Probing the Quantum Vacuum at the High-Intensity Frontier

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Talk title vs. conference title?

Many-particle systems under extreme conditions

A view on the classical vacuum

A view on the classical vacuum

Absence of anything

The birth of the quantum vacuum

Field quantization
 Dirac's "theory of the positron"
 Heisenberg-Euler theory

. . .

(BORN, HEISENBERG, JORDAN 1925)

(DIRAC 1928)

(EULER,KOCKEL 1935; EULER 1936; HEISENBERG-EULER 1936)

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Two light quanta can ... be transformed
into two other light quanta through the
- so to speak - virtual possibility of
pair building ....
The question ... is principally of
great importance, ... problem of how
Maxwell's equations are modified by
Dirac's theory.
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Heisenberg's evaluation of Euler's thesis

(TRANSLATION: W. DTTRICH 2014)









$$\lambda_{\rm C} = \frac{\hbar}{mc} \approx 3.9 \cdot 10^{-13} \,\mathrm{m}$$

 $\tau_{\rm C} = \frac{\lambda_{\rm C}}{c} \approx 1.3 \cdot 10^{-21} \,\mathrm{s}$

 \triangleright Heisenberg's uncertainty principle, e.g. $\Delta x \cdot \Delta p \geq \frac{\hbar}{2}$



Quantum field theory: possibility of everything

 \triangleright Heisenberg's uncertainty principle, e.g. $\Delta x \cdot \Delta p \geq \frac{\hbar}{2}$



Quantum vacuum $\hat{=}$ Many-particle system \checkmark

Probing the quantum vacuum

▷ Quantum electrodynamics (QED): e.g., light-by-light scattering*

www.



 \implies violation of Maxwell's superposition principle in vacuo

* 3) P. Debye, in einer mündlichen Diskussion mit Herrn Prof. Heisenberg.

Probing the quantum vacuum

▷ Heisenberg-Euler effective action

Folgerungen aus der Diracschen Theorie des Positrons.

Von W. Heisenberg und H. Euler in Leipzig.

Mit 2 Abbildungen. (Eingegangen am 22. Dezember 1935.)

Aus der Diracschen Theorie des Positrons folgt, da jedes elektromagnetische Feld zur Paarerzeugung neigt, eine Abänderung der Maxwellschen Gleichungen des Vakuums. Diese Abänderungen werden für den speziellen Fall berechnet, in dem keine wirklichen Elektronen und Positronen vorhanden sind, und in dem sich das Feld auf Strecken der Compton-Wellenlänge nur wenig ändert. Es ergibt sich für das Feld eine Lagrange-Funktion:

$$\begin{split} \mathfrak{L} &= \frac{1}{2} \left(\mathfrak{G}^{\mathfrak{g}} - \mathfrak{Y}^{\mathfrak{g}} \right) + \frac{e^3}{h \, c} \int\limits_{0}^{\infty} e^{-\eta} \frac{\mathrm{d} \, \eta}{\eta^3} \left\{ i \, \eta^2 \left(\mathfrak{G} \, \mathfrak{Y} \right) \right\} \cdot \frac{\cos \left(\frac{\eta}{|\mathfrak{G}_k|} \, \sqrt{\mathfrak{G}^2 - \mathfrak{Y}^2 + 2 \, i \left(\mathfrak{G} \, \mathfrak{Y} \right)} \right) + \mathrm{konj}}{\cos \left(\frac{\eta}{|\mathfrak{G}_k|} \, \sqrt{\mathfrak{G}^2 - \mathfrak{Y}^2 + 2 \, i \left(\mathfrak{G} \, \mathfrak{Y} \right)} \right) - \mathrm{konj}} \\ &+ |\mathfrak{G}_k|^2 + \frac{\eta^3}{3} \left(\mathfrak{Y}^2 - \mathfrak{G}^2 \right) \right\} \cdot \\ \left(\mathfrak{G}_k \, \mathfrak{Y} \, \operatorname{Kraft} \, \mathrm{auf} \, \mathrm{das} \, \mathrm{Elektron.} \\ \left| \mathfrak{G}_k \, \mathfrak{g}_k \, \mathfrak{g}_k - \mathfrak{g}_k^2 - \mathfrak{g}_k^2 \right|^2 = \frac{m^2 \, c^3}{e \, \hbar} = \frac{1}{\pi 137^{\alpha}} \frac{e}{(e^3/m \, c^3)^2} = \, \mathrm{sKritische} \, \mathrm{Feldstärke}^*. \end{split} \right) \end{split}$$

