Comprehensive test of the Brink-Axel hypothesis in the energy range of the pygmy dipole resonance

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Briefly about the gamma-ray strength function

- ► Gamma-ray strength function (GSF) is an average property of excited nuclei (by analogy with the level density).
- ▶ By the general definition, the GSF $f_{XL}(E_{\gamma})$ is the distribution of the average reduced width $\overline{\Gamma}$ for transitions of XL multipole type (X = E for electric, X = M for magnetic) over gamma-ray energies E_{γ} :

$$f_{i \to f, XL}^J(E_\gamma) = \frac{\bar{\Gamma}_{i \to f, XL}^J}{E_\gamma^{2L+1}} \cdot \rho^J(E_f)$$

where we consider transition(s) from the state(s) i to the state(s) f of a certain spin and parity J, E_{γ} is the energy of an absorbed/emitted photon, $\rho^{J}(E_{f})$ is the density of levels around excitation energy E_{f} .

- *i.e.* proportional to the reduced matrix element squared per unit excitation energy interval → transition probability.
- GSF from the photoabsorption experiments "upward " \overrightarrow{f} strength, from the γ -emission "downward " \overleftarrow{f} strength.



Schematic representation of the GSF in terms of γ -transitions.



Generalized Brink-Axel hypothesis

- Photoabsorption cross section (and, therefore, GSF) of the giant electric dipole resonance (GDR) is independent of the detailed structure of the initial state (D. M. Brink).
- Further generalized to both γ-absorption and emission processes (P. Axel).
- Additional independence of the initial and final spins of the states involved (only dipole selection rules).

Generalized Brink-Axel hypothesis

► The "upward " \overrightarrow{f} GSF should be identical to the "downward " \overleftarrow{f} GSF.



Schematic representation of the GSF in terms of γ -transitions.



Motivation: why is the Brink-Axel hypothesis important?

- ▶ Provides significant simplification of many problems in nuclear physics:
 - \triangleright Widely used to calculate the low-energy E1 and M1 strengths (e.g. Oslo method).
 - \triangleright Calculation of Gamow-Teller and Fermi transition strengths in β -decay and electron capture.
 - \triangleright (n, γ) cross-section calculations for the astrophysical r-process in extreme environments and nuclear reactors.

The validity of the Brink-Axel hypothesis is still to be tested for different energy regions and systems!



Comparison of the GSF from the photoabsorption cross section data f^{σ} and photon emission process $\langle f^p \rangle$. J. Isaak *et al.*, Physics Letters B 788 (2019) 225 - 230



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Motivation: why PDR?

- ▶ The PDR has another particular astrophysical effect:
 - ▷ Influence of the PDR on neutron capture rates and resulting abundances in the r-process.
 - ▷ The r-process is responsible for the production of $\approx 50\%$ of elements heavier than iron in the Universe.
 - \triangleright Appearance of the PDR increases probability of the (n, γ) reaction.







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BA hypothesis

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Motivation: why PDR?



GSF function for ¹²⁰Sn extracted from the Oslo-type of experiment $(p,p'\gamma)$, (p,p') and (γ, n) experiments in the Oslo-type of experiment $(p,p'\gamma)$, (p,p') and (γ, n) experiments of $(p,p'\gamma)$.

BA hypothesis

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Motivation: why Sn?



Comparison of the predictions for the total γ -strength functions with the OCL experimental measurements for $^{116-119,121,122}$ Sn.

H. K. Toft et al., Phys. Rev. C 83, 044320 (2011)



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Experimental techniques used in the present work

The Brink-Axel hypothesis is tested in the energy range below the neutron separation energy for 116,120,124 Sn.





Oslo method: Experiment on ^{116,120,124}Sn at the Oslo Cyclotron Laboratory



The principal scheme of the experiment.

Principal goal of the experiment is to extract:

- 1. Nuclear level density: number of levels per unit energy
- 2. γ -ray strength function ~ reduced transition probability

^{120,124}**Sn:**

- Performed in February-March, 2019 by means of OCL facilities with the OSCAR LaBr3(Ce)γ-detector array and charged particle SiRi detector.
- ▶ 126° - 140° angles are covered.
- Study of the (p,p'γ) reaction with 16 MeV proton beam on highly pure ^{120,124}Sn target.
- ▶ proton- γ coincidences were extracted.

