



Summary of "Physics of Photocathodes for Photoinjectors" workshop BNL, October 3-5, 2023

Luca Cultrera

EWPPA – workshop – HZDR, September 17-19, 2024





Special Thanks

Local Organizing Committee

- L. Cultrera (co-Chair)
- M. Gaowei (co-Chair)
- S. Capp

Support staff

- C. Weaver
- D. Basozmen

Scientific Committee

- Luca Cultrera
- Oksana Chubenko
- Dimitre Dimitrov
- Mengjia Gaowei
- Carlos Hernandez-Garcia
- Kevin Jensen
- Siddharth Karkare
- Jared Maxson
- Nathan Moody
- Pietro Musumeci
- John Smedley
- Marcy Stutzman
- Theodore Vecchione



P3 workshop 2023, BNL



75 Registered participants



7 sessions, 37 talks



15 posters



P3 workshop 2023, BNL





7 plenary sessions

- Introductory session
- High charge High average current
- Free Electron Laser (two sessions)
- Theory (two sessions)
- Novel concepts (two sessions)
- Ultrafast Diffraction and microscopy
- Spin polarized sources

I will try to give an objective and comprehensive summary

Talks are uploaded to the following indico website

Photocathode Physics for Photoinjectors Workshop (3-October 5, 2023) · Indico (bnl.gov)



Introductory session



EWPAA 2022 Highlights

Lee Jones (STFC, DL), On behalf of the EWPAA Organizing Committee



Accelerator R&D Roadmap and it's impact on photocathode research

- New accelerator roadmap emphasise interest to further development of Energy Recovery Linacs (ERL) as high efficiency sustainable accelerator facilities.
- High average current, in 100 mA range, electron injectors are one of the most critical components of ERLs and additional efforts should be concentrated on further development of the injectors.
- Major efforts should be concentrated on the following:
 - Optimisation technology of existing photocathode materials
 - · Development new photocathode materials which could operate with existing laser system
 - Design photocathode plug which may be compatible with different types of the guns
- Development of laser systems which would allow for delivery required current with UV photocathodes
- Development of gun technologies which would allow for providing beams compatible with existing accelerator structures without excessive bunching schemes. Potential to reach this goal have.
 - High field QWR SRF gun
 - High field QWR NCRF gun
 - Elliptical cavity SRF gun
- There is no operational injector which can demonstrate 100 mA of polarised current: Limited by photocathode lifetime
- Potential to deliver this current (if GaAs lifetime by activation with Cs-alkali metal layer is successful),
- DC photocathode guns
- · SRF photocathode guns



Electron Sources for Particle Accelerators

YINE SUN Advanced Photon Source Argonne National Lab.

Oct. 3, 2023

Photocathode Physics for Photoinjectors Workshop Stony Brook, New York

Argonne 📤

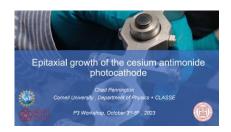
April 2023: e[±] Source Roadmap Working Group Report

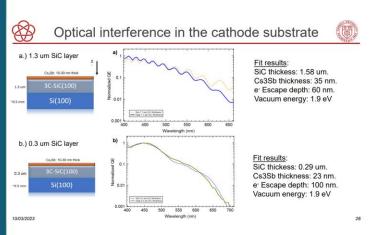
https://indico.fnal.gov/event/59123/

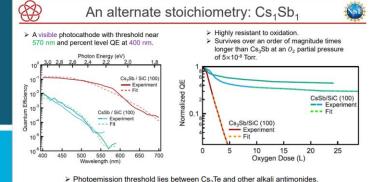
Year	Near-term (<5 years)	Mid-term (5~10 years)	Long-term (10~20 vears)
e ⁻ Cathode	Photocathodes with 1% QE and 30 n	Cryogenic temperatures and athodes in DC guns (20mA pol. and 100 mA unpol neV MTEs Photocathanising photocathodes (robust surfaces, nano-struc	very high fields operation .) odes with 1% QE and 5 meV MTEs
e ⁻ Gun	DC gun beam ~1-10 mA polarized NCRF: cryo gun at 250 MV/m; x-band gun, CW Polarized GaAs in an SRF photogun	10 ⁻¹⁴ Torr vacuum for long GaAs lifetime and Low Frequency rf gun	DC gun beam 10~20 mA polarized
e ⁻ Injector	Control laser profile, limit nonlinear SC induced emittance growth: beer can (mid); elliptical (far) NCRF, SRF accelerating cavities: fully RF symmetrized fields to eliminate emittance growth to 10% (near), 1%(mid), 0.1%(far) Partition phase space: RFBT-EEX for damping ring free (mid), linear LPS (long) High Charge Drive Bunch Trains: charge-balanced, equal energy bunches duration 5-25 nsec.		
e ⁺ polarized	SC undulators Compton-based sources - h Bremmstrahlung polarized positron source devi	igh flux circularly polarized gamma-rays R&D	class polarized e+ source
e ⁺ unpolarized		d acceleration sections ources for accelerator and ultrafast science (also	polarized)
	10 ⁻¹⁴ Torr vacuum for long GaAs lifetime Photocathode Physic	Routine 10's mA GaAs beams cs for Photoinjectors Workshop Yine	25