Inre Entwicklungsglieder für (gegen $|\xi_k|$) kleine Felder beschreiben Prozesse der Streuung von Licht an Licht, deren einfachstes bereits aus einer Störungsrechnung bekannt ist. Für große Felder sind die hier abgeleiteten Feldgeleichungen von den Maxwell schen sehr verschieden. Sie werden mit den von Born vorgeschlagenen verglichen. (Heisenberg, Euler 1936; Weisskopf 1936; Schwinger 1951)



- strong field nonlinearities
- ⊳ quantum vacuum

 \simeq medium

critical field: extreme conditions!

$$E_{
m cr}\simeq 1.3 imes 10^{18} V/m$$

▷ vacuum instability:

pair production

Probing the quantum vacuum

Heisenberg-Euler effective action

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Ihre Entwicklungsglieder für (gegen [4],]) kleine Felder beschreiben Prozense der Streuung von Licht an Licht, deren einfachstes bereits aus einer Störungsrechnung bekannt ist. Für große Felder sind die hier abgeleisten Feldgleichungen von den Max well schen zehr verschieden. Sie werden mit den von Born vorgeschlagenen verglichen. (HEISENBERG, EULER 1936; WEISSKOPF 1936; SCHWINGER 1951)

 $1 2 \times X 2$

$$S_{\text{eff,HE}}[F] = \sum_{n=0}^{\infty} \bigcup_{i=1}^{2} \bigcup_{j=1}^{2} \bigcup_{i=1}^{3} \bigcup_{j=1}^{3} \bigcup_{j=1}^{3} \bigcup_{j=1}^{3} \bigcup_{i=1}^{3} \bigcup_{j=1}^{3} \bigcup_{j=1}^{$$

conceptual cornerstone of QFT

many predicted phenomena still await their discovery

high energy/momentum

strong fields / high intensity

high energy/momentum



strong fields / high intensity

high energy/momentum





strong fields / high intensity

high energy/momentum





Example: Light Propagation in a *B* field.

▷ quantum Maxwell equation

(Heisenberg,Euler'36;Weisskopf'36)

$$0 = \partial_{\mu} \left(F^{\mu\nu} - \frac{1}{2} \frac{8}{45} \frac{\alpha^2}{m^4} F^{\alpha\beta} F_{\alpha\beta} F^{\mu\nu} - \frac{1}{2} \frac{14}{45} \frac{\alpha^2}{m^4} F^{\alpha\beta} F_{\alpha\beta} \tilde{F}^{\mu\nu} \right)$$

Example: Light Propagation in a *B* field.

 \triangleright quantum Maxwell equation for a "light probe" $f^{\mu\nu}$

$$0 = \partial_{\mu} f^{\mu\nu} - \frac{8}{45} \frac{\alpha^2}{m^4} F_{\alpha\beta} F^{\mu\nu} \partial_{\mu} f^{\alpha\beta} - \frac{14}{45} \frac{\alpha^2}{m^4} \tilde{F}_{\alpha\beta} \tilde{F}^{\mu\nu} \partial_{\mu} f^{\alpha\beta}$$

Phase and group velocity

$$v_{\parallel} \simeq 1 - \frac{14}{45} \frac{\alpha^2}{m^4} B^2 \sin^2 \theta_B$$
$$v_{\perp} \simeq 1 - \frac{8}{45} \frac{\alpha^2}{m^4} B^2 \sin^2 \theta_B$$



(Toll'54) (Baier,Breitenlohner'67;Narozhniy'69) (Adler'71)