¹¹⁶**Sn:**

Older experiment with the previous setup: NaI scintillator detectors+silicon particle telescope, ¹¹⁷Sn(³He, αγ)¹¹⁶Sn reaction with 38 MeV ³He beam.



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The Oslo method



- Unfolding: extraction of the original spectra from the raw $p-\gamma$ coincidence data on the base of the known detector response.
- ► First generation method: subtraction of secondary and higher order γ -transitions.
 - Factorization of the first generation matrix P as $P(E_x, E_\gamma) \sim \rho(E_x - E_\gamma) \cdot \mathcal{T}(E_\gamma).$
 - Subsequent extraction of the level density $\rho(E_x E_\gamma)$ and radiative transmission coefficient $\mathcal{T}(E_\gamma)$ (Brink-Axel hypothesis).
- $\mathcal{T}(E_{\gamma})$ is used for estimation of the dipole γ -strength function $f(E_{\gamma}) = \mathcal{T}(E_{\gamma})/2\pi E_{\gamma}^3$
 - ▶ Normalization of the level density and γ RSF.



Extraction of the GSF for a certain initial excitation energy



Since the first generation matrix P is given by $P(E_x, E_\gamma) \sim \rho(E_x - E_\gamma) \cdot \mathcal{T}(E_\gamma)$, it is possible to map $\mathcal{T}(E_x, E_\gamma)$ and therefore $f(E_x, E_\gamma)$ as functions of initial excitation energy E_i :

$$f(E_i, E_{\gamma}) = \frac{N(E_i) \cdot P(E_i, E_{\gamma})}{\rho(E_i - E_{\gamma})},$$

where $N(E_i)$ is the normalization factor:

$$N(E_i) = \frac{\int_0^{E_i} \mathcal{T}(E_\gamma) \rho(E_i - E_\gamma) dE_\gamma}{\int_0^{E_i} P(E_i, E_\gamma) dE_\gamma}$$



The Shape method



▶ The number of counts in each diagonal is given by:

$$N_D \sim f(E_\gamma) E_\gamma^3 \sum_{J_i = J_f + 1}^{J_i = J_f + 1} \sigma(E_i, J_i) g(E_i, J_i),$$

where $\sigma(E_i, J_i)$ is the cross-section to populate a certain initial state with spin J_i at a given excitation energy E_i , and $g(E_i, J_i)$ is a known spin distribution function.

▶ For the current analysis the g.s. $(J_f^P = 0^+)$ and 1st excited state diagonal $(J_f^P = 2^+)$ are chosen (for the same E_i):

$$f(E_{\gamma 1,2}) \frac{N_{D1,2}}{E_{\gamma 1,2}^3 \sum_{J_i = J_f + 1}^{J_i = J_f + 1}} g(E_i, J_i)$$

▶ Pairs of $f(E_{\gamma 1})$ and $f(E_{\gamma 2})$ for each excitation energy are "sewed "together:



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(p,p') experiment at the RCNP



► Total gamma strength functions for stable tin isotopes were also extracted from the forward-angle inelastic proton scattering data at relativistic beam energies (295 MeV protons). Eperiment was performed at the Research Center for Nuclear Physics (RCNP) in Osaka, Japan. [S. Bassauer et al., Phys. Rev.C 102, 034327 (2020).]





Main results: GSF for 116,120,124 Sn





Main results: GSF extracted for several initial excitation energies for $^{120}\mathbf{Sn}$





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Main results: GSF extracted for several initial and final excitation energies for $^{120}{\rm Sn}$





Main results: Oslo method + Shape method + (p,p') for ¹²⁰Sn





Main results: Oslo method + Shape method + (p,p') for ¹²⁴Sn





Summary

- ► Comparison of the GSF extracted from the Oslo-type experiments and the relativistic Coulomb excitation experiment in forward-angle inelastic proton scattering demonstrates good agreement in the energy range between ≈ 5.5 MeV and the neutron separation energy.
- Experiments based on ground state photoabsorption provide the same information on GSFs in nuclei as Oslo-type experiments.
- ▶ The GSFs seem to be somewhat independent of the energies and spins of initial and final states (as shown by the Shape method).

the generalized Brink-Axel hypothesis holds for ^{116,120,124}Sn isotopes in the energy range below the neutron separation energy!

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The assumptions made in the calculations of (n,γ) reactions relevant to r-process nucleosynthesis are verified.



Thank you for your attention!



Porter-Thomas fluctuations





Porter-Thomas fluctuations





Porter-Thomas fluctuations





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