10/3/2023

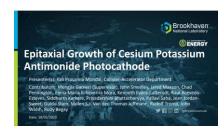
High Charge / High Average Current

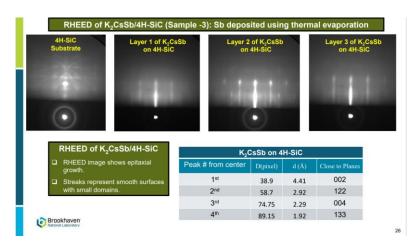


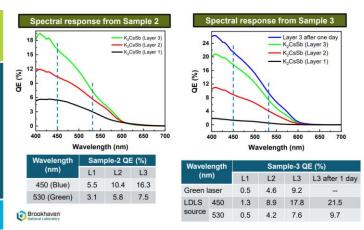




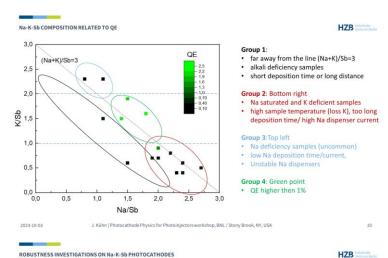
- C.T. Parzyck, C.A. Pennington et al, "Atomically smooth films of CsSb: a chemically robust visible light photocathode" arXiv:2305.19553

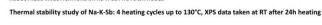


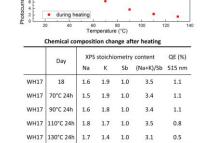












- · Photocurrent decreases with higher sample temperature
- Decrease of photocurrent may due to the loss of K content
- · X-ray also contribute to the loss of

J. Kühn | Photocathode Physics for Photoinjectors workshop, BNL / Stony Brook, NY, USA

High Charge / High Average Current





Demonstration of Thermal Limit to Mean Transverse Energy from Cesium Antimonide Photocathode Arizona State University

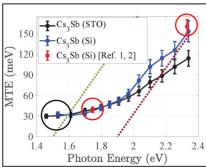
Experiment Measurements Growth ħω: 1.45 - 2.33 eV Thickness: ~40 nm Substrates: Si and Nb doped ST Repetition rate: 500 kHz Pulse length: 150 fs Transfer QE Via UHV transfer line in PEEM

Mean Transverse Energy

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Measured: PEES, MTE & QE



- At $\hbar\omega = 1.5$ eV. MTE = 30 meV (~25 meV at 300 K)
- At $\hbar\omega$ = 1.8 eV. MTE = 40 meV and at $\hbar\omega$ = 2.3 eV, MTE = 150 meV (comparable to previously reported
- The dotted line is the plot for (excess energy)/3 considering Φ = 1.5 eV (green) and Φ = 1.9 eV (brown)
- MTE doesn't scale as 1/3rd of excess energy (Scattering before emission)

Physical Review Special Topics-Accelerators and Beams 18.11 (2015): 11340: 2) Applied Physics Letters 99.15 (2011).

3% - 4%

 $\lambda = 530 \text{ nm}$

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(Hexagonal)

Comparison between the reference experimental [1] and

cubic compounds.

However, the observed two major peaks in our KSb-thick

cathodes do not match those of the K₃Sb hexagonal or





Comparison between the DFT Simulation and Experimental Data for K₃Sb

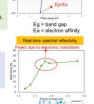


Motorized filter wheel

(Cubic)







2) Photocathodes production To support 24/7 operations, cathode production and exchange systems were developed which

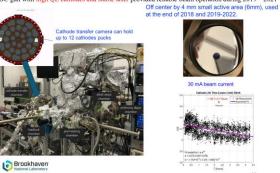
include two cathode deposition systems, three multi-cathode (up to 12 cathodes) vacuum suites and a mechanism allowing for cathode exchange in RHIC tunnel in about 1 hour.

DC gun with high QE cathodes and stable laser provided reliable beam operation during 2019 - 2021

LEReC high current experience

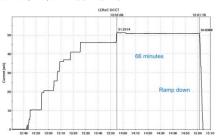
iaofeng Gu, Alexei Fedotov, Dmitry Kayran, Jore Brookhaven National Laboratory

with alkali cathode



3) Higher Current Test

- 1. 320kV: if the HVPS Faults are not caused by the voltage, could try 350kV later.
- 66 minutes: The current was ramped down prior to next polarization measurements in RHIC which typically results in high losses in the location of the Gun.
- 50 mA: limited by the injection dump power (25kW);







Brookhaven

Free Electron laser

SHINE (14xt

R&D of Very High Frequency (VHF) gun for SHINE project at Tsinghua University

Yingchao Du

Department of Engineering Physics, Tsinghua University Beijing, China,100084

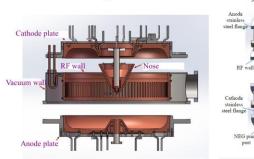
dych@tsinghua.edu.cn

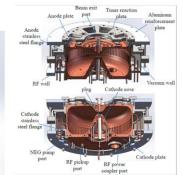


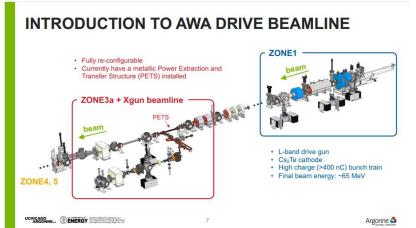


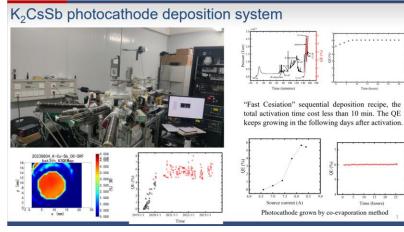
WHF gun RF design and optimization

Mechanical design









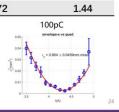
) ドギ大学 Beam testing

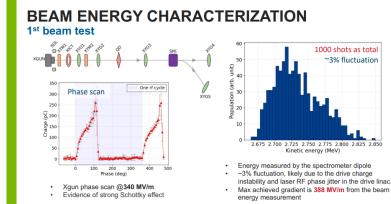
• Emittance measurement and optimization(preliminary results)