(DITTRICH,HG'98; DITTRICH,HG'00)

→ magnetized quantum vacuum induces birefringence

▷ detection schemes: PVLAS, BMV, Q&A, OSQAR

cf. birefringence in uniaxial cristals

⊳ e.g., calcite



An upcoming experiment

Probing the quantum vacuum at extreme conditions

vacuum birefringence of a probe wave induced by a high-intensity pulse (pump) Jena SFB-TR18 scheme: (Heinzl, Liesfeld, Amthor, Schwoerer, Sauerbrey, Wipf 2006)



Vacuum Birefringence: a flagship experiment for

HI JENA HELMHOLTZ

▷ Theory of pulse collisions

(Karbstein, HG, Reuter, Zepf'15)

(KARBSTEIN, SUNDQVIST'16; KARBSTEIN, MOSMAN'21)

High definition X-ray polarimetry

(Marx et al. 2011,Schulze 2015& 2018)



(Schlenvoigt, Heinzl, Schramm, Cowan, Sauerbrey'15)



Theory of pulse collisions: discernible Photons

▷ Photons above background

(KARBSTEIN, BLINNE, HG, ZEPF 2019)

$$rac{dN_{
m signal}}{d\Omega} \geq \mathcal{P} \; rac{dN_{
m background}}{d\Omega}$$

... useful tool to identify sweet spots in large parameter space

▷ Signal photons from vacuum emission picture

(KARBSTEIN, SHAISULTANOV'15)

$$dN_{
m signal} = rac{d^3k}{(2\pi)^3} |\mathcal{S}_{VE}(ec{k})|^2$$

applicable for any pump/probe pulse shapes amenable to numerical simulations (BLINNE,HG,KARBSTEIN KOHLFÜRST,ZEPF'15)

Example: colliding PW-class pulses

 \triangleright collision under $\vartheta_2 = 135^{\circ}$

(HG,KARBSTEIN,KOHLFÜRST'17)



"Canonical approach": background suppression by X-ray polarimetry

▷ high-definition x-ray polarimetry using Bragg reflections

(Marx et al. 2011,Schulze 2015ず 2018)



polarization purity (HED@XFEL) 2022: $P < 8 \times 10^{-11}$

(SCHULZE ET AL. 2022)

"3-pulse collision scheme"



🙂 sizable signal

- $\bigcirc \omega_{\text{signal}} \simeq |\omega_1 \pm \omega_2 \pm \omega_3|$
- 🙂 angular separation of signal
- " prediction for signal divergence
 - experimentally more challenging

XFEL+optical:

(King, Hu, Shen'18)





"Dark-field scheme"

▷ annular probe beam

(Karbstein,Mosman'20)

- low-background hole in forward direction
- quantum signal maximum on axis





▷ proof of concept experiment: $\mathcal{P} = 10^{-8}$ (Karbstein,Ullmann,Mosman,Zepf'22)

 \implies HIBEF design study

Vacuum Birefringence @ HIBEF/HED collaboration

> Dark-field scheme

(MICHAL ŠMÍD, ET AL. 2023)



▷ Simulated shadow factor: $P = 1.6 \times 10^{-12}$ ▷ road map:

- Priority access project 5438: beam time in March 2024, dark-field test X-ray only
- Autumn 2024: combining X-ray with Relax lasers,
- \implies full experiment 2025(?)

Why quantum vacuum physics?

Quantum effective actions:

exploring the Heisenberg-Euler paradigm measuring effective field theories near thresholds

▷ vacuum birefringence:

fundamental effect of QED

 \sim (g – 2), Lamb shift, etc.

precision explores unknown territory

$$\Delta v \simeq \frac{6}{45} \left(1 + \frac{25}{4} \frac{\alpha}{\pi} + \dots \right) \frac{\alpha^2}{m^4} I$$

▷ quantum vacuum as a building block





Alfvén, magneto-acustic modes nonlinear/shock-wave propagation higher-harmonic generation self-focusing high-intensity metrology

▷ laboratory tests of astrophysical strong-field environments





(Mignani et al. 2016) (Capparelli et al. 2017)

polarization properties of pulsars, magnetars, etc.

