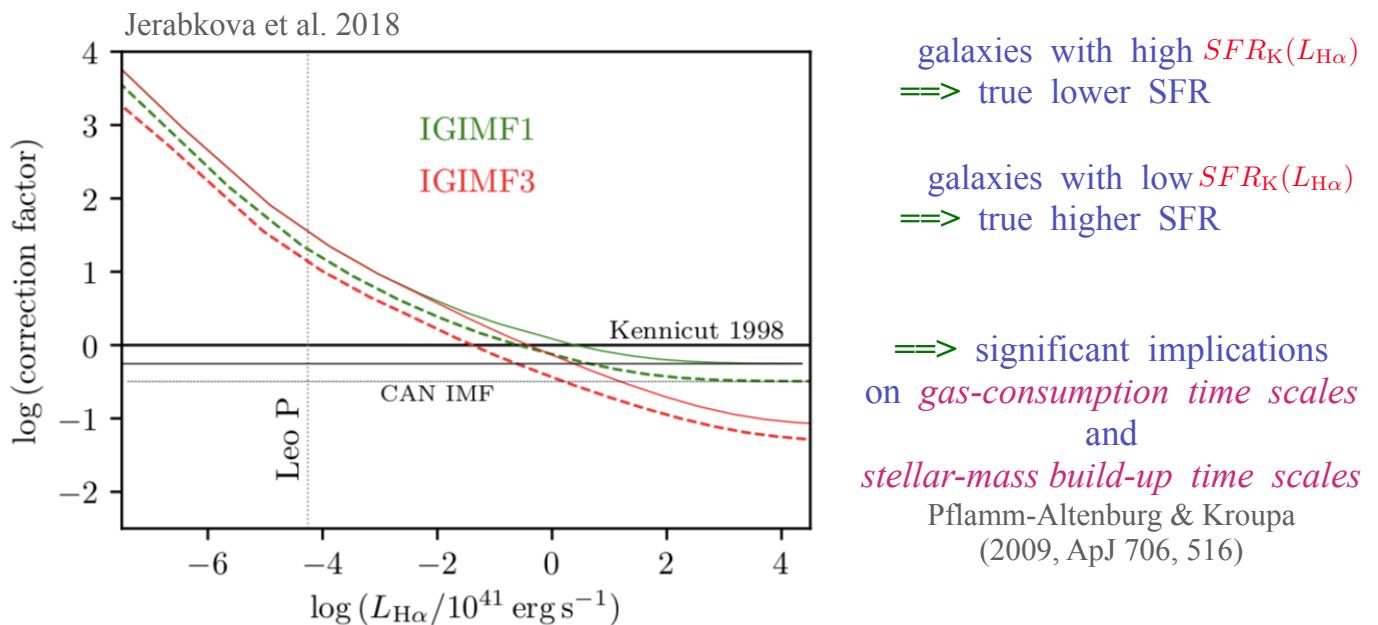
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SFRs

## Observed SFRs now depend on the SFR

$$\text{correction factor} = \frac{SFR_{\text{IGIMF}}(L_{\text{H}\alpha})}{SFR_{\text{K}}(L_{\text{H}\alpha})}$$

"Kennicutt" - measured assuming invariant standard IMF



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Rates of SNIa  
and  
neutron star mergers

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The shape of the IMF and thus of the gwIMF changes systematically.

Thus the rate of SN1a per long-lived star needs to be recalibrated, given the observed rate.

This has been achieved for the first time correctly in

## Downsizing revised: Star formation timescales for elliptical galaxies with an environment-dependent IMF and a number of SNIa

Show affiliations

Yan, Zhiqiang  ; Jeřábková, Tereza ; Kroupa, Pavel

2021A&A...655A..19Y

Number of SN1a per star must increase with SFR.

Consistent with independent observational data :

**Table 1.** Observational estimations on the time-integrated number of SNIa per stellar mass formed.

$N_{\text{Ia}} \cdot 10^3 [M_{\odot}^{-1}]$	Reference
Volumetric and field galaxies	
2.3 ± 0.6	Maoz et al. (2011)
1.3 ± 0.15	Maoz et al. (2012)
0.485 ± 0.065	Perrett et al. (2012)
0.43 <sup>+0.04</sup> <sub>-0.13</sub>	Graur & Maoz (2013)
0.98 <sup>+0.3</sup> <sub>-0.76</sub>	Rodney et al. (2014)
1.3 ± 0.1	Maoz & Graur (2017)
1.6 ± 0.3	Maoz & Graur (2017)
2 ± 1	Heringer et al. (2019)
4 <sup>+2</sup> <sub>-1</sub>	Heringer et al. (2019)
Galaxy clusters	
4.4 <sup>+1.5</sup> <sub>-1</sub>	Maoz et al. (2010)
5.4 ± 2.3	Maoz & Graur (2017)
3.1 <sup>+1.1</sup> <sub>-1.0</sub>	Freundlich & Maoz (2021)