Bunch charge	Projected emittance (95%) $(um \cdot rad)$	Slice emittance (95%) $(um \cdot rad)$	Bunch length (mm rms)
10 pC	0.16	0.15	0.49
50 pC	0.41	0.38	1.15
100 pC	0.85	0.72	1.44

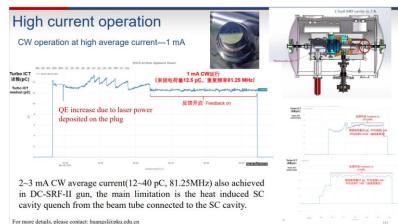




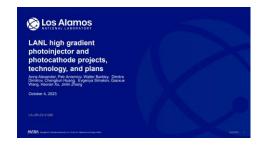




Argonne 🔷



Free Electron laser



Carie: high gradient photoinjector capability

· Goal: demonstrate operation of high QE cathodes in a high-gradient RF injector

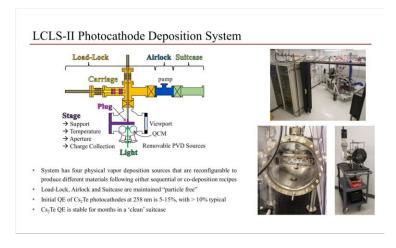
Why?

- · ICS 44 keV xray source
- UED



1_{//}s(Cu/Cu-Ag)

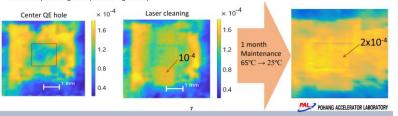
Some of the Photocathode Science at SLAC 2023 Photocathode Physics for Photoinjectors Workshop Theodore Vecchion SLAC NATIONAL ACCELERATORY





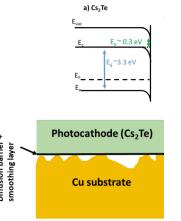
PAL-XFEL s-band photocathode gun and injector

- Maintaining the normalized emittance at 135 MeV
- Average 0.4 u at 250 pC (variations between 0.35~0.45u)
- QE at 253 nm, 1~2x10-4, Vacuum 2x10-11 mbar
- Twice laser cleaning since 2016 and 2022,
- Low surface damage using ~300 ps long IR laser cleaning
- Half a day cleaning and operation right away



High gradient photocathode design:

- High bandgap + high(er) Ea
- · Ultrathin cathodes
- Heterostructures?
- · Lossy/imperfect preferred?
- · Does photoemission facilitate breakdown?
- · What happens to the injector when photocathodes break down?



SLAC's Grand SRF Photocathode Challenge

We have $\varepsilon_n = 0.4 \,\mu\text{m}$ at 100 pC, 1 mm, 100 MeV (in theory) We want $\varepsilon_{\rm u}$ < 0.1 μm at 100 pC, 1 mm, 100 MeV

Photocathode must also not generate particles or contaminate the cavity

Intrinsic emittance

Dark current

Quantum Efficiency

Temporal Response

1/e Lifetime

 $\varepsilon_{int} \leq 0.3 \ \mu m/mm \ future$ $\varepsilon_{int} \leq 0.2 \ \mu m/mm \ eventually$

QE drops near threshold

QE drops with temperature

Operate near threshold for emission Operate at low temperature Reduce surface roughness

Increase chemical uniformity

Use semiconductor photocathode QE must be sufficient to generate 100 pC without multiphoton contributions but QE > 0.1% is not strictly a requirement

< 50 ps, bunched downstream to « 1 ps

> 1 week (operational issue)

« 1 ps at the photocathode is not necessary

Visible or IR wavelength preferred for laser shaping

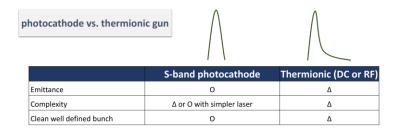
Motohiro Suyama, Hamamatsu

Challenge: Which to use?

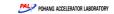
Spectral Response Curves

Conventional: Cs₃Sb More exotic: Cs,Sb:Na,KSb Novel: Na₂O

Wavelength [nm]

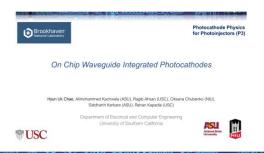


- Low emittance provides headroom for injection and readiness for the lower emittance storage ring upgrade
- Higher QE is preferred, then low power laser is easy to handle, more flexible to bunch patterns
- * Transform-limited laser pulses ensure the well-defined bunch profile

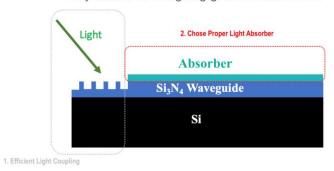




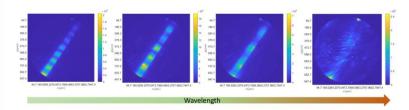
Novel concepts



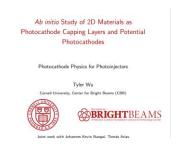
Key factors of designing good Photocathode



Wavelength Dependent Photoemission



- The period of the pattern increases as we go from $\lambda = 520$ nm to $\lambda = 532$ nm.
- These transverse patterns are formed due to interference between these co-propagating modes thereby generating beating patterns with significant evanescent intensities that cause the electron emission.