▷ search for "new physics"

all degrees of freedom contribute to the quantum vacuum



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▷ sensitivity to weakly interacting, light particles axions, dark matter candidates ...

Going to extremes

Extreme conditions: Strong-field limit?

Leading-log resummation of perturbation theory 1PI Effective action in magnetic fields

$$\Gamma[B] o \infty, \quad ext{for} \ rac{lpha}{3\pi} \ln rac{eB}{m^2} o \mathcal{O}(1)$$

LL resummation translates

high-*p* Landau pole \rightarrow high-*B*

 \implies artefact of single-scale RG?



(RITUS'75;DITTRICH,REUTER'85)

Extreme conditions: Strong-field limit?

▷ singularities for large electric field expected

and observed already at 1-loop order: $Im\Gamma[E] \sim \alpha (eE)^2 e^{-\frac{m^2}{eE}}$

(SAUTER'31, SCHWINGER'51)

BUT (constant) magnetic fields do not transfer energy to charged particles

 \implies no heuristic physical reason to translate $ext{high-}p \ o \ ext{high-}B$



Hypothesis: Quantum scale symmetry

(Gell-Mann,Low'54; Weinberg'76; Wetterich'19)

▷ If strong-field limit exists, there will be a regime where

$$eB \gg m_i^2, \qquad i=e,\mu,\tau,\pi^{\pm},\ldots$$

▷ This suggests that action density

$$rac{\Gamma}{\Omega} \hspace{.1in}
ightarrow \hspace{.1in}$$
 scale invariant form in units of an RG scale $\hspace{.1in} k \ggg m_i$

Does such an RG fixed point ("fixed function") exist as a solution to the RG equations?

Functional RG flow equation

▷ Flow equation for the 1PI effective action

$$k\partial_k\Gamma_k = \frac{1}{2} \operatorname{Tr} k\partial_k R_k (\Gamma_k^{(2)} + R_k)^{-1}$$



▷ We explore the ansatz (in dim'less quantities)

$$\Gamma = \int_{x} w_{k}(\mathcal{F}, \mathcal{G}^{2}), \quad \mathcal{F} = \frac{1}{4} F_{\mu\nu} F^{\mu\nu}, \ \mathcal{G} = \frac{1}{4} F_{\mu\nu} \widetilde{F}^{\mu\nu}$$

Evidence for quantum scale symmetry: $w_k(\mathcal{F}, 0)$ needs to exists globally for $\mathcal{F} > 0$ + RG scale independence

Functional RG flow equation

▷ fixed-function equation

(LAPORTE,LOCHT,PEREIRA,SAUERESSIG'23; HG,SCHIRRMEISTER'23)

$$k\partial_k w_k \stackrel{!}{=} 0 = -4w_k + (4+\eta)(w'_k \mathcal{F} + 2\dot{w}_k \mathcal{G}^2) - \frac{1}{32\pi^2} \int_{\mathcal{Y}} y^2 (2y^2r' + \eta r)Y_k$$

Solutions depend parametrically on photon anomalous dimension η \vartriangleright Diagrammar:



Functional RG flow equation

▷ fixed-function equation

(LAPORTE,LOCHT,PEREIRA,SAUERESSIG'23; HG,SCHIRRMEISTER'23)

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Solutions depend parametrically on photon anomalous dimension η \triangleright Diagrammar:



Global fixed function



scale-invariant solution exists globally

 \implies Check: perturbative small- η expansion consistent with one-loop strong-field limit.

Conclusions

▷ Probing the quantum vacuum at the high-intensity frontier

understanding the ground state of nature

"The vacuum is a medium"

exploring uncharted territory of fundamental physics

... high amplitude vs. high energy

- first discoveries and emerging opportunities
- \implies Golden opportunities for CPA lasers + XFEL!



QED quantum vacuum under ultra-extrem conditions first nonperturbative evidence for stability (global existence)