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: University

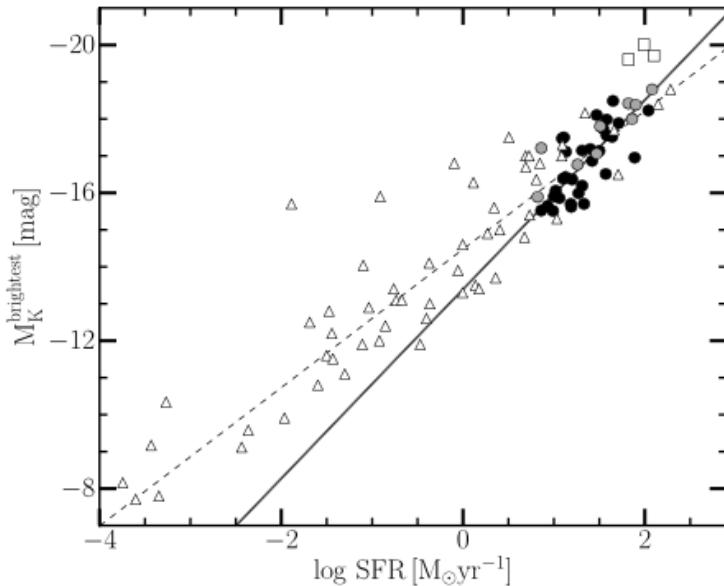
The shape of the IMF and thus of the gwIMF changes systematically.

Probably similar for neutron-star -- neutron-star mergers

Number of them per star increases with SFR ?

The process which enhances the population of tight WD--WD and NS--NS binaries is dynamical processing in high-density massive long-lived star clusters

Larger SFR ==> more massive clusters :



**Figure 4.**  $M_K^{\text{brightest}}$ –SFR relation with data from literature (the triangles; Adamo et al. 2011 and references therein) added to the present work (symbols as in Figure 1). The solid line is our best fit of Equation (3) and the dashed line is the fit from Weidner et al. (2004) to the optical V-band data after a constant  $V - K = 2$  conversion.

Weidner et al. (2004);  
Randriamanakoto,  
Escala et al. (2013)

A good correlation of the most luminous cluster and the number of SSCs with  $M_K$  magnitude brighter than  $-15$  shows that a size-of-sample effect is broadly consistent with the observed  $M_K^{\text{brightest}}$ –SFR relation. On the other hand, the observed scatter in the relation is surprisingly small, and we show that it can be explained with random sampling effects *only if* the LF of SSCs is very steep at the bright end, steeper than usually observed. Hence, physical reasons determining the luminosity of the brightest SSC from host properties, and/or internal cluster effects, likely play a role as well. We derived a relation tying

# Top-heavy IMF and "quasars"

# Consider a UCD with some $10^7 M_\odot$ today

Initially, it may have had:

- A mass of some  $10^8 M_\odot$ ...
- ... but a half-mass radius of only a few pc! (expansion through mass-loss!)
- ... A population of  $10^6$  O-stars...
- ...with a total luminosity of  $>10^{11} L_\odot$ .