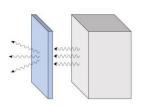
Why Scattering States are Important

Electron Transparency and Reflectivity

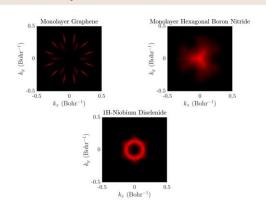
 Mean-transverse energy (MTE) and quantum efficiency (QE) are modulated by capping layers

 We need electron transparency and reflectivity

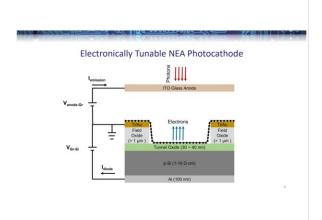
 Transmission and reflection can be taken directly from scattering states

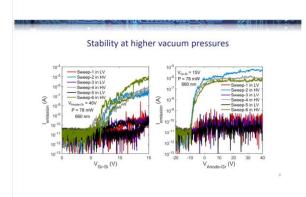


Results: Summary of Photoemission





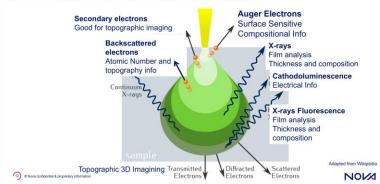




Novel concepts



Electron Beams Have Many Capabilities to Address Many of the Needs For Metrology and Inspection



E-Beam Source Requirements for Metrology

Electron Source

- ✓ Energy Spread ≤ 0.3 eV
- √ High Brightness > 10⁸A / (V m²)
- ✓ Focus to high current densities
- √ Spot Size: <100 nm²
 </p>

Multibeam > 20 >20 ✓ (or smaller for in device testing) Cost Cost ✓ Stability: <1%
</p> ✓ Long Lifetime Photocathodes have some clear advantages over state-of-the-art electron sources. 101/31

Temp

Stability

Schottky

TFE

≥ 0.7

1800

<1%

~30x107

 A/m^2

Photo-

cathode

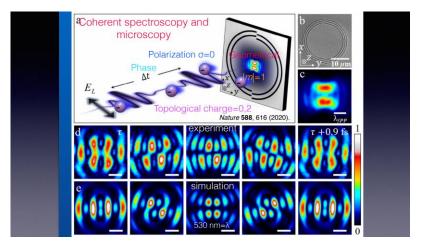
0.3

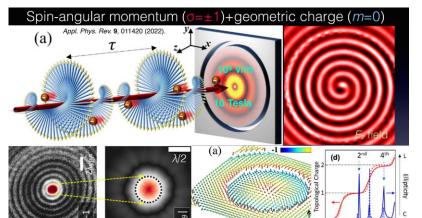
294

<1%

>30x10⁷









Optical Near Field Electron Microscopy

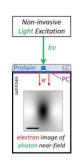
John Smedley, SLAC Guido Stam, ULEI Sense Jan van der Molen, ULEI

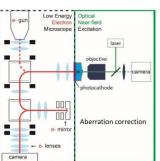




Positioning ONEM in LEEM



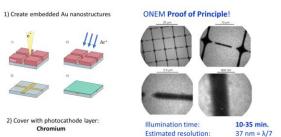






UV-ONEM using LED (275 nm)







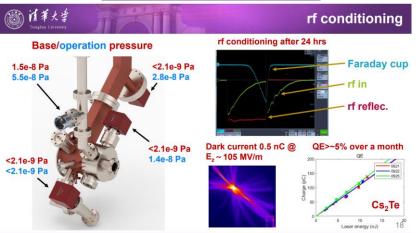
UED/UEM

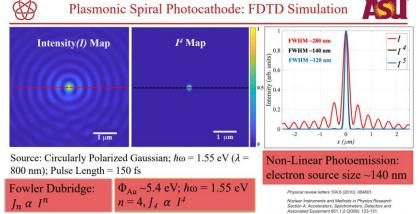




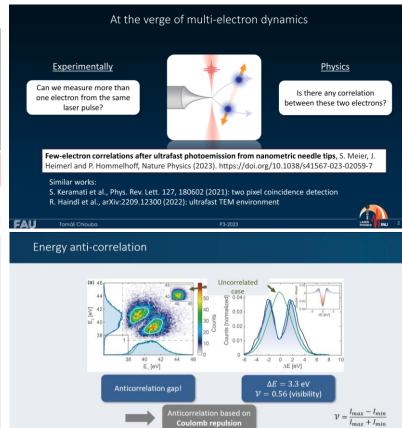
Bright Electron Beams from Plasmonic Spiral
Photocathode
Alimohammed Kachwala
Arizona State University



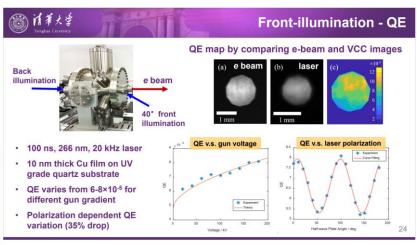


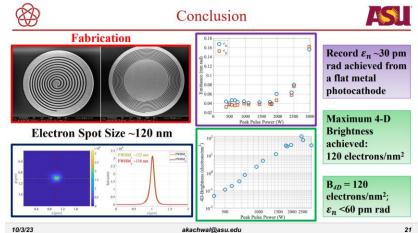


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[1] Yudin, G.L., Ivanov, M.Y., Physical Review A 64(1) (2001)





10/3/23

UED/UEM





Dr. Lee Jones Senior Accelerator Physicist

Accelerator Science and Technology Centre STFC Daresbury Laboratory

TESS: The Transverse **Energy Spread Spectrometer**

and Recent Experimental Results to

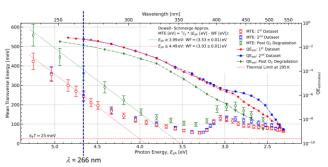
TESS Leading Particulars:

- · Laser-driven plasma broadband lightsource
- · Twin grating monochromator
- Usable flux from λ = 236 nm to > 800 nm
- · Filter wheel to avoid space charge
- · Beam profile camera at working distance Optimise source spot for each λ
- · Sapphire viewport window
- · Maximise UV transmission
- TESS base pressure typically 3 × 10⁻¹¹ mbar
- · LN, Cooling loop
- · Piezoelectric leak valve for controlled degradation studies
- TESS linked to GaAs PPF at 4 × 10⁻¹² mbar
 - . Storage carousel with up to 6 photocathodes
- · Vacuum suitcase for sample transfers between systems (< 1 × 10⁻¹⁰ mbar)

L. Jones et al., Rev. Sci. Instrum. 93, 113314 (2022)

2-Stage Microchannel Plate TESS XHV Vacuum System Kodial Window Credit: D.P. Juarez-Lopez & H.E. Scheibler

TESS Characterisation of DL CsTe Photocathode (~ 5% QE @ 266 nm)



- MTEs for both datasets broadly agree, and follow Dowell-Schmerge approximation (WF ~ 3.9 eV)
- O₂ degradation increases MTE (WF ~ 3.5 eV) and reduces estimated QE in comparison to 'clean' surface
- Thermal floor not reached, either for the dominant Cs-Te emitters or the Cs/CsO emitters
- Demonstrates the need for multi-wavelength QE monitor during deposition
 H.M. Churn et al., Proc. IPAC 2023, TUPN030, 1404-1407

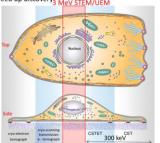
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Promise of MeV Microscopy

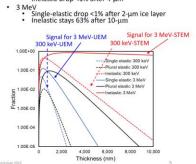
Benefit of increasing electron energy to MeV

- Life science application: 3D image intact thick bio-samples studying cell-biology and microbiology in cellular context
- · No need of cryo-FIB (focused ion beam) slicing thick cells
- · Limit to 10-20 lamellae/day, "blindly" select target
- Speed up discovery MeV STEM/UEM



Top and side views of a eukaryotic cell. S.G. Wolf, et al., Cellular Imaging, Springer

- . Why MeV? Phase contrast (TEM) & amplitude contrast (STEM)
- 300 keV
 - Single-elastic drop <1% after 1-µm ice layer
 Inelastic drop <1% after 4-µm

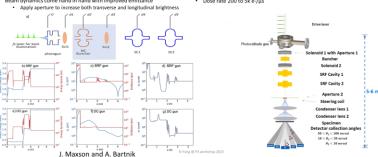


Challenges of implementing MeV STEM

- 1. Reduce beam emittance to ~ 2 pm with current of 30 750 pA
- 2. Photocathodes & beam dynamics need to improve brightness by >1000

Probe size: $\sigma_d = 2 nm$; divergence: $\sigma_\theta \le 1 mrad$; emittance: $\epsilon_{eeo} \le 2 pm$

- · Reduce laser spot size and MTE from photocathode
- Increase OE >10-5
- Beam dynamics come hand in hand with improved emittance
- Preliminarily MeV-STEM column design (reversal of TEM column)
- Assume 2 µm spot at cathode
- Dose rate 200 to 5k e-/us





Theory

Modeling Optical Interference Effects for Optimization of Electron Emission Properties from Thin Film Semiconductor Photocathodes

D. A. Dimitrov, A. Alexander, C. Huang, N. Moody, V. Paylenko. E. Simakov, G. Wang, H. Yamaguchi Los Alamos National Laboratory, Los Alamos, NM 87545, USA K. L. Jensen, University of Maryland, College Park, MD 20741, USA J. Smedley
SLAC National Accelerator Laboratory, Menlo Park, CA 94025, USA LA-UR-23-31152

Optical interference effects in a thin film photocathode

- Light E-field in photocathode: $E_{ph}(x) = E_0(t_1e^{i\kappa_1x} + r_1e^{-i\kappa_1x})$
- The E-field r; and t; coefficients depend on indices of refraction $\hat{n}_i = n_i(\omega) + ik_i(\omega)$ in different material layers and on L.
- For normal light incidence, r_i and ti are derived in K. L. Jensen et al., J. Appl. Phys. 128, 115301 (2020).

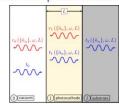


Figure 10: Schematic of incident light on a photocathode thin film of thickness L deposited on a (metal) substrate.

Absorption and transport: thin film photocathode

$$f_{tf}(\omega, E, \cos(\theta)) = \frac{\int_0^L |E_{ph}(x)|^2 e^{-x/(\lambda(E)\cos(\theta))} dx}{\int_0^L |E_{ph}(x)|^2 dx}$$

Comparison to QE experimental data (#) from Cs₃Sb on Ag and Si, for $\lambda = 450$ nm and 532 nm.

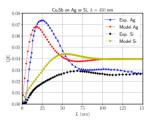


Figure 17: For $\lambda = 450 \text{ nm} (2.755 \text{ eV})$, the QE from Cs₃Sb on Ag is highest for Cs₃Sb film thickness near 21 nm.

D. A. Dimitrov et al. Slos Alomos

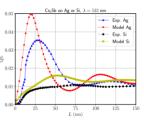
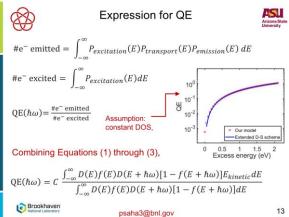
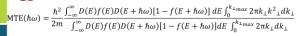


Figure 18: The extended MM for QE shows similar functional behavior with film thickness for both substrates.