====> Quasar-like objects Jerabkova et al. 2017

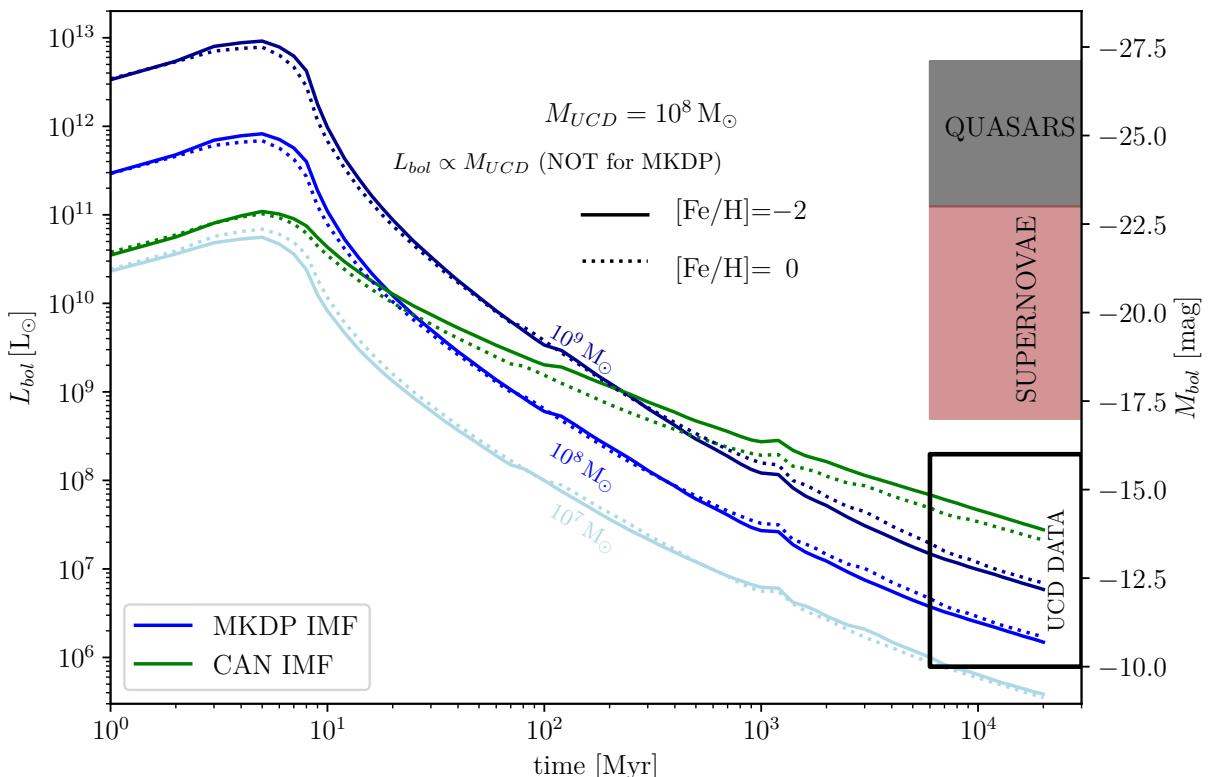
luminosities of  $10^{12} L_{\text{Sun}}$ ,  
few-pc radius,  
fluctuating on a monthly time-scale,  
life-time 50 Myr.