Expression for MTE

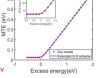


By substituting the value of $k_{\perp_{max}}$, where $k_{\perp_{max}} = \sqrt{\frac{2m}{\hbar^2}}$ ($E_{kinetic}$)

$$\text{MTE}(\hbar\omega) = \frac{\hbar^2}{2m} \frac{\int_{-\infty}^{\infty} D(E) f(E) D(E + \hbar\omega) [1 - f(E + \hbar\omega)] E_{kinetic}^2 dE}{\int_{-\infty}^{\infty} D(E) f(E) D(E + \hbar\omega) [1 - f(E + \hbar\omega)] E_{kinetic} dE}$$

Under the assumption of a constant DOS



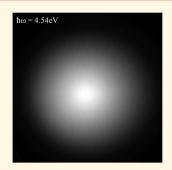


Electron Scattering Processes in Photoemission: The Franck-Condon Effect W. Andreas Schroeder Louis Angeloni and Ir-Jene Shan Physics Department, University of Illinois Chicago Department of Energy

Metallized diamond(001) emission

Photocathode Physics for Photoinjectors (P3) Workshop





Two emission signals:

(i) Outer large MTE signal from strong optical phonon assisted Franck-Condon effect

$$\mathbf{p}_{vac.} = \mathbf{p}_{band} + n\mathbf{q}_{phonon}$$

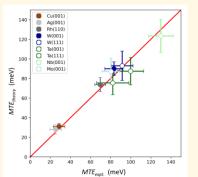
- Emission of Boltzmann tail of electron distribution from NEA upper CB of diamond ($\gamma \approx -1eV$)
- High emission efficiency over PE barrier due to momentum resonance $(T \rightarrow 1)$
- (ii) Inner signal with MTE < 100meV from electrons directly emitted from NEA upper CB of diamond
- Lower emission efficiency as p_{vac.} ≠ p_{band}; less than ~5% of total signal

BOTH with Gaussian spatial beam profiles

J.D. Rameau et al., Phys. Rev. Lett. 106, 137602 (2011)

MTE: Theory vs. Experiment





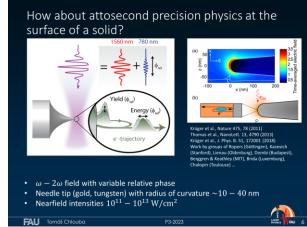
- Theoretical MTE:
- $MTE_{theory} = MTE_{e-e} + MTE_{band}$... $MTE_{band} \ge \left(\frac{m_T^*}{m_0}\right) k_B T_e$; $m_T^* < m_0$ $MTE_{band} \leq 1.5k_BT_e$; $m_T^* > m_0$
- Emission processes not independent: BOTH emissions from virtual excited band states
- ⇒ Convolution in momentum space
 - ∴ Addition of p_T variances (i.e., MTEs)
 - ... Consistent with observation of a single Gaussian electron beam profile

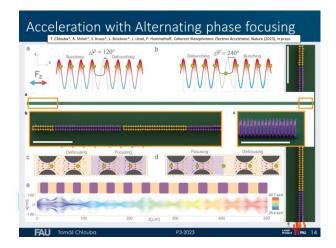
Brookhaven

psaha3@bnl.gov

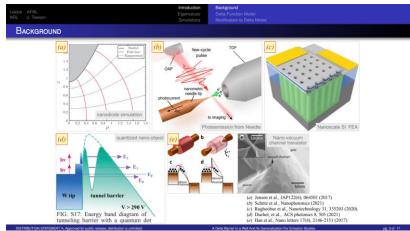
Theory

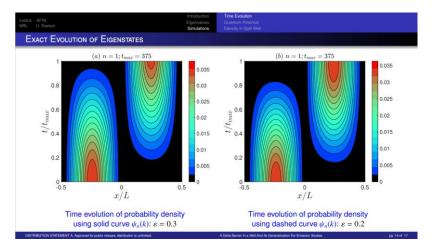








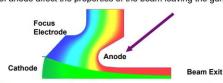








▶ Basic question: How do small changes in position or potential of anode affect the properties of the beam leaving the gun?



▶ Conventional solution: Trial and error. Do many simulations with different anode potentials or positions to understand sensitivities. Also leads to selecting the best (optimized) solution based on some performance metric.





Distribution A: Approved for Release. Distribution is unlimited



Mean Displacement: 2D Axisymmetric Electron Gun

- Manufacturing sensitivity to AK-Gap axial offset

Results:

- The test case where the direct perturbation of a voltage change worked as expected.

- The reverse-beam case was oddly sensitive to changes.

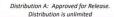






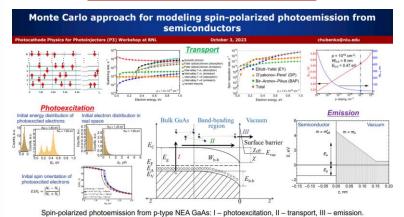


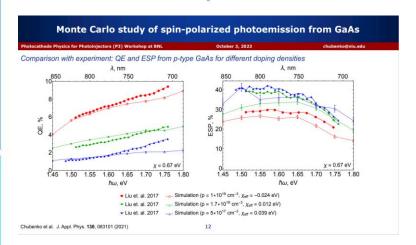




Spin polarized

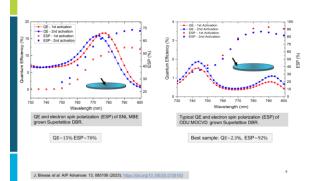




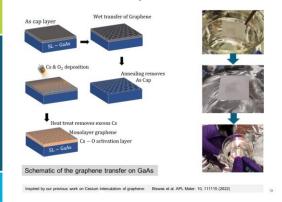




GaAs/GaAsP superlattice DBR: QE, ESP



Transfer of Graphene onto GaAs

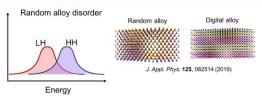


Strained superlattice InAlGaAs/AlGaAs spin-polarized photocathodes implemented with both random and digital alloying

Aaron Engel¹, Marcy Stuzman², Jason Dong¹, Christopher Palmstrøm^{1,3}



Approach to reduce alloy disorder and improve uniformity





- · Digital alloying minimizes random alloy disorder
- · Broader high spin-polarization window
- · InAlGaAs digital alloys (near lattice matched to InP) should have
- · Better optical emission than random allovs
- · Better uniformity than random alloys



D. Song,...,Y.T. Lee J. Cryst. I. J. Fritz, J. F. Klem, M. J. Hafich, A. J.

C.S. Wang,..., A.C. Gossard, L.A. Howard Appl. Phys. Lett. 66, 2825 (1995) Coldren J. Cryst. Growth 277, 13 (2005)

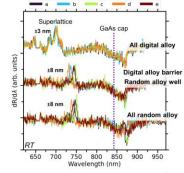
Digital alloys improve uniformity for cavity resonance

- · Uniformity problematic in the past
- · Essential for resonant cavity with DBR
- · Broader features in digital alloy well
- · Digital alloy well + barrier makes significant improvement to uniformity



MATERIALS

Uniformity is adequate... now DBRs





Spin polarized



Cornell Laboratory for Accelerator-based Sciences and Education (CLASSE)

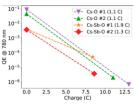
Polarized Photocathodes for **Future Applications**

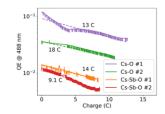
Samuel J. Levenson Photocathode Physics for Photoiniectors Workshop Stony Brook, NY, USA October 5, 2023

Alternative NEA Coatings: Cs-Sb-O High current studies in the HERACLES beamline

- · No improvement in the charge lifetime at 488 nm, modest improvement at 780 nm
- . Sb improves chemical poisoning, ion back bombardment less so
- · Relative QE change smaller for Cs-Sb-O samples

Now, we want to optimize the thicknesses of these activation layers w.r.t. operational lifetime Current project: HERACLES runs with varying activation layer thicknesses

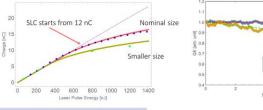


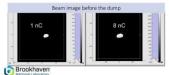


- J. K. Bae, M. Andorf, A. Bartnik, A. Galdi, L. Cultrera,
- J. Maxson and I. Bazarov, AIP Adv., 2022, 12, 095017.



Bunch Charge and Cathode Lifetime





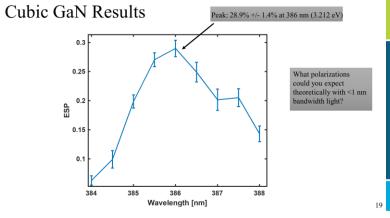
- •Bulk GaAs with 785 nm polarized laser; Gun operates at 300 kV.
- •7.5 nC bunch charge polarized beam, 5000 pulses/s ~37.5 uA; ·With anode bias, we didn't observe QE drop. Without anode bias 1/e lifetime is 63 hrs. Dominated by the
 - outgassing from FC. E. Wang et.al PHYSICAL REVIEW ACCELERATORS AND BEAMS 25, 033401 (2022) 12

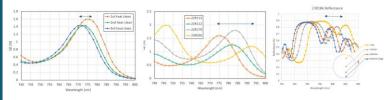
Anode +3000V; τ=-1000 hrs

Anode 0V; r=63 hrs

Third figure is provided by A. Master

DBR samples resonance frequency variation





- · Vary the DBR resonance frequency(~20 nm):
 - o Iterations of heat clean and QE (same sample, same location, same temperature)
 - o Different samples

Brookhaven

o Same sample on different location

A Dynasil Company



Strained Layer Epitaxy CdTe/ZnTe

$Cd_{(1-x)}Zn_{x}Te$	Lattice Parameter (A)	% Strain
x=0 (CdTe)	6.48	0.00
x=0.1	6.451	0.45
x=0.2	6.422	0.90
x=0.3	6.371	1.68
x=0.4	6.298	2.81
x=0.5	6.25	3.55
x=0.6	6.218	4.04
x=0.7	6.183	4.58
x=0.8	6.128	5.43
x=0.9	6.051	6.62
x=1 (ZnTe)	6.1	5.86

Preliminary investigation for CdTe indicate:

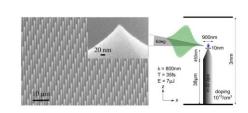
- · 1.5eV bandgap permits excitation at 780nm;
- · ~2% strain has the necessary valence band splitting (0.1eV) to realize 100% polarization;
- · p-type can be achieved with antimony/arsenic doping;
- · has excellent absorption and spin relaxation time;
- · negative electron affinity has been demonstrated with Cs-O surface treatments.

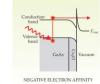
What is needed:

- DFT calculations of band structure
- Monte Carlo simulations of transport properties

rmd.dvnasil.com

GaAs Nanotip Activation





- · ALD can be used for nanotip activation
- · RMD does have a proprietary Cs ALD precursor (metal organic)
- On-site Cs2O surface activation is possible

Swanwick, Michael E., et al. "Nanostructured ultrafast silicon-tip optical field-emitter arrays." Nano letters 14.9 (2014): 5035-5043