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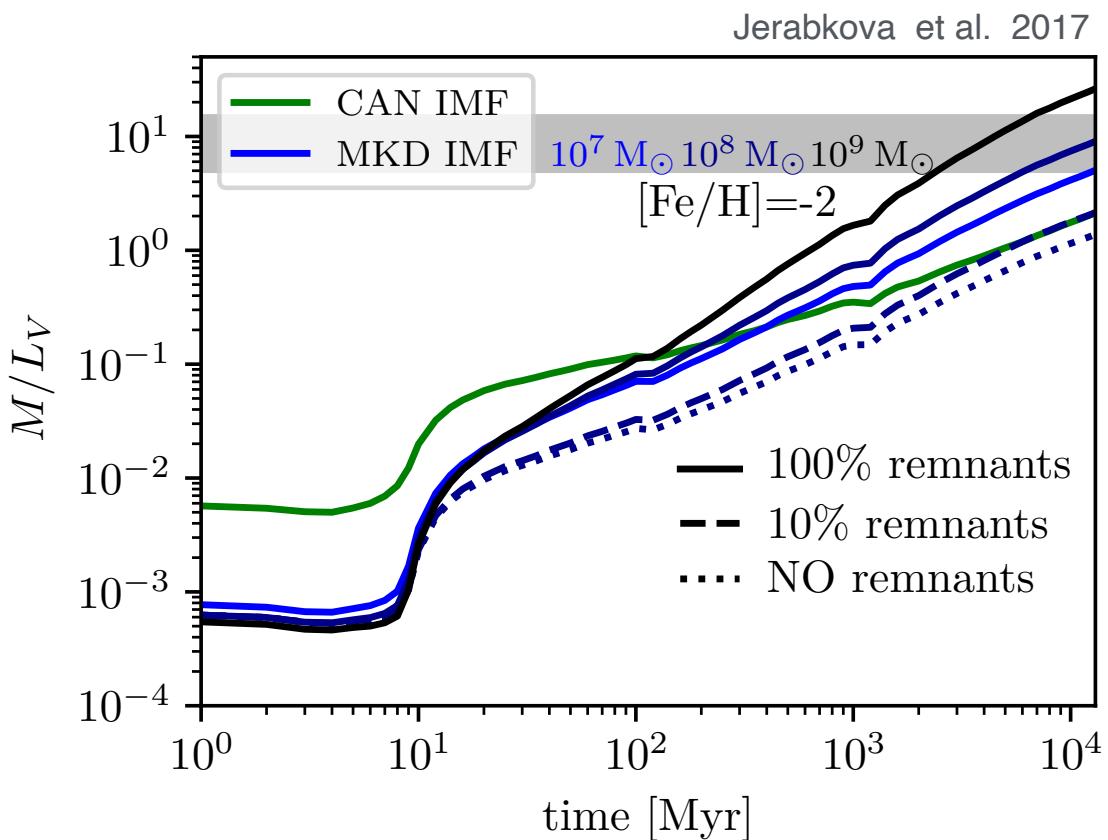
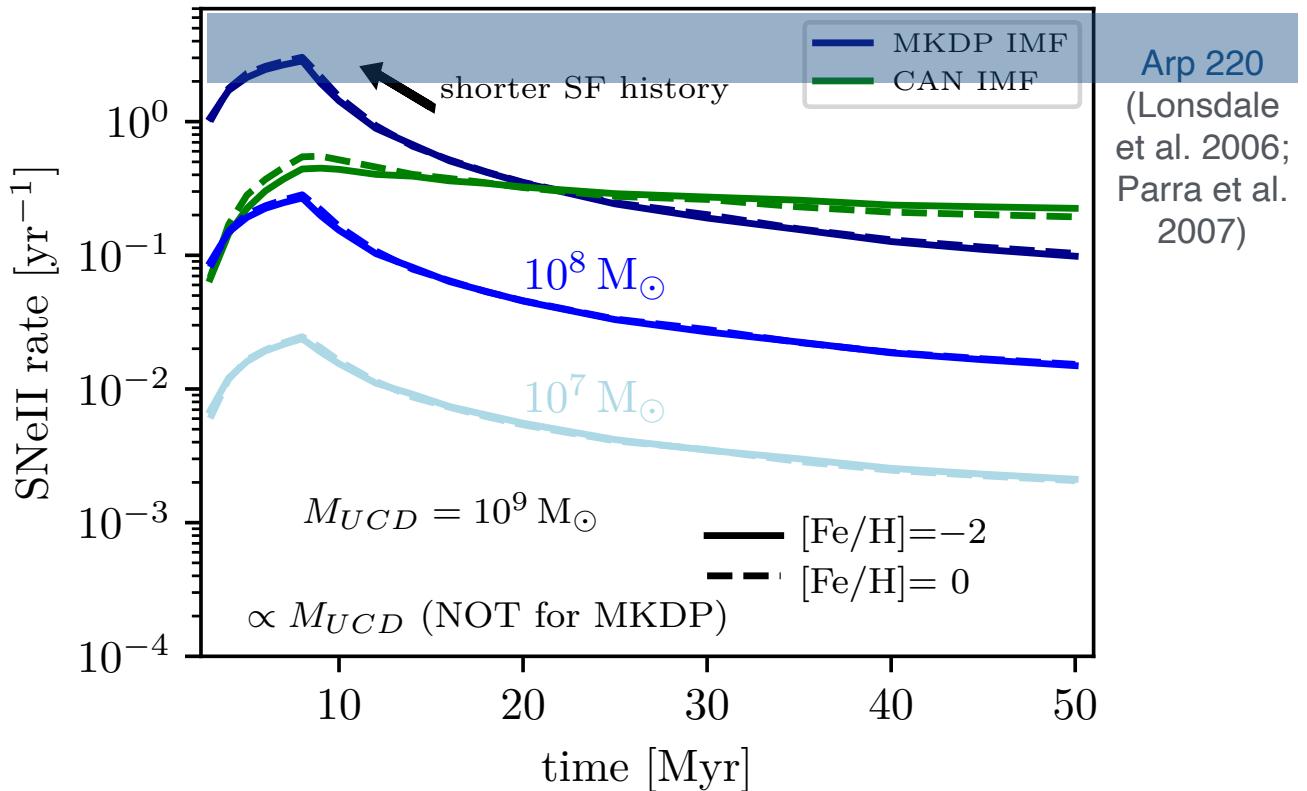
====> Quasar-like objects Jerabkova et al. 2017

The redshift dependent photometric properties are calculated as predictions for *James Webb* Space Telescope (JWST) observations.



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## Conclusions

Significant evidence for density and Z dependent IMF in molecular cloud cores (pc-scale, one Myr-time scale).

The galaxy-wide IMF (the gwIMF) changes with SFR, as expected. IGIMF theory agrees with observations.

Significant implication for measured SFRs (and gas-consumption- and stellar-mass build-up times) !

Testbed: extremely star-bursting clusters (UCDs) at high-z.

**Are some/most quasars at very high z merely young UCDs ?**  
Jerabkova et al. 2017

Formation of SMBHs      Kroupa et al. 2020  
Lecture tomorrow