P3 2025: Welcome to the Desert

*ARIZONA STATE UNIVERSITY





Beus CXFEL Lab

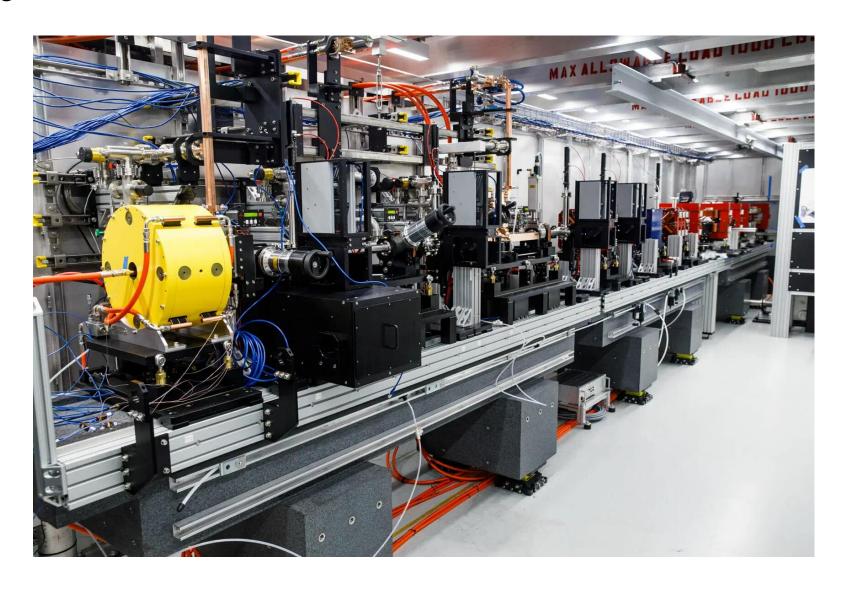


10 MeV linac based Incoherent Compton X-ray source – first x-rays

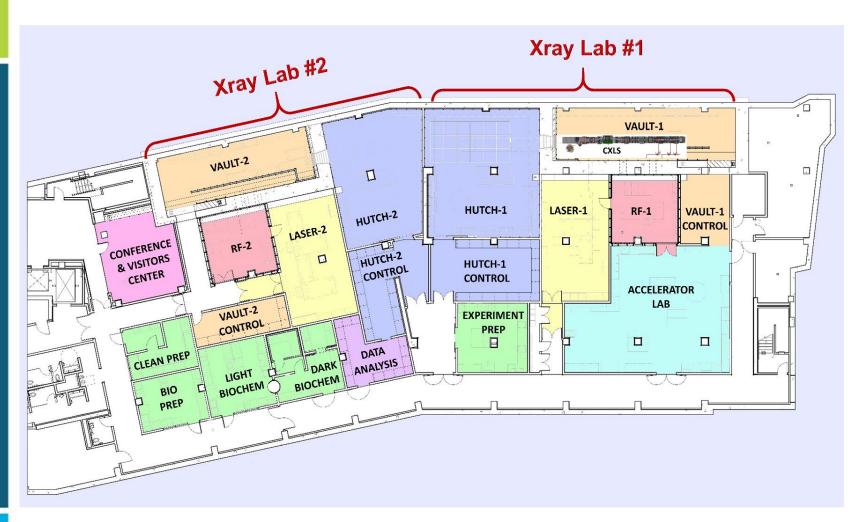
\$90M NSF X-ray facility

Coherent X-ray Compton source research





Beus CXFEL Lab



Unique ultra-stable photoinjector facility

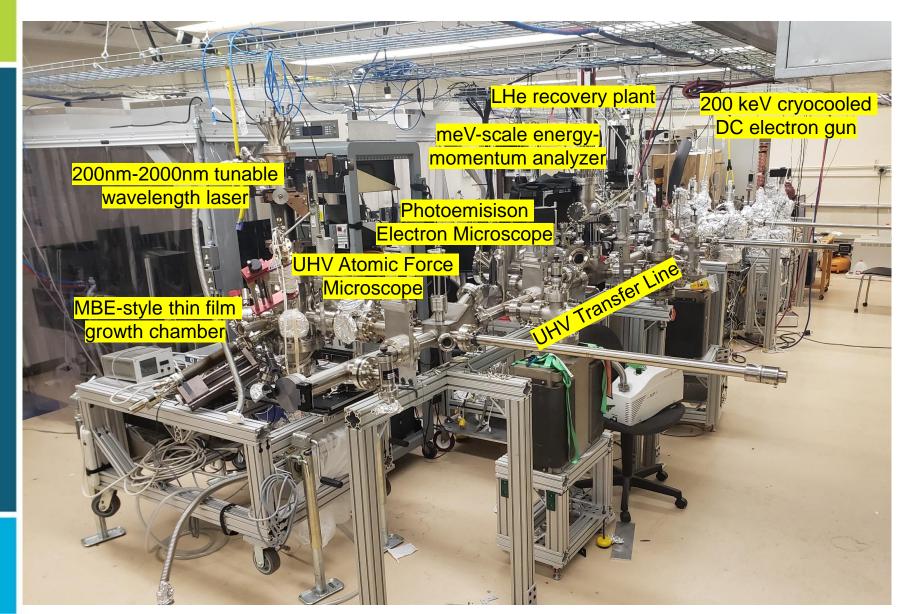
2 vaults:

- Vibration free
- EM free
- Radiation shielded

Large Dedicated X-ray 'hutches' + prep labs



Photoemission and Bright Beams Lab



UHV thin film growth

Atomic scale surface characterization

Fundamental photoemission physics

Electron Beam Physics

Ultrafast Electron Diffraction (soon)

John M. Cowley Center for High Resolution Electron Microscopy

Nion UltraSTEM100



Sub-Å spatial resolution Sub-10 meV energy spread

JEOL ARM200F





- Vibrational Spectroscopy in EM
- Aloof EM
- Lorentz EM ...



Nanofab and Eyring Materials Center









Large University Campus

- Strong expertise in nanoscience and ultrafast physics and upcoming accelerator physics
- >60,000 students







P3 2025: See you next year

But will it still be P3...?









Thank you!!!

Thanks to all authors from which I have taken slides and to Siddharth Karkare who provided information on next P3 workshop at ASU

Supported by DE-SC00012